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**Soundness and completeness of quantum
root-mean-square errors**

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Classical Root-Mean-Square Error

- **Definition.** If a quantity A is indirectly measured by measuring another quantity M with the joint probability distribution $\mu(a, m)$ of A and M , the root-mean-square error $\varepsilon_G(\mu)$ of the measurement of A is defined by

$$\varepsilon_G(\mu) = \left(\sum_{a,m} (m - a)^2 \mu(a, m) \right)^{1/2}.$$

- **Theorem.** $\varepsilon_G(\mu) = 0$ if and only if $\mu\{M = A\} = 1$, where

$$\mu\{M = A\} = \sum_{(a,m): m=a} \mu(a, m).$$

Universal Models for Quantum Measurements

- **Definition.** $M = (\mathcal{K}, \xi, U, M)$: Measuring Process for the system described by $\mathcal{H} \Leftrightarrow$

\mathcal{K} = a Hilbert space, modeling the state space of the probe

ξ = a unit vector on \mathcal{K} , modeling the initial state of the probe

U = a unitary on $\mathcal{H} \otimes \mathcal{K}$, modeling the measuring interaction

M = a self-adjoint operator on \mathcal{K} , modeling the meter observable

- For any $A, B \in \mathcal{O}(\mathcal{H})$, the measuring process M from time 0 to τ determines

$$A(0) = A \otimes I, \quad B(0) = B \otimes I, \quad M(0) = I \otimes M,$$

$$A(\tau) = U^\dagger A(0)U, \quad B(\tau) = U^\dagger B(0)U, \quad M(\tau) = U^\dagger M(0)U.$$

