

Periodic solutions bifurcated from figure eight choreography: non-planar eight and non- symmetric eight

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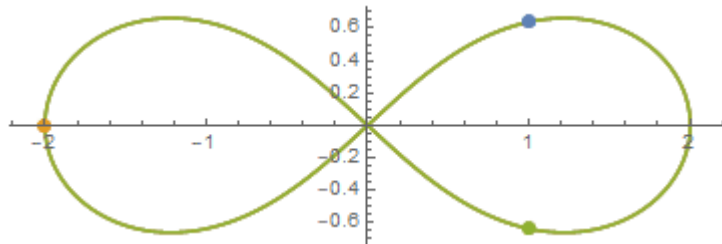
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ICIAM 2023 TOKYO

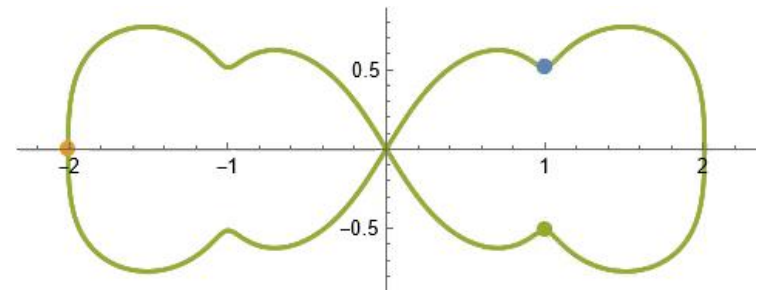
Bifurcations of figure eight choreography

We are investigating bifurcation of the figure eight choreography under homogeneous potential $-\frac{1}{r^a}$ by changing the power a , or Lenard-Jones(LJ) type potential $\frac{1}{r^{12}} - \frac{1}{r^6}$ by period T .



homogeneous potential ($a=1$)

$$-\frac{1}{r^a}$$

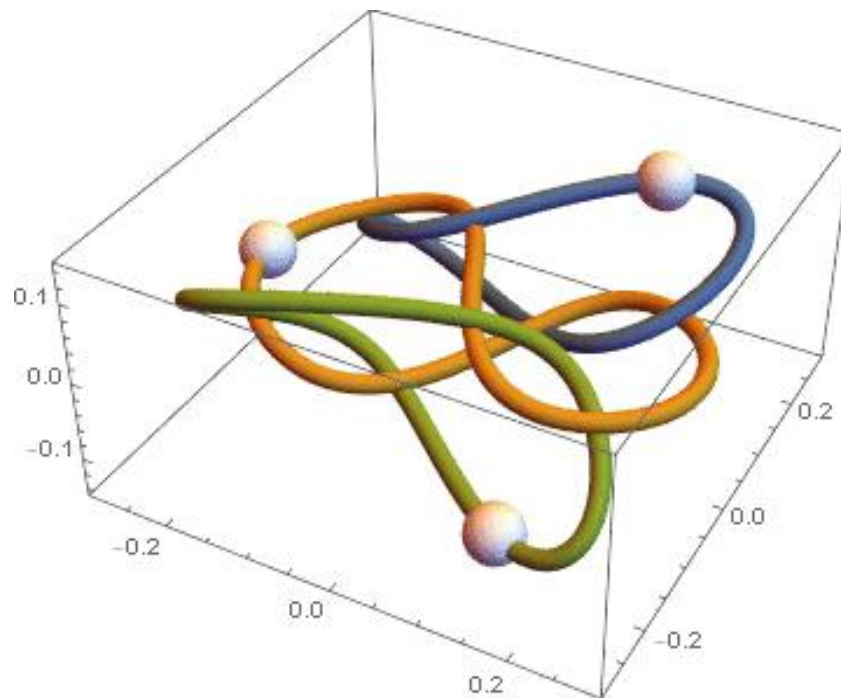


LJ potential ($T=50$)

$$\frac{1}{r^{12}} - \frac{1}{r^6}$$

Non-planar bifurcation solution

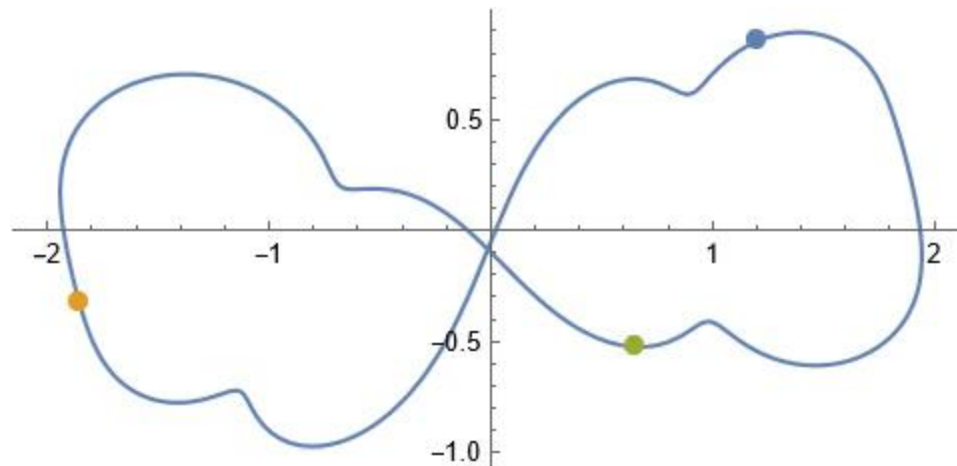
In this talk, we show two interesting bifurcation solutions, a non-planar bifurcation solution under homogeneous system, and



($a = 0.75$)

Figure-eight with no spatial symmetry

an eight-shaped choreography having no spatial symmetry under Lenard-Jones system.



Contents

My talk consists of the following eight parts:

1. Necessary condition of bifurcation
2. Symmetry of figure eight choreography
3. Symmetry group and irreducible representations
4. Bifurcations of figure eight to the planar motion
5. Bifurcations from planar bifurcations
6. Numerical calculations
7. Bifurcation of figure eight to the non-planar motion
8. Future work

1. Necessary condition of bifurcation

Necessary condition of bifurcation

At the bifurcation point,
eigenvalue λ of linear operator

$$H(q) = -\frac{d^2}{dt^2} + \frac{\partial^2 L}{\partial q^2}$$

crosses zero, where L is
Lagrangian for three bodies.

There, the eigenfunction ϕ
represents bifurcation solution
as

$$q^{(b)} \rightarrow q + h\phi$$

where $q^{(b)}$ is the bifurcation
solution, q the original solution,
and h a real number.

$$L(q, \dot{q}) = \frac{1}{2} \sum_i \dot{q}_i^2 - U(q)$$

$$U(q) = u(r_{12}) + u(r_{23}) + u(r_{31})$$

$$u(r) = -\frac{1}{ra} \quad \text{or} \quad -\frac{1}{r^6} + \frac{1}{r^{12}}$$

Necessary condition of bifurcation

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$$q^{(b)} \rightarrow q + h\phi$$

where $q^{(b)}$ is the bifurcation
solution, q the original solution,
and h a real number.

Since bifurcation solution
 $q^{(b)}$ and original solution q
satisfy eq. of motion:

$$\frac{d^2}{dt^2} q^{(b)} = \frac{\partial L}{\partial q} \Big|_{q^{(b)}},$$

$$\frac{d^2}{dt^2} q = \frac{\partial L}{\partial q}.$$

Thus, the difference $\phi =$
 $q^{(b)} - q$ satisfies

$$\frac{d^2}{dt^2} \phi = \frac{\partial L}{\partial q} \Big|_{q^{(b)}} - \frac{\partial L}{\partial q}$$

$$\rightarrow \frac{\partial^2 L}{\partial q^2} \phi$$

2. Symmetry of figure eight choreography

Eigenfunction of $H(q)$

Suppose the operator

$$H(q) = -\frac{d^2}{dt^2} + \frac{\partial^2 L}{\partial q^2}$$

commute with the all elements of the group $G = \{g_1, g_2, \dots, g_n\}$.

Then, the eigenspace of eigenvalue λ

$$H(q)\phi = \lambda\phi$$

is represented by unitary irreducible representations of G , especially, orthogonal irreducible representations since the eigenfunctions are real.

Symmetry of figure eight choreography

Actually, figure eight choreography $q(t)$ is invariant under the following three operations C, M, S :

- **C : choreographic**

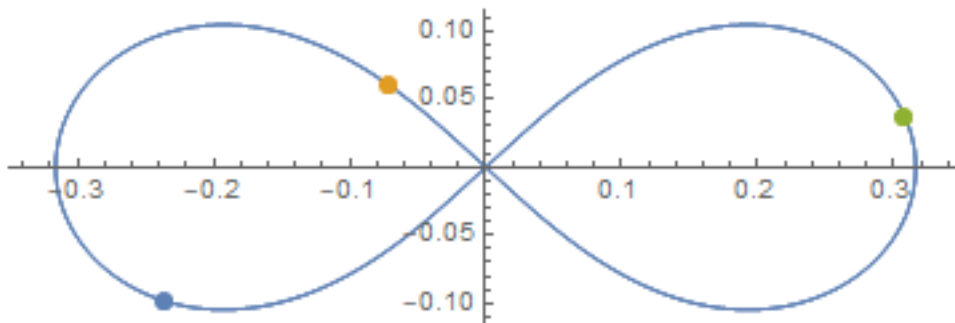
$$Cq(t) = \sigma q\left(t - \frac{T}{3}\right), \sigma: \text{cyclic permutation of bodies}$$

- **M : x inversion with time shift**

$$Mq(t) = \mu_x q\left(t - \frac{T}{2}\right), \mu_x: \text{inversion of } x \text{ coordinates}$$

- **S : space and time inversion with exchange two bodies**

$$Sq(t) = -\tau q(-t), \tau: \text{exchange of body 1 and 2}$$



Choreographic means all bodies move in the same orbit. Thus, time shift $T/3$ is cyclic permutation of bodies.

Symmetry of figure eight choreography

Actually, figure eight choreography $q(t)$ is invariant under the following three operations C, M, S :

- C : choreographic

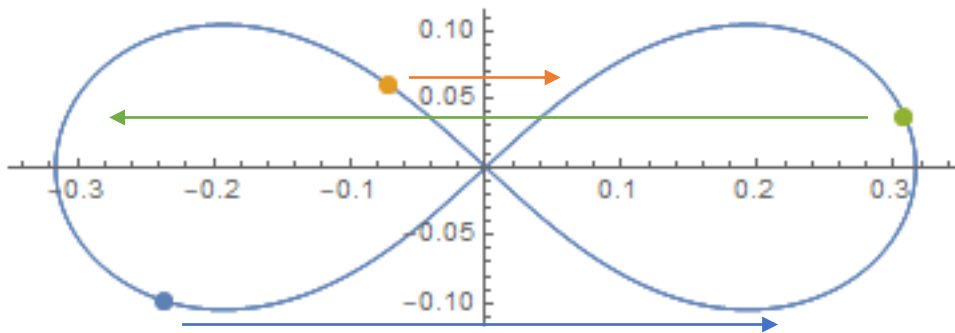
$$Cq(t) = \sigma q(t - \frac{T}{3}), \sigma: \text{cyclic permutation of bodies}$$

- M : x inversion with time shift

$$Mq(t) = \mu_x q(t - \frac{T}{2}), \mu_x: \text{inversion of } x \text{ coordinates}$$

- S : space and time inversion with exchange two bodies

$$Sq(t) = -\tau q(-t), \tau: \text{exchange of body 1 and 2}$$



M : x inversion with time shift $T/2$.