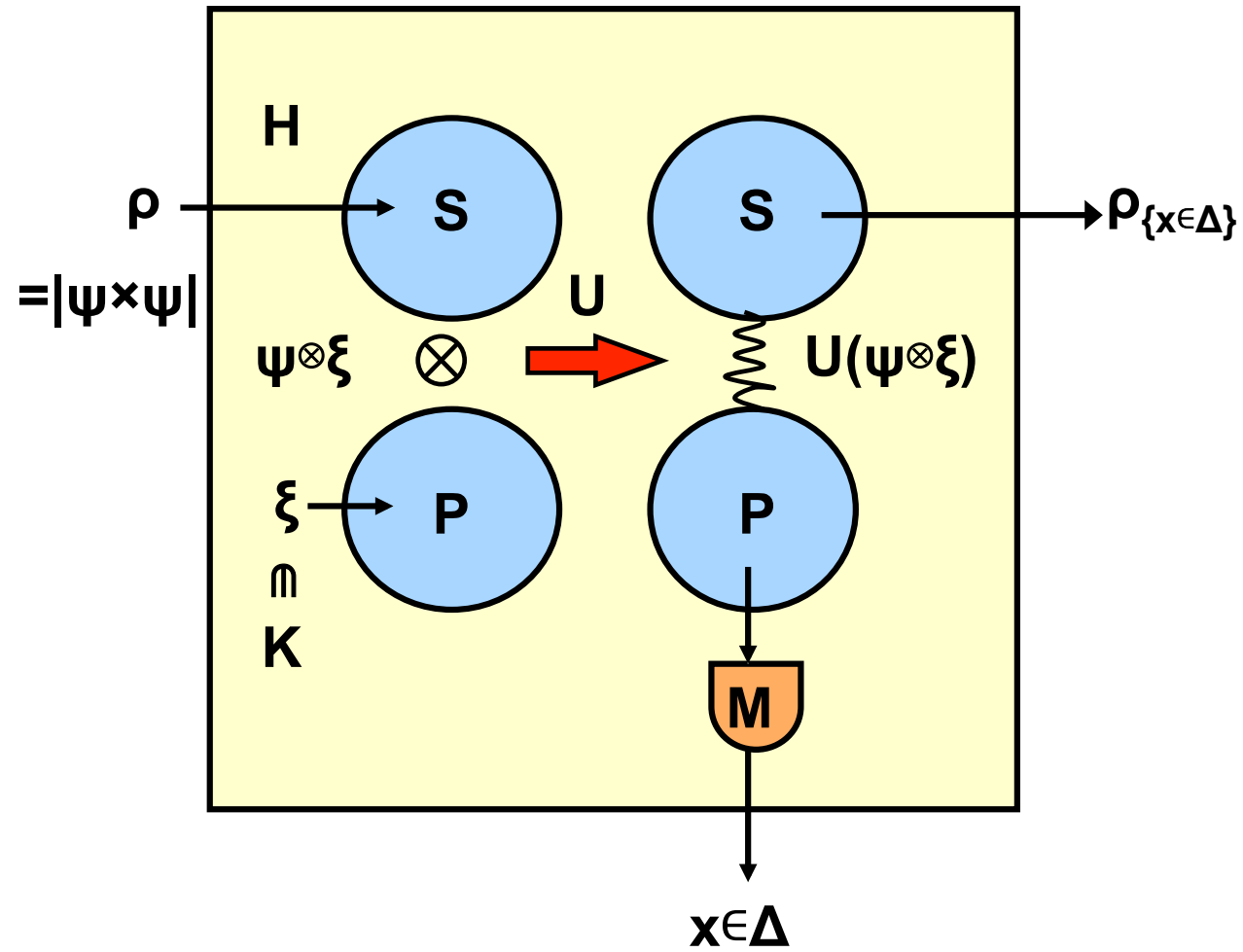


Figure 1: Measuring Preprocess

- **When is the measurement $M = (\mathcal{K}, \xi, U, M)$ of A in ψ considered accurate? —The joint probability distribution μ of $M(\tau)$, $A(0)$ exists and satisfies $\mu\{M(\tau) = A(0)\} = 1$.**

State-Dependent Commutativity

- **Definition:** X and Y commute in Ψ ($X \leftrightarrow_{\Psi} Y$) iff $[f(X), g(Y)]\Psi = 0$ for all polynomials $f(X), g(Y)$.

- **Definition:** A joint probability distribution (JPD) of observables X, Y in state Ψ is a probability distribution $\mu_{\Psi}^{X,Y}(x, y)$ on \mathbb{R}^2 satisfying

$$\langle \Psi | f(X, Y) | \Psi \rangle = \sum_{x,y} f(x, y) \mu_{\Psi}^{X,Y}(x, y)$$

for any polynomial $f(X, Y)$.

- **Theorem:** There exists a JPD of A, B in Ψ .

$$\Leftrightarrow A \leftrightarrow_{\Psi} B$$

$\Leftrightarrow \Psi$ is a superposition of common eigenstates of A and B , i.e.,

$$\Psi = \sum_{x,y} c_{x,y} |A = x, B = y\rangle .$$

Weak Joint Distribution

- **Definition:** The weak joint distribution (WJD) $\nu_{\Psi}^{A,B}(x, y)$ of observables A, B in state Ψ is defined by

$$\nu_{\Psi}^{A,B}(x, y) = \langle \Psi | P^A(x) P^B(y) | \Psi \rangle .$$

- **Remark:** The WJD can be measured by weak measurement and post-selection, i.e.,

$$\nu_{\Psi}^{A,B}(x, y) = \langle P^B(y) \rangle_{w, A=x, \Psi} \mu_{\Psi}^A(x).$$

Quantum Perfect Correlation

- **Definition:** A and B are perfectly correlated in Ψ ($A =_{\Psi} B$)
 $\Leftrightarrow A \leftrightarrow_{\Psi} B$ and the JPD $\mu_{\Psi}^{A,B}$ satisfies

$$\sum_x \mu_{\Psi}^{A,B}(x, x) = 1.$$

- **Theorem (MO 2005):**

$$\begin{aligned} A =_{\Psi} B &\Leftrightarrow \nu_{\Psi}^{A,B}(x, y) = 0 \quad \text{if } x \neq y \\ &\Leftrightarrow \left\langle P^B(y) \right\rangle_{w, A=x, \Psi} = \delta_{x,y}. \end{aligned}$$

- **Remark:** Perfect correlation is experimentally accessible.
- **Theorem (MO 2005):** The relation $=_{\Psi}$ is an equivalence relation among observables. In particular, if $A =_{\Psi} B$ and $B =_{\Psi} C$, then $A =_{\Psi} C$.