Symmetry of figure eight choreography

Actually, figure eight choreography $q(t)$ is invariant under the following three operations C, M, S :

- C : choreographic

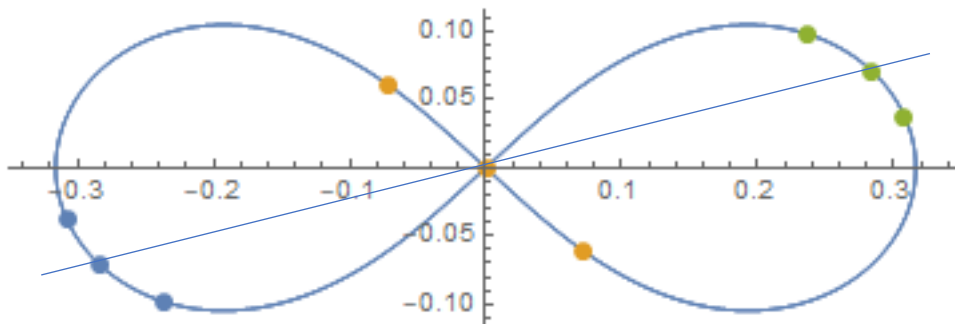
$$Cq(t) = \sigma q\left(t - \frac{T}{3}\right), \sigma: \text{cyclic permutation of body}$$

- M : x inversion with time shift

$$Mq(t) = \mu_x q\left(t - \frac{T}{2}\right), \mu_x: \text{inversion of } x \text{ coordinates}$$

- **S : space and time inversion with exchange body**

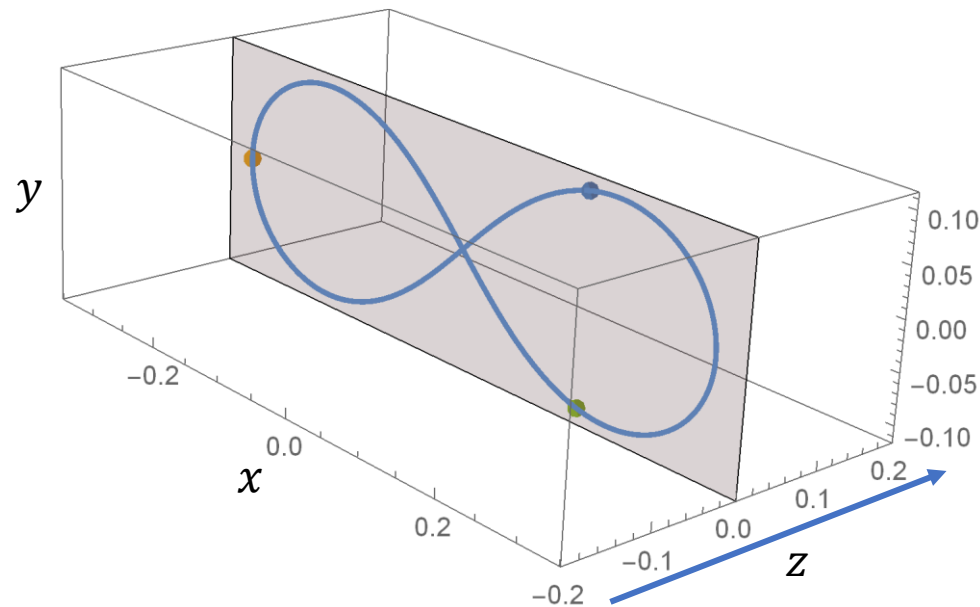
$$Sq(t) = -\tau q(-t), \tau: \text{exchange of body 1 and 2}$$



S : space and time inversion with exchange body, since space inversion is time inversion with exchange of two bodies.

Symmetry of figure eight choreography

Moreover, planar motion in x-y plane is invariant under z-inversion μ_z in three dimension.



3. Symmetry group and irreducible representations

Dihedral group D_{6h}

Therefore, $H(q)$ for figure eight choreography q commutes with the four operations

C, M, S and μ_z ,

and their combinations (products).

From the algebraic relations

$$C^3 = 1, M^2 = 1, S^2 = 1, \mu_z^2 = 1$$

$$CS = SC^{-1}, MS = SM, CM = MC,$$

$$C\mu_z = \mu_z C, S\mu_z = \mu_z S, M\mu_z = \mu_z M$$

group generated by $\{C, M, S, \mu_z\}$ is isomorphic to dihedral group D_{6h} with order $|D_{6h}| = 24$.

Irreducible representations of D_{6h}

D_{6h} has

8 one dimensional orthogonal irreducible representations and
 $(C, M, S, \mu_z) = (1, \pm 1, \pm 1, \pm 1)$

4 two dimensional orthogonal irreducible representations

$$(C, M, S, \mu_z) = \left(\begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}, \pm \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}, \pm \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \right)$$

Planar and non-planar bifurcation

If representation of μ_z is 1 or unit matrix

\Rightarrow eigenfunctions are symmetric in μ_z , that is, planar. Thus,

4 one dimensional orthogonal irreducible representations and

$$(C, M, S, \mu_z) = (1, \pm 1, \pm 1, \color{yellow}{+}1)$$

2 two dimensional orthogonal irreducible representations

$$(C, M, S, \mu_z) = \left(\begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}, \pm \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}, \color{yellow}{+} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \right)$$

are planar.

Planar and non-planar bifurcation

The rest half are non-planar:

4 one dimensional orthogonal irreducible representations and

$$(C, M, S, \mu_z) = (1, \pm 1, \pm 1, -1)$$

2 two dimensional orthogonal irreducible representations

$$(C, M, S, \mu_z) = \left(\begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}, \pm \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}, - \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \right)$$

Symmetry of eigenfunction and bifurcation solution

Moreover, symmetry of i 'th eigenfunction ϕ_i is also given by the representation:

It is a subgroup $G_e(i)$ consisting of the elements with i 'th diagonal element of the representation is one

$$G_e(i) = \{ g \mid (g)_{ii} = 1 \}$$

Then, the symmetry of the bifurcation solution $q^{(b)}$ is given by $G_e(i)$ (Fujiwara et al 2020).

4. Bifurcation of figure eight to the planar motion

4 one-dimensional planar representations

$$(1, 1, -1, 1)$$

$$(1, -1, 1, 1)$$

$$(1, -1, -1, 1)$$

$$(1, 1, 1, 1)$$

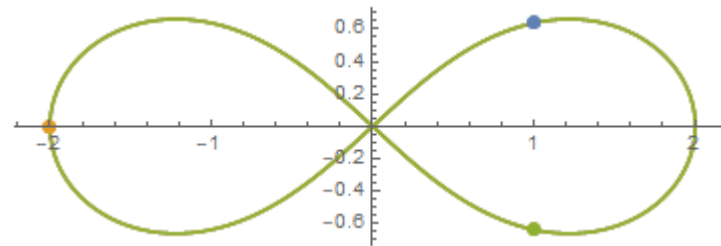
2 two-dimensional planar representations

$$\left(\begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}, + \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \right)$$

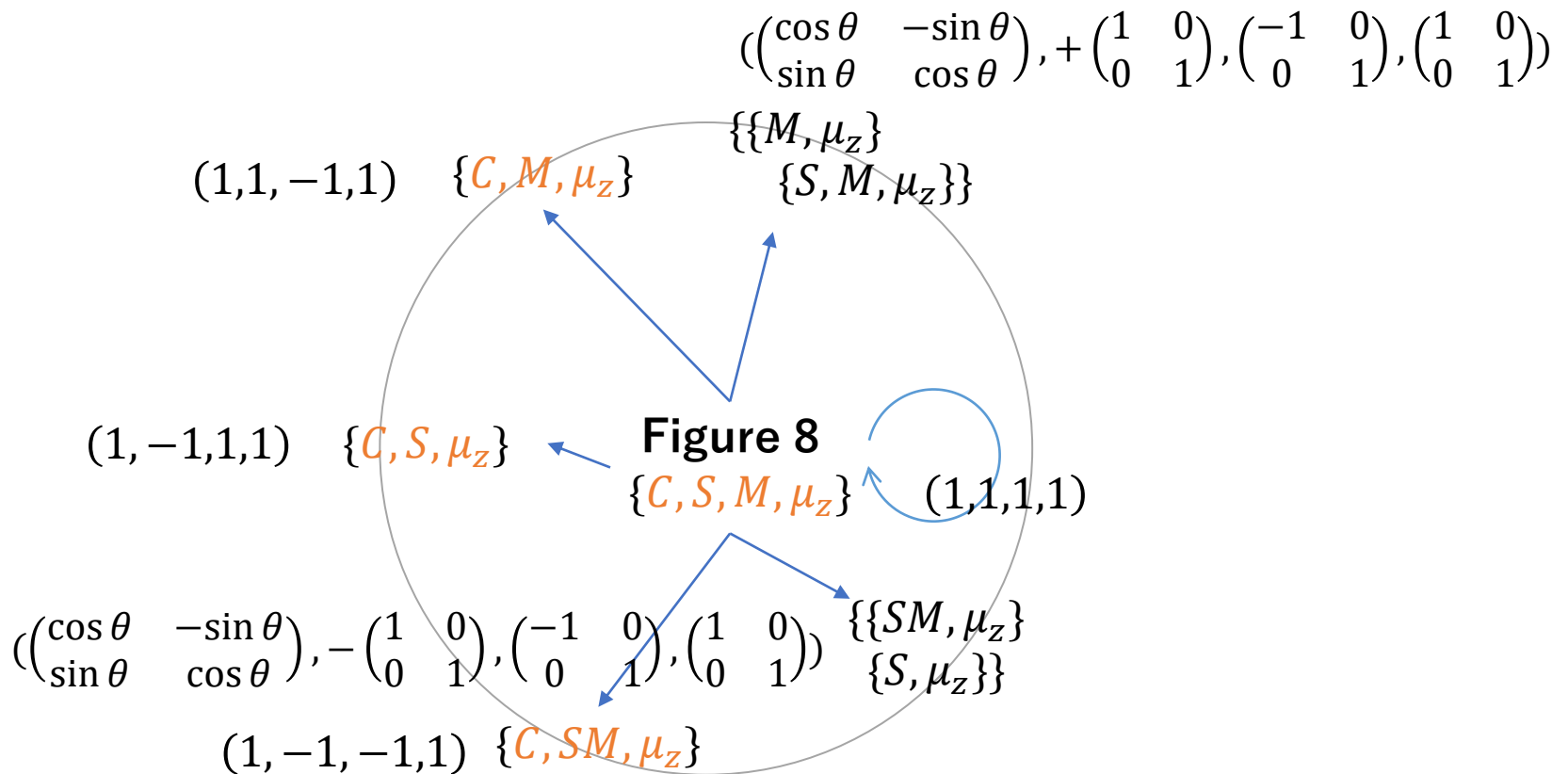
$$\left(\begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}, - \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \right)$$