State-Dependent Accurate Measurements

- **Definition (MO 2005):** A measuring process $M = (\mathcal{K}, \xi, U, M)$ *accurately measures* A in $\psi \Leftrightarrow$

$$A(0) =_{\psi \otimes \xi} M(\tau).$$

- **Remark:** The above condition is operationally accessible, since it is equivalent to

$$\sum_{x \neq y} |\nu_{\psi \otimes \xi}^{A(0), M(\tau)}(x, y)| = 0.$$

- **Theorem** A measuring process $M = (\mathcal{K}, \xi, U, M)$ *accurately measures* A in $\psi \Leftrightarrow$

$$\psi \otimes \xi = \sum_x c_x |A(0) = x, M(\tau) = x\rangle,$$

where $|c_x|^2 = |\langle \psi | A = x \rangle|^2$.

Noise Operator based Quantum Root-Mean-Square Error

- **Definition.** For any measuring process $M = (\mathcal{K}, \xi, U, M)$ the noise operator $N(A)$ for measuring A is defined by

$$N(A) = M(\tau) - A(0).$$

- The NO based QRMSE for measuring A in Ψ is defined by

$$\varepsilon_{NO}(A, \psi) = \langle N(A)^2 \rangle^{1/2} = \|[M(\tau) - A(0)]\psi \otimes \xi\|.$$

- **Theorem (Lund-Wiseman 2010)** The NO based QRMSE $\varepsilon_{NO}(A, \psi)$ can be measured by weak measurement and post-selection

$$\varepsilon_{NO}(A, \psi)^2 = \sum_{a,m} (m - a)^2 \text{Re } \nu_{\psi \otimes \xi}^{M(\tau), A(0)}(m, a).$$

Requirements for Quantum RMS Errors

- (i) **Device-independent definability:** The error measure should be definable by the POVM of the measuring process, the observable to be measured, and the state of the object.
- (ii) **Correspondence principle:** The error measure should be identical with the classical rms error if the joint probability distribution of $M(\tau)$ and $A(0)$ exists.
- (iii) **Soundness:** The error measure should take the value zero for any accurate measurements.
- (iv) **Completeness:** The error measure should take the value zero if and only if the measurement is accurate.

Device-Independent Definability

- The NO based QRMSE ε_{NO} , satisfies the device-independent definability.
- The POVM of M: $\Pi(x) = \langle \xi | E^{M(\tau)}(x) | \xi \rangle$
- Moment operator of POVM Π : $m^{(n)}(\Pi) = \sum_{x \in \mathbb{R}} x^n \Pi(x)$
- $\varepsilon_{NO}(A, \psi)$ satisfies

$$\varepsilon_{NO}(A, \psi)^2 = \text{Re} \langle \psi | m^{(2)}(\Pi) - 2m(\Pi)A + A^2 | \psi \rangle .$$

Correspondence Principle

- The NO based QRMSE ε_{NO} satisfies the correspondence principle.
- If $M(\tau)$ and $A(0)$ commute in $\psi \otimes \xi$, there exists the joint probability distribution $\mu(a, m) = \mu_{\psi \otimes \xi}^{A(0), M(\tau)}(a, m)$ and we have

$$\varepsilon_{NO}(A, \psi)^2 = \sum_{a, m} (m - a)^2 \mu(a, m).$$

- Note that any error notions having been proposed based on the distance between the probability distributions do not satisfy the Correspondence Principle.

Soundness

- The NO based QRMSE ε_{NO} satisfies the soundness condition.
- If $A(0) =_{\Psi \otimes \sigma} M(\tau)$ then

$$\mu_{\psi \otimes \xi}^{M(\tau), A(0)}(m, a) = 0 \quad \text{for } m \neq a,$$

and hence

$$\varepsilon_{NO}(A, \psi)^2 = \sum_{a, m} (m - a)^2 \mu_{\psi}^{M(\tau), A(0)}(m, a) = 0.$$

- The NO based QRMSE satisfies the device-independent-definability, the correspondence principle, and the soundness.

Locally Uniform Quantum Root Mean Square Error

- For any $t \in \mathbb{R}$, define

$$\varepsilon_t(A, \psi) = \varepsilon_{NO}(A, e^{-itA}\psi).$$

- The locally uniform rms error is defined by

$$\bar{\varepsilon}(A, \psi) = \sup_{t \in \mathbb{R}} \varepsilon_t(A, \psi).$$

- **Theorem:** (1) If $A(0)$ and $M(\tau)$ commute in $\psi \otimes \xi$, then

$$\bar{\varepsilon}(A, \psi) = \varepsilon_{NO}(A, \psi).$$

(2) $\bar{\varepsilon}$ satisfies all the requirements (i)–(iv).

(3) $\varepsilon_{NO}(A, \psi) \leq \bar{\varepsilon}(A, \psi)$.

(4) If $A(0)^2 = M(\tau)^2 = I$, then $\bar{\varepsilon}(A, \psi) = \varepsilon_{NO}(A, \psi)$.

(5) The relation

$$\bar{\varepsilon}(Q, \psi)\bar{\varepsilon}(P, \psi) \geq \frac{\hbar}{2}$$

is violated.

Example

$$A = \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}, \quad M = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}, \quad |\psi\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

with $\Pi(y) = P^M(y)$. Then we have

$$\varepsilon_{NO}(A, \Pi, \psi) = 0,$$

but the measurement is not accurate, since A and Π are not identically distributed as $\langle \psi | P^A(2) | \psi \rangle = 1/2$ but $\langle \psi | \Pi(2) | \psi \rangle = 0$.

We have

$$\varepsilon_t(A, \Pi, \psi) = 2|\sin t|, \quad \text{and} \quad \bar{\varepsilon}(A, \Pi, \psi) = 2, \quad (1)$$

despite of the relation $\varepsilon_{NO}(A, \Pi, \psi) = 0$, the relation $\bar{\varepsilon}(A, \Pi, \psi) = 2$ correctly indicate that the above measurement of A is not accurate.

Simultaneous Measurements

- Let $M = (\mathcal{K}, \xi, U, M)$ be a measuring process. For any function f define $f(M) = (\mathcal{K}, \xi, U, f(M))$
- **Definition.** A simultaneous measurement for A, B in Ψ is defined as a pair of measuring processes $(f(M), g(M))$.
- **Definition:** A simultaneous measurement $(f(M), g(M))$ for A, B in ξ is *accurate* iff

$$\begin{aligned} f(M(\tau)) &=_{\Psi \otimes \xi} A(0), \\ g(M(\tau)) &=_{\Psi \otimes \xi} B(0). \end{aligned}$$

Uncertainty Relations for Simultaneous Measurements

- **Definition.** The error $(\bar{\varepsilon}(A), \bar{\varepsilon}(B))$ of a simultaneous measurement $(f(M), g(M))$ for A, B in ξ is defined by $\bar{\varepsilon}(A) = \bar{\varepsilon}(A, \psi, f(M))$ and $\bar{\varepsilon}(B) = \bar{\varepsilon}(B, \psi, g(M))$.
- **Theorem.** $(f(M), g(M))$ is accurate if and only if $(\bar{\varepsilon}(A), \bar{\varepsilon}(B)) = (0, 0)$.
- **Theorem.** Let $C_{AB} = \frac{1}{2} | \langle \psi | [A, B] | \psi \rangle |$. The following relations hold
 - (i) $\bar{\varepsilon}(A)\bar{\varepsilon}(B) + \sigma(B)\bar{\varepsilon}(A) + \sigma(A)\bar{\varepsilon}(B) \geq C_{AB}$.
 - (ii) $\sigma(B)^2\bar{\varepsilon}(A)^2 + \sigma(A)^2\bar{\varepsilon}(B)^2 + 2\bar{\varepsilon}(A)\bar{\varepsilon}(B)\sqrt{\sigma(A)^2\sigma(B)^2 - C_{AB}^2} \geq C_{AB}^2$.