Figure 8

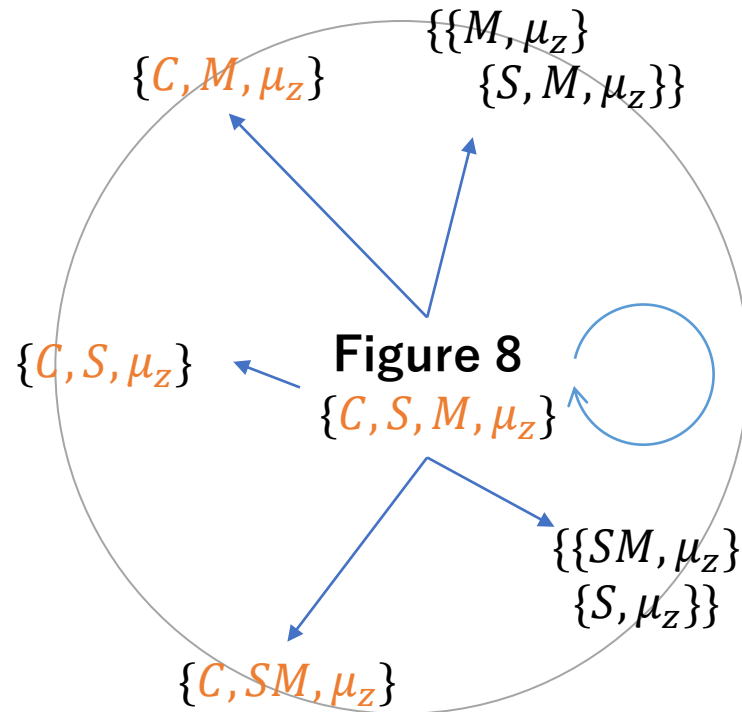
$\{C, S, M, \mu_z\}$



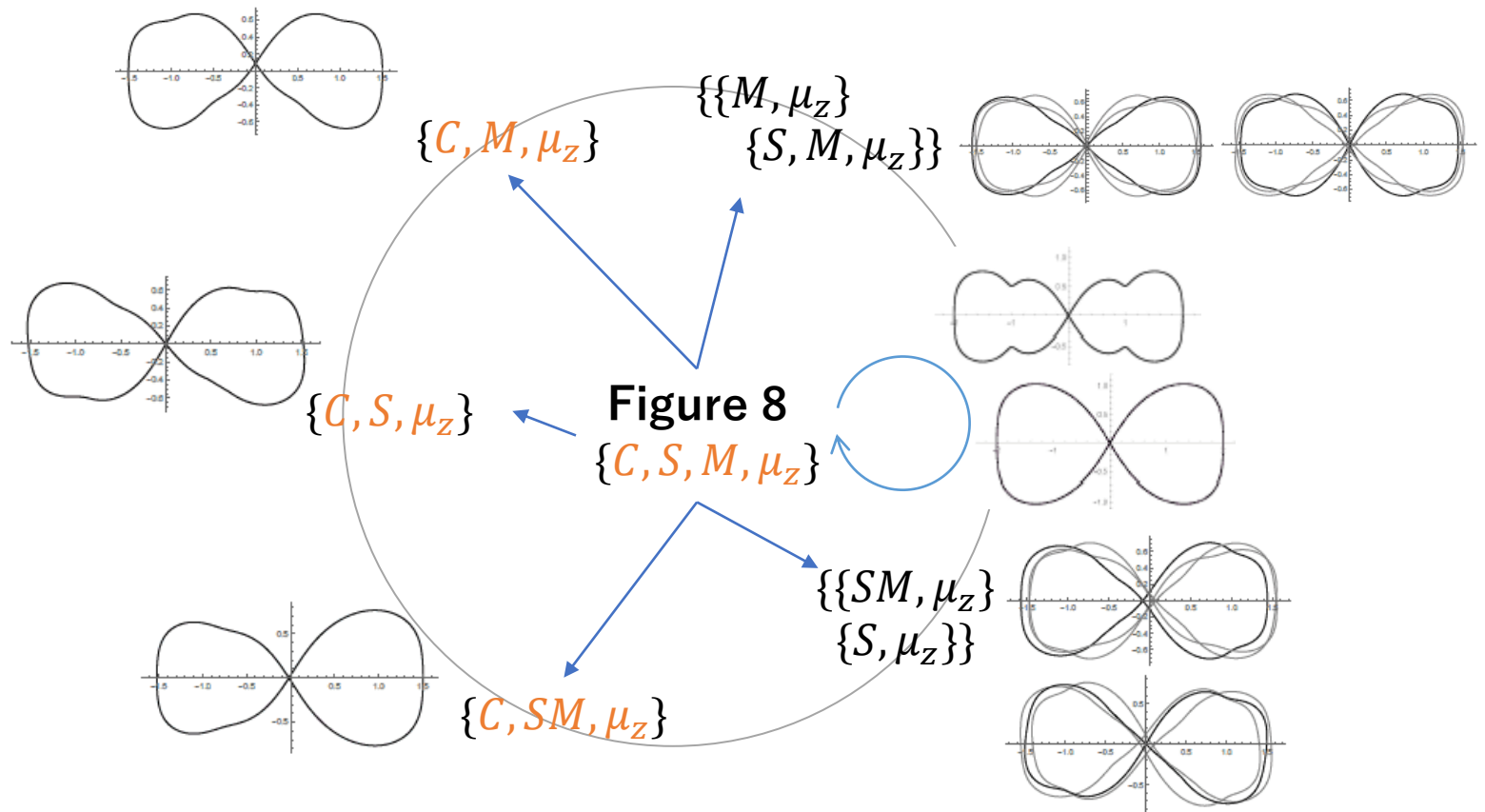
Thus, the figure 8 choreography bifurcates to planar solutions represented by four one dimensional and two two-dimensional representations.



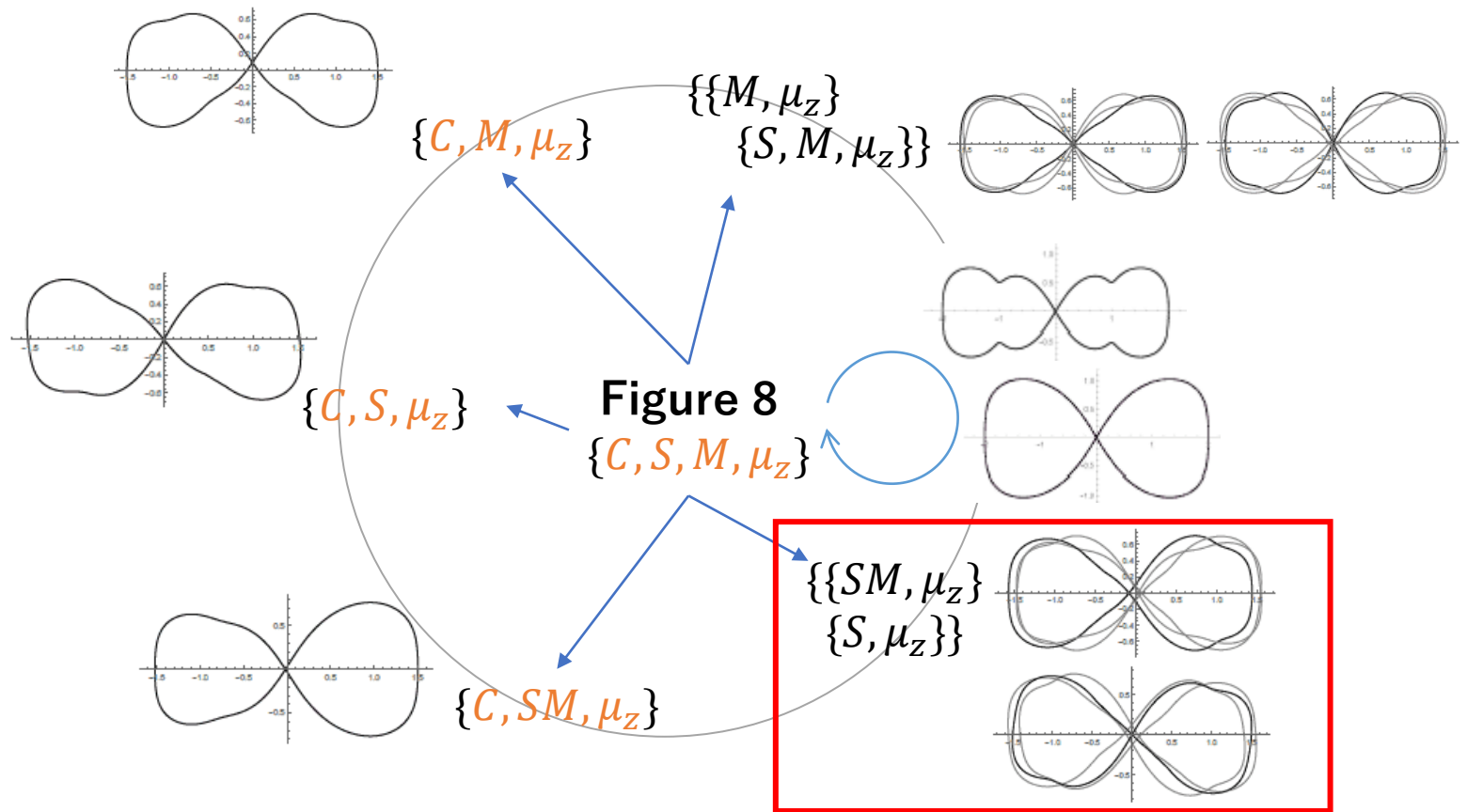
In this figure, symmetries are shown in curly parenthesis,



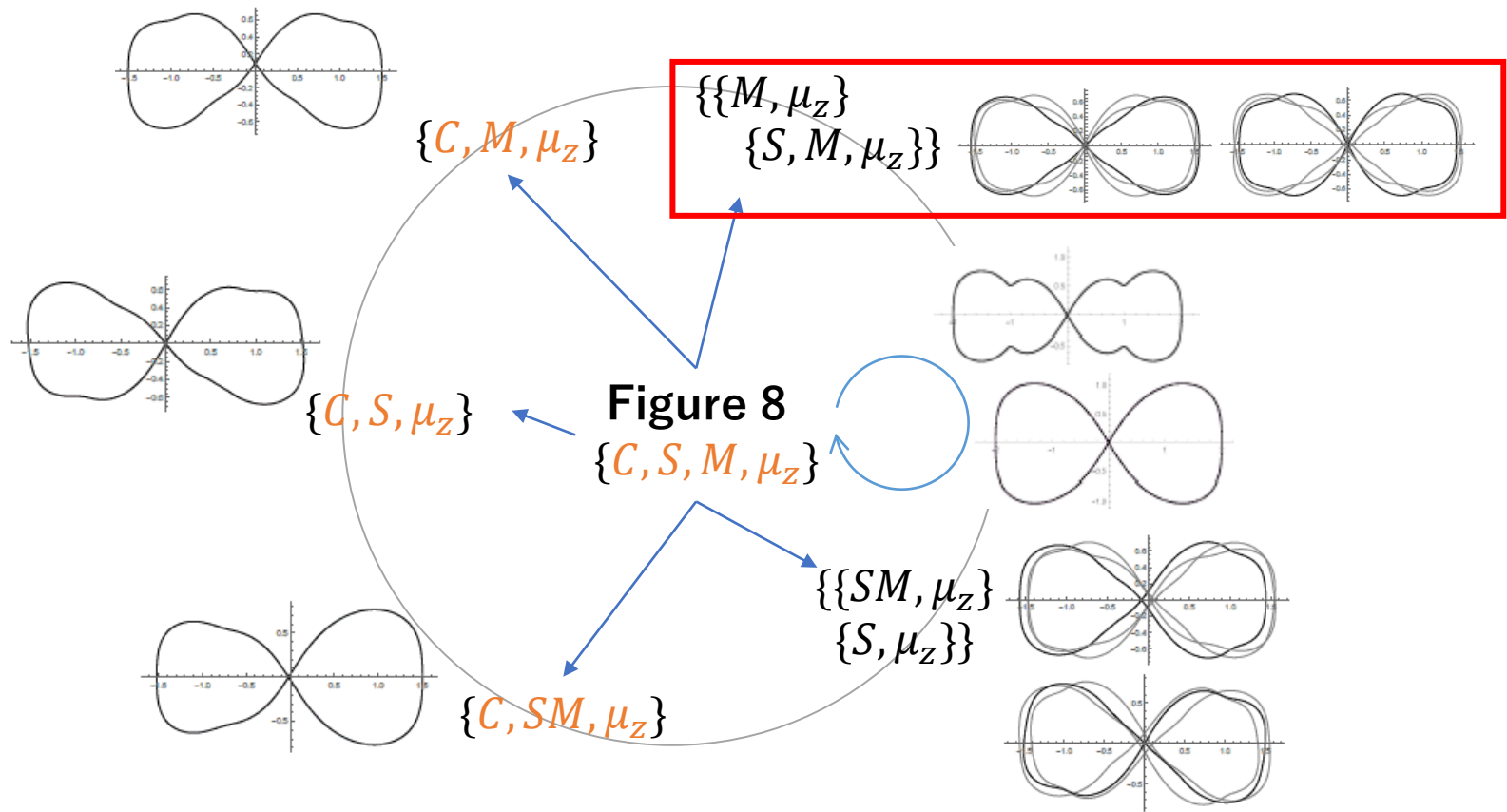
allows show bifurcations, and choreographies are emphasized by orange.



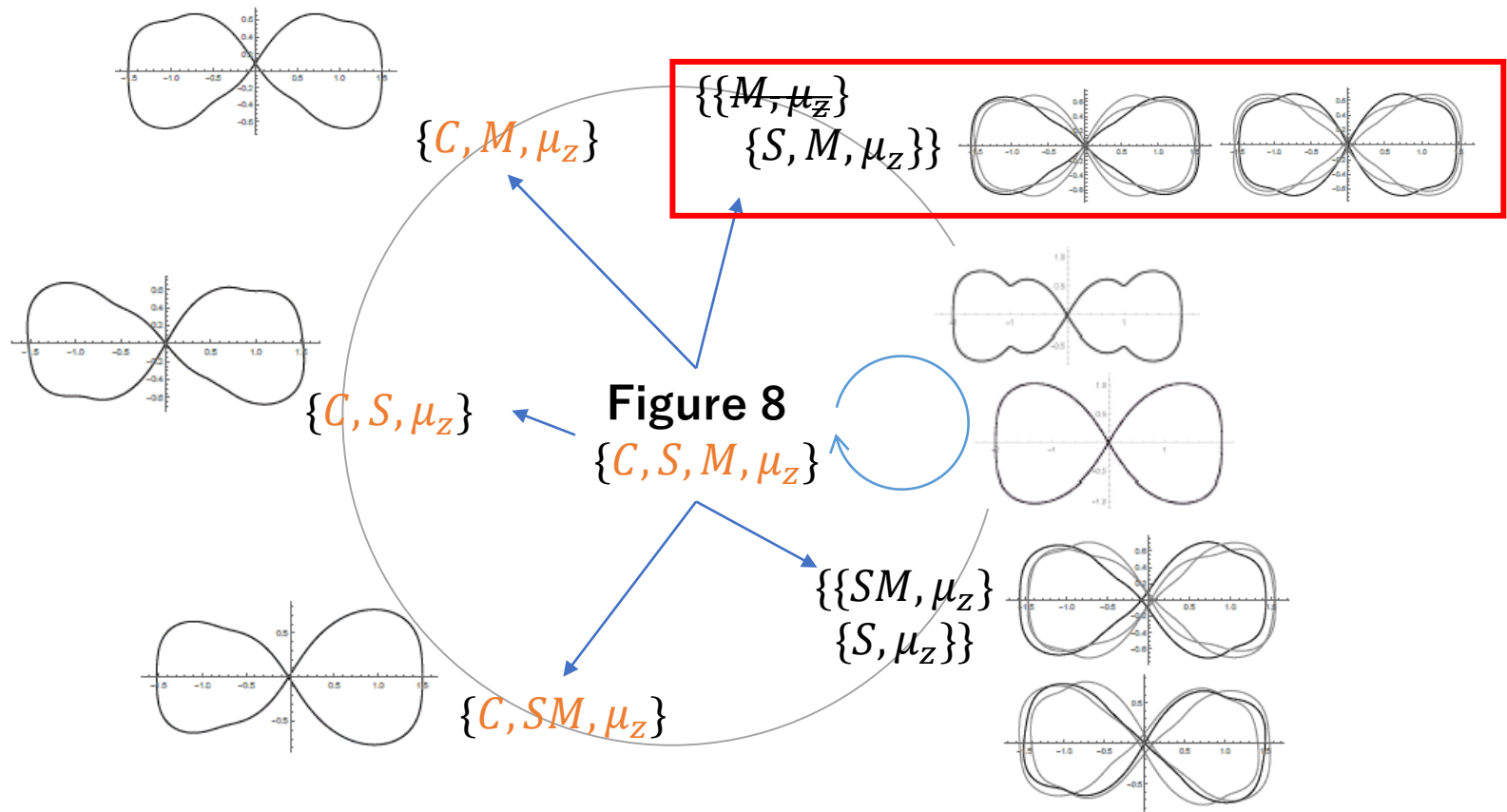
In 2019, we found all these bifurcations numerically in LJ system.



Note that two dimensional representations yield two bifurcations having two symmetries in curly parenthesis like $\{SM, \mu_z\}$ and $\{S, \mu_z\}$.

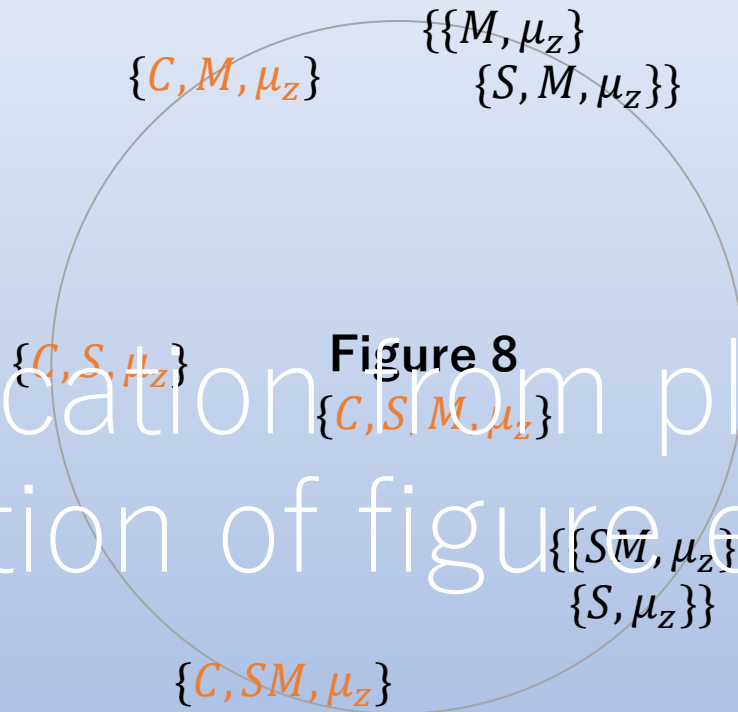


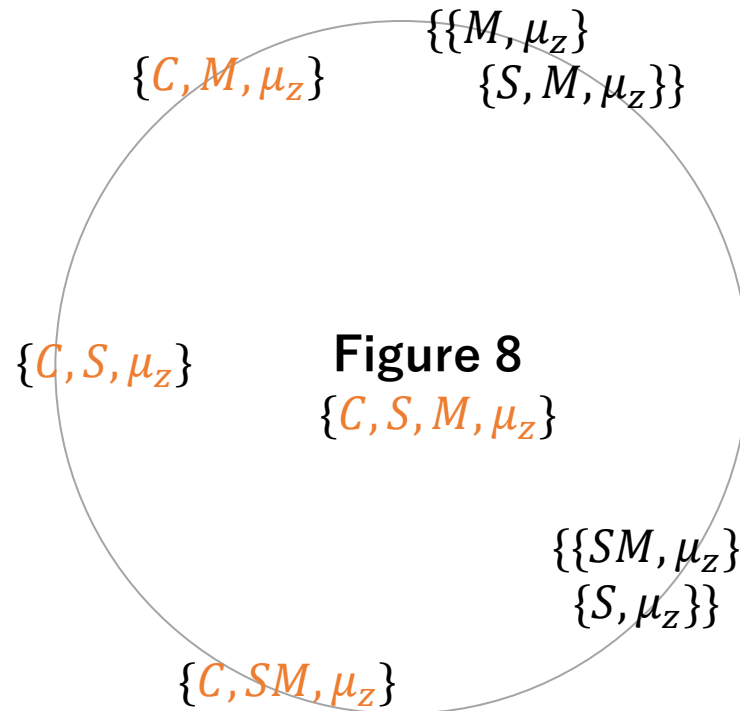
However, if one symmetry is subset of the other as $\{M, \mu_z\}$ and $\{S, M, \mu_z\}$, both bifurcations have the same larger symmetry, as $\{S, M, \mu_z\}$.



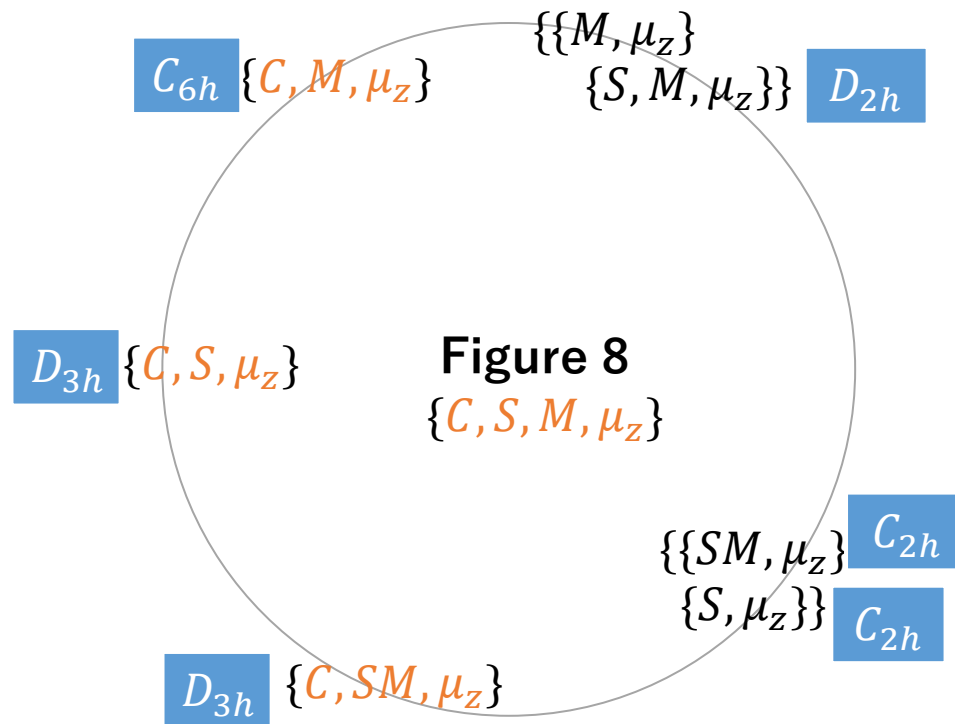
Thus, we strike out smaller symmetry.

5. Bifurcation from planar bifurcation of figure eight

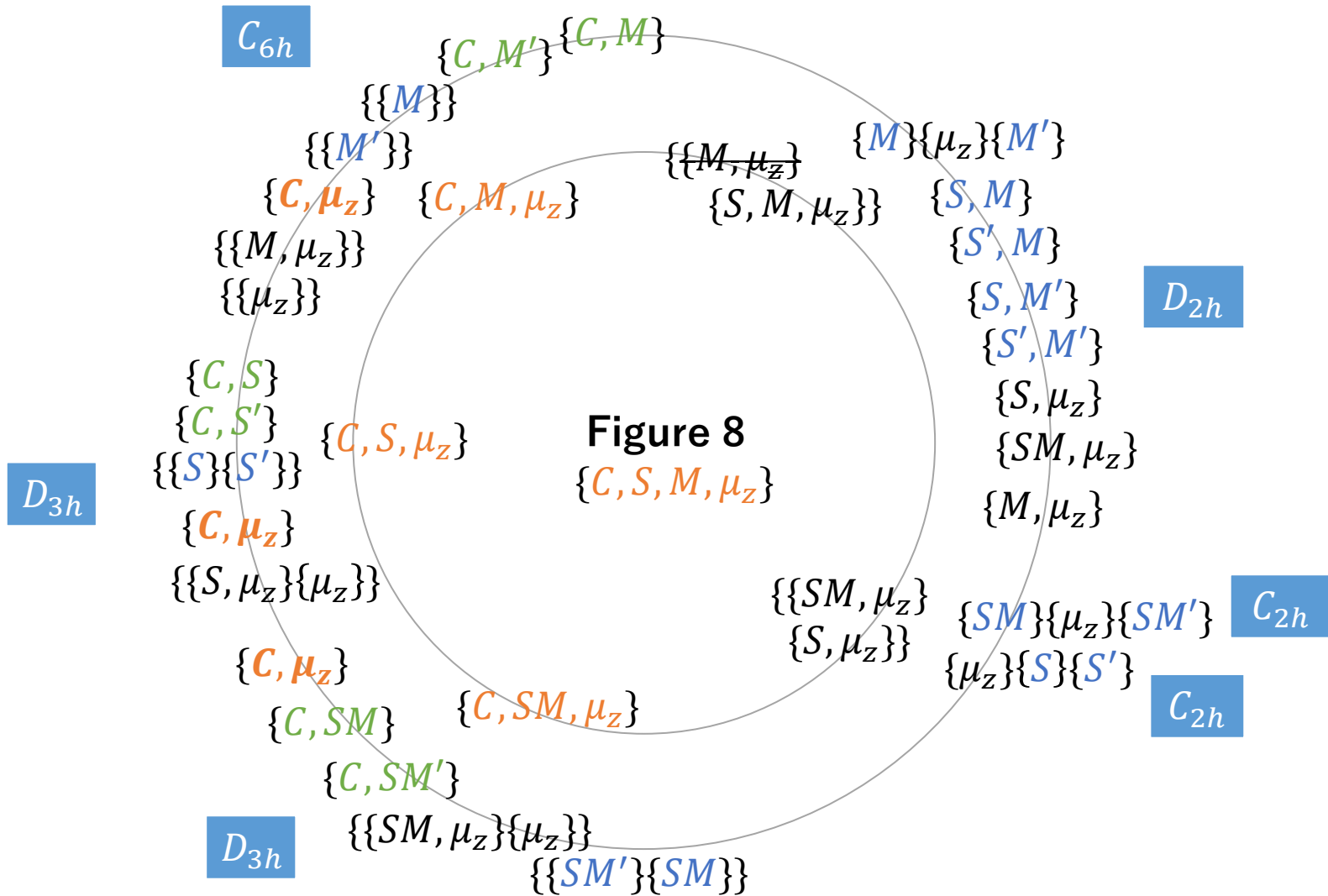




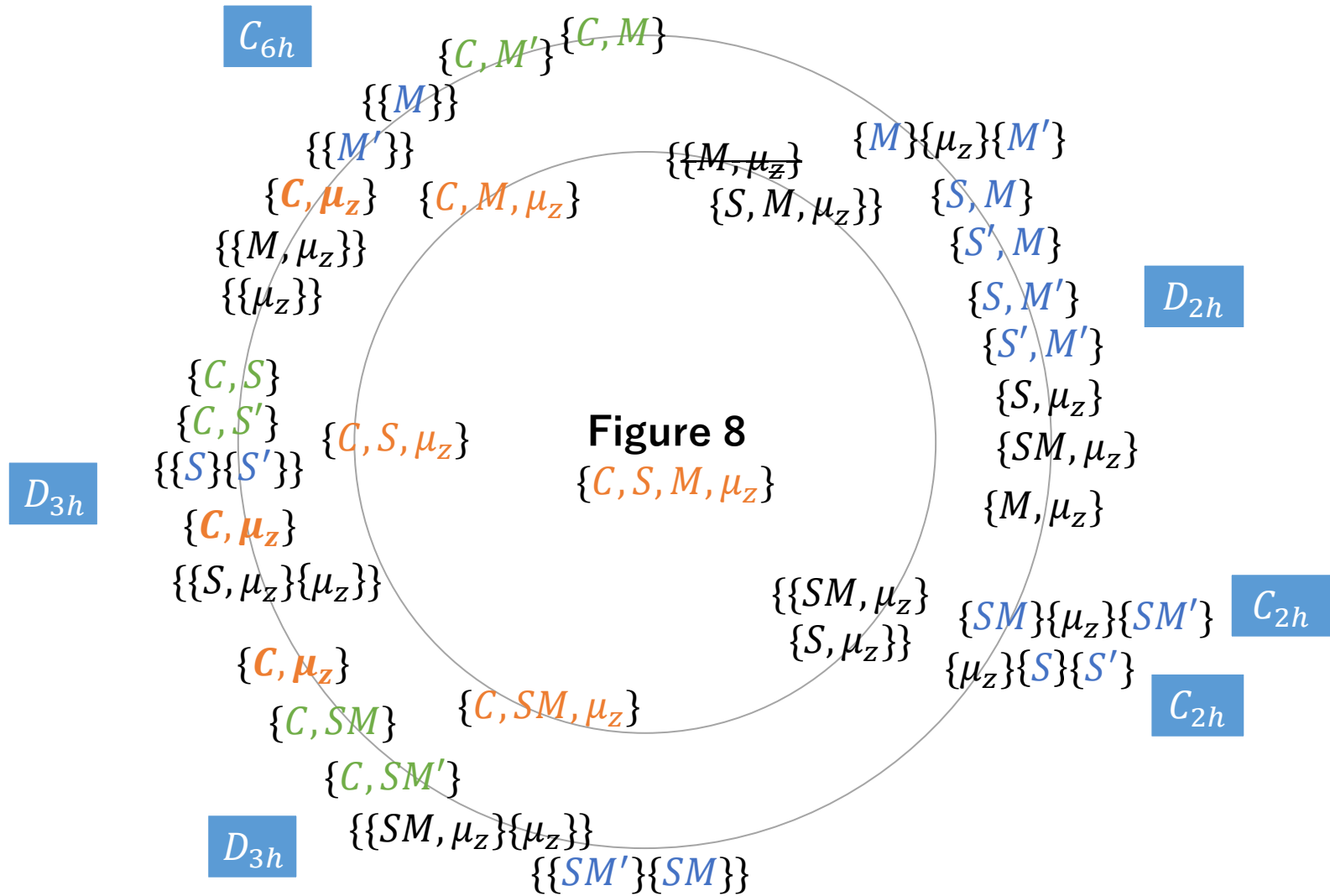
Next, we consider the bifurcation from these bifurcations. We first consider the symmetry group for each solutions,



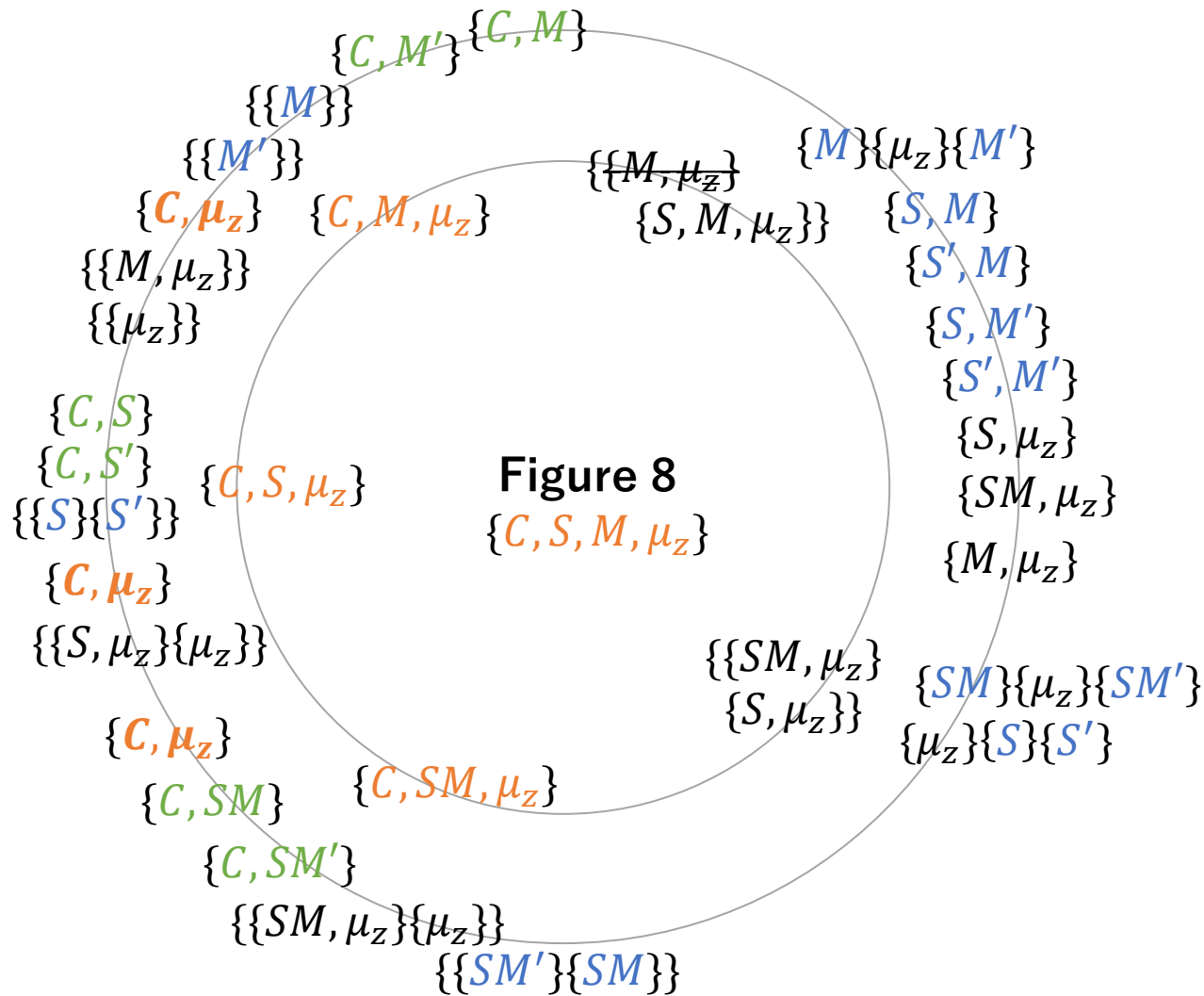
From their orthogonal irreducible representations,



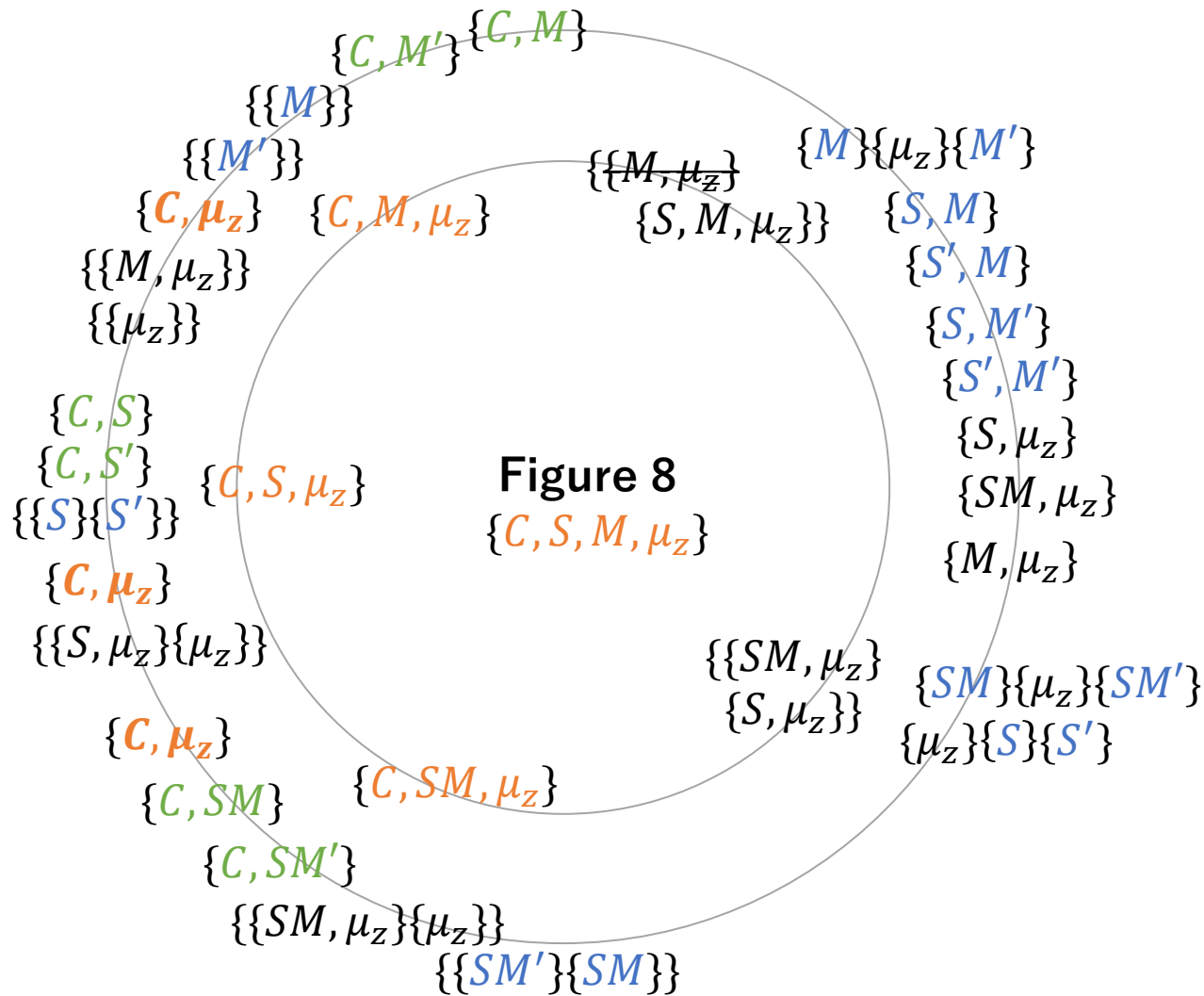
we obtain bifurcation from bifurcations, shown on the second ring.



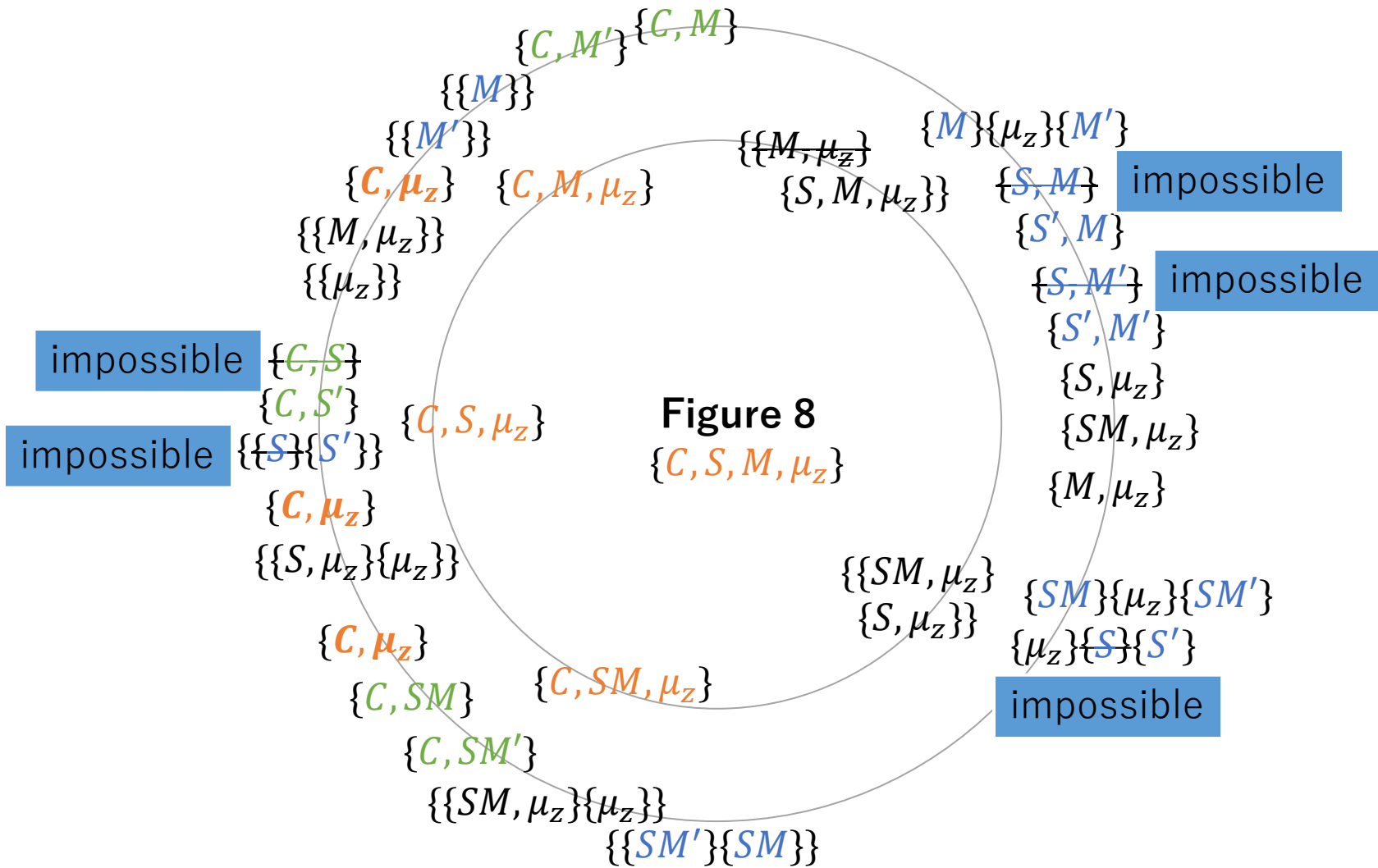
Here, a non-planar choreographies are colored by green, non-planar by blue, and prime is μ_z .



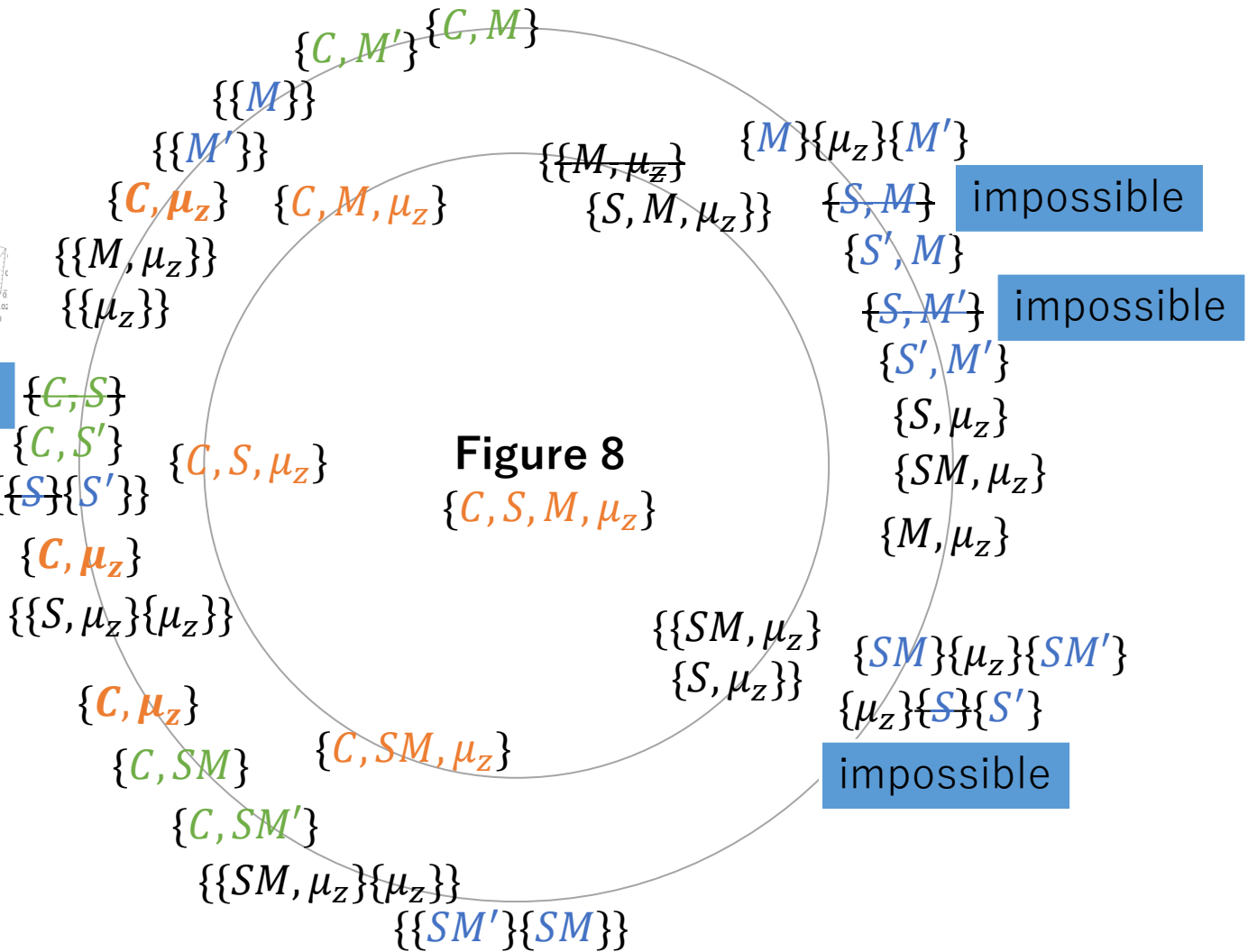
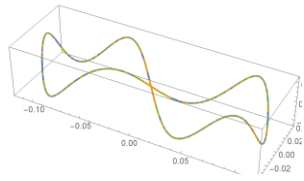
However, several of them can not be solutions because of angular momentum restrictions.



If angular momentum is zero, three body motion is planar. On the other hand, operation S changes sign of the angular momentum.



Thus, bifurcation with S but without μ_z is impossible. (Fujiwara 2020)



For example, this choreography is impossible.

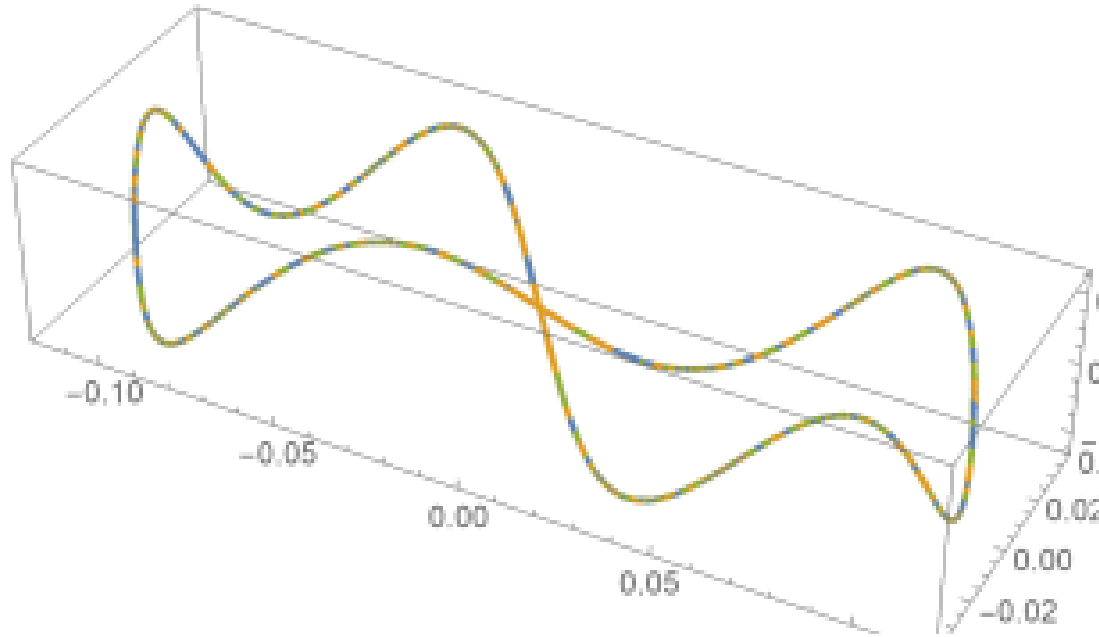
$\{C, M'\}$ $\{C, M\}$

impossible

possible

possible

impossib

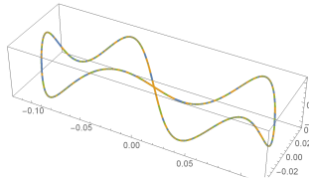


$\{C, SM'\}$

$\{SM, \mu_z\}$ $\{\mu_z\}$

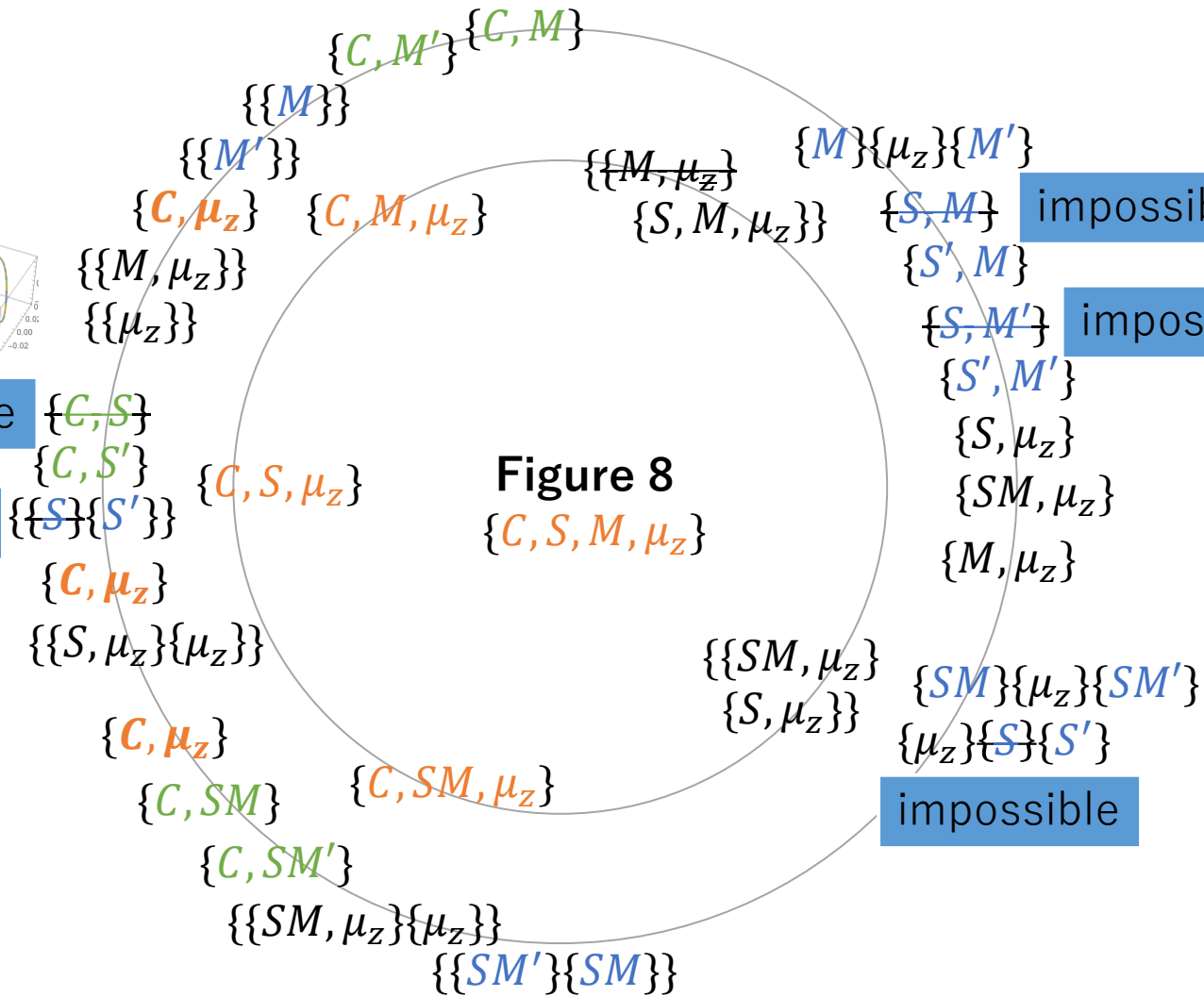
$\{SM'\}$ $\{SM\}$

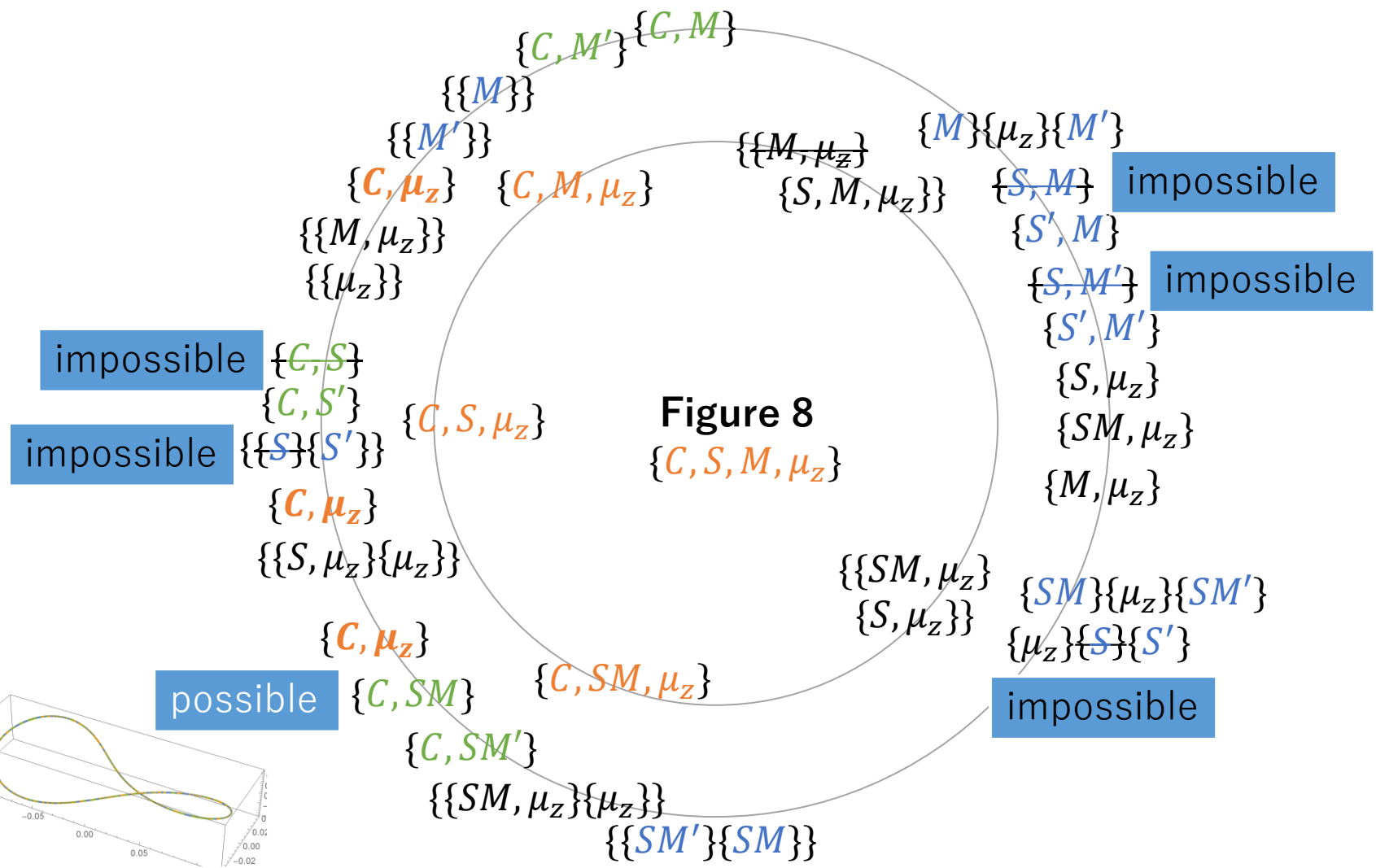
$\{M'\}$



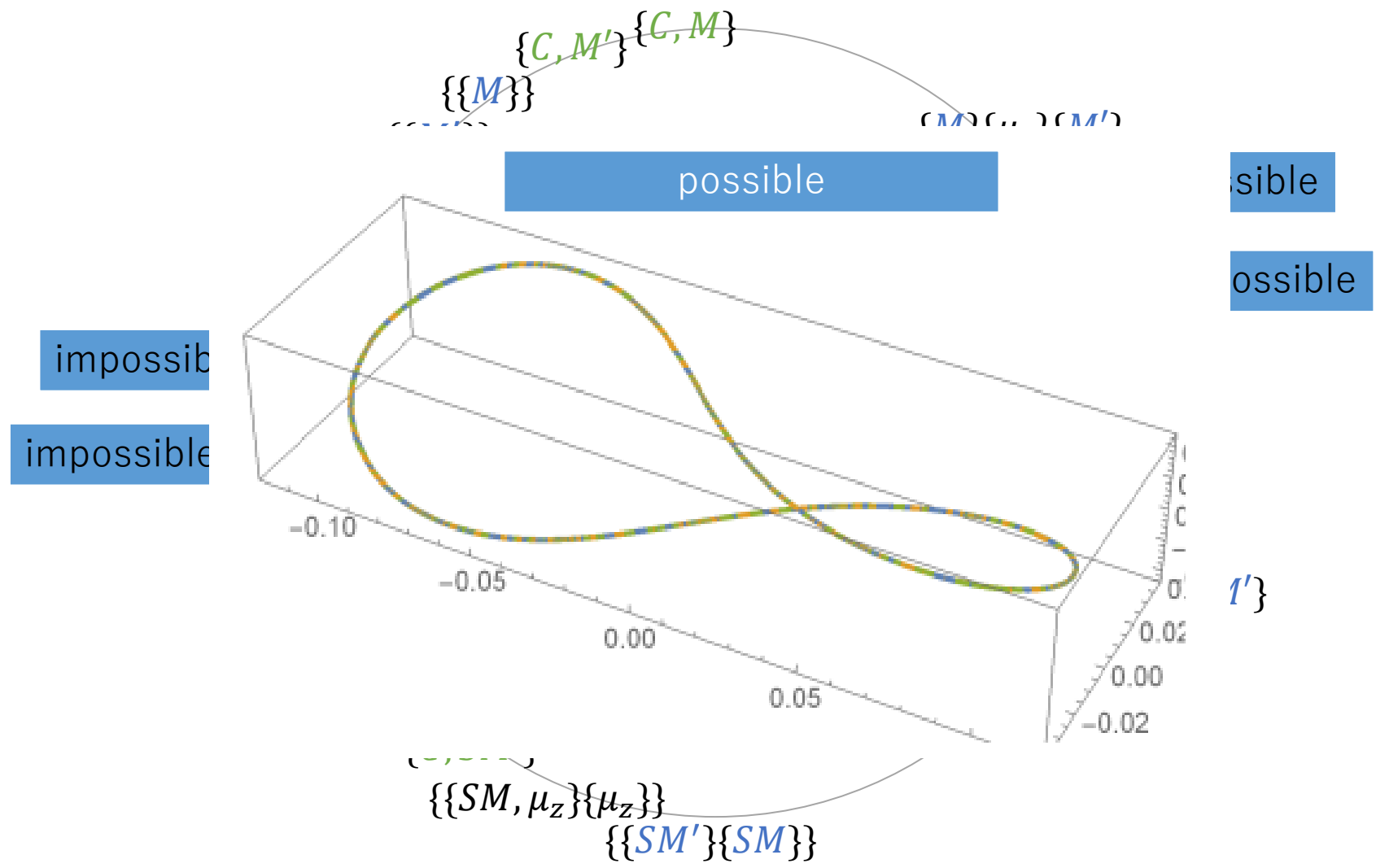
impossible

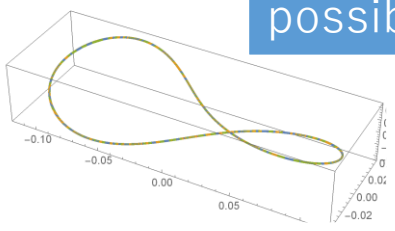
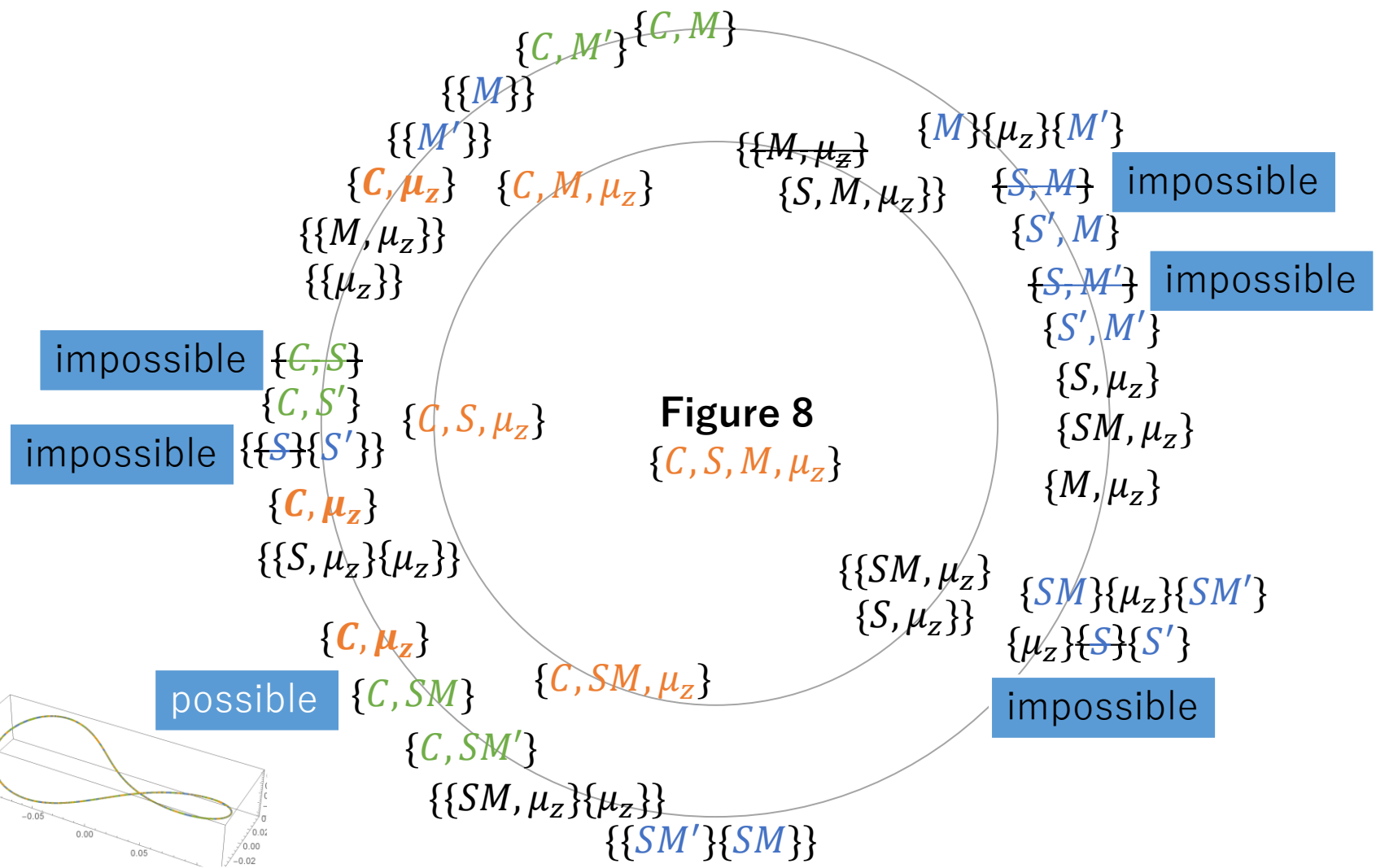
impossible

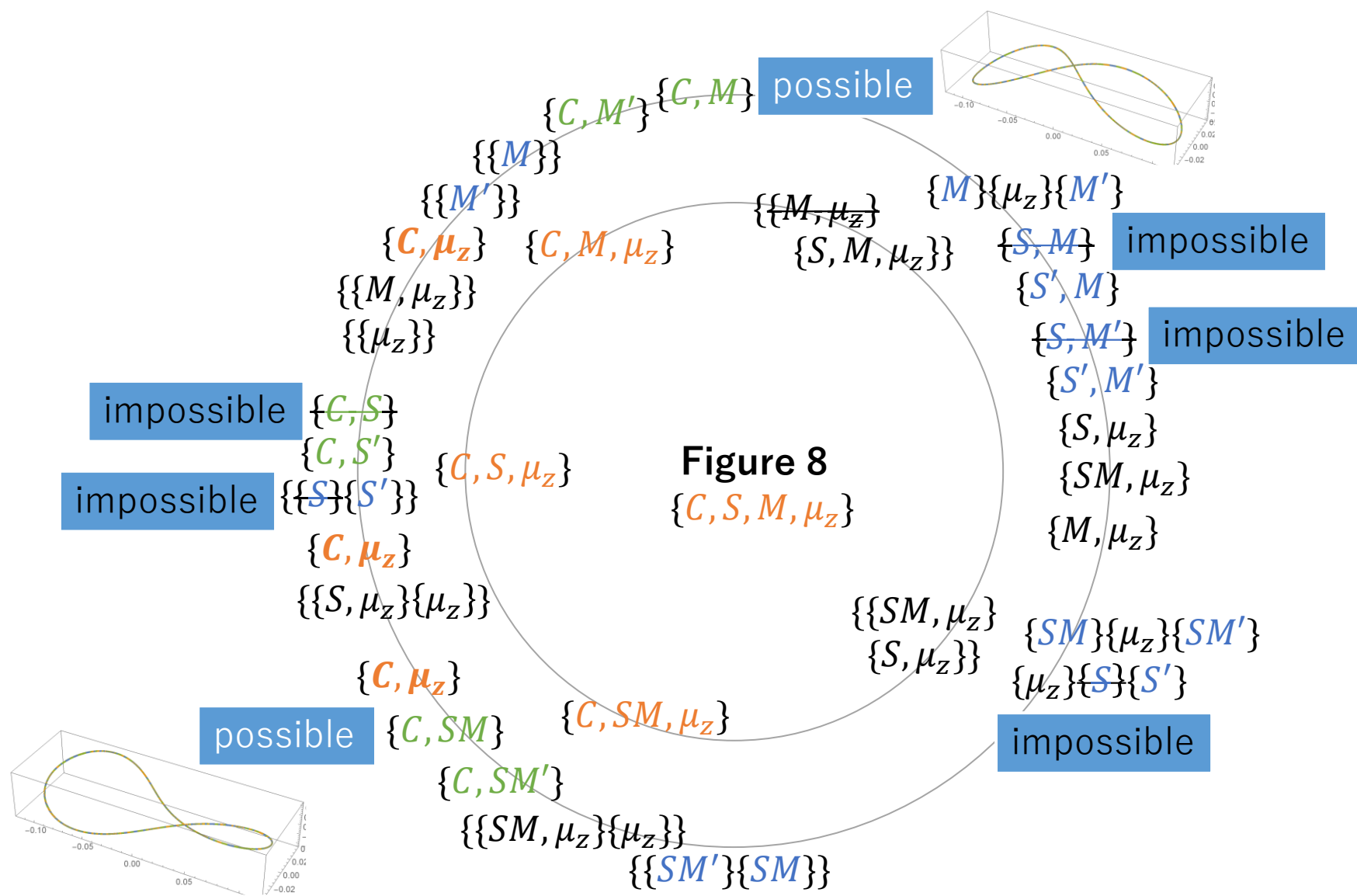




While, for example, this is a possible non-planar choreography.







Also this is a possible non-planar choreography.

$\{\{M\}\}$ $\{C, M'\}$ $\{C, M\}$

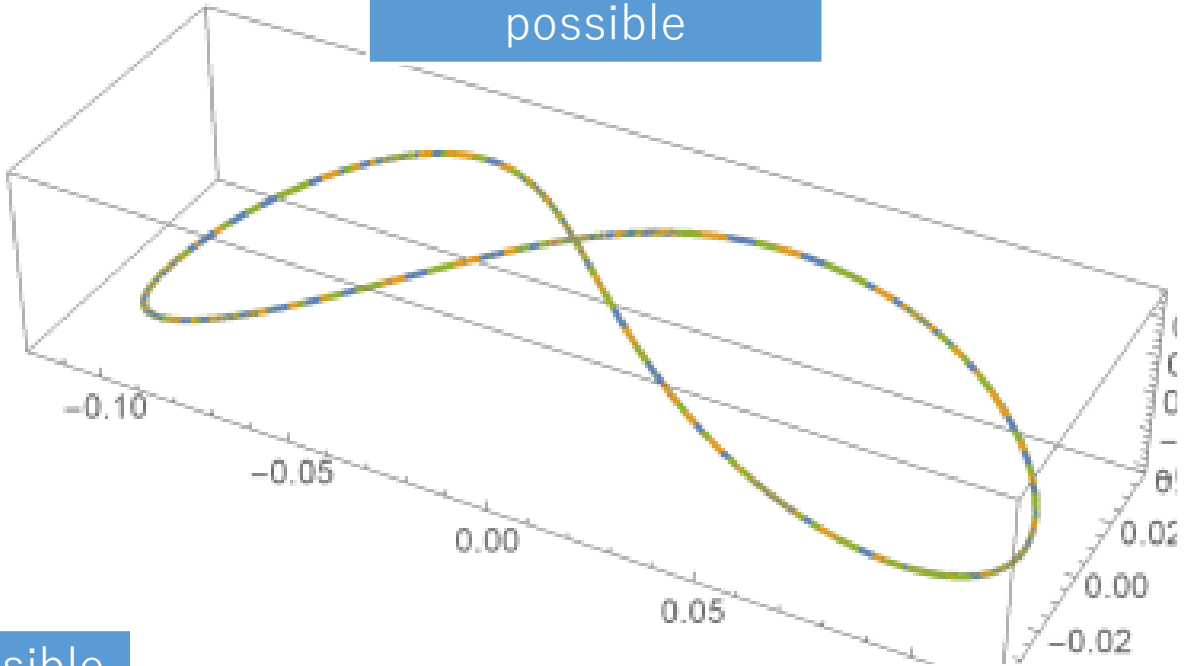
possible

possible

possible

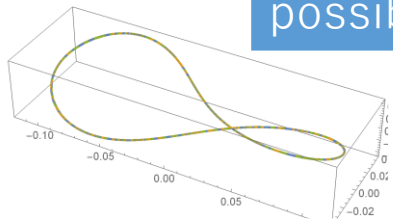
impossible

impossible

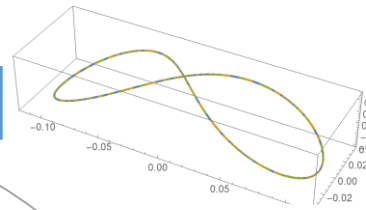


$\{I'\}$

possible



$\{C, SM'\}$
 $\{\{SM, \mu_z\}\{\mu_z\}\}$
 $\{\{SM'\}\{SM\}\}$



possible

$\{C, M'\}$ $\{C, M\}$

$\{M\}$

$\{M'\}$

$\{C, \mu_z\}$

$\{C, M, \mu_z\}$

$\{M, \mu_z\}$

$\{M\}$ $\{\mu_z\}$ $\{M'\}$

$\{S, M, \mu_z\}$

$\{S, M\}$

impossible

$\{M, \mu_z\}$

$\{S', M\}$

$\{\mu_z\}$

$\{S', M'\}$

impossible

impossible

$\{C, S\}$

Figure 8

impossible

$\{S\}$ $\{S'\}$

$\{C, S, \mu_z\}$

$\{C, S, M, \mu_z\}$

$\{S, \mu_z\}$

$\{C, \mu_z\}$

$\{SM, \mu_z\}$

$\{S, \mu_z\}$ $\{\mu_z\}$

$\{SM, \mu_z\}$

$\{SM\}$ $\{\mu_z\}$ $\{SM'\}$

$\{C, \mu_z\}$

$\{S, \mu_z\}$

$\{\mu_z\}$ $\{S\}$ $\{S'\}$

possible

$\{C, SM\}$

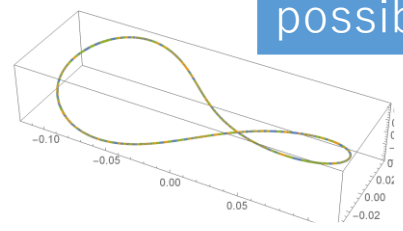
$\{C, SM, \mu_z\}$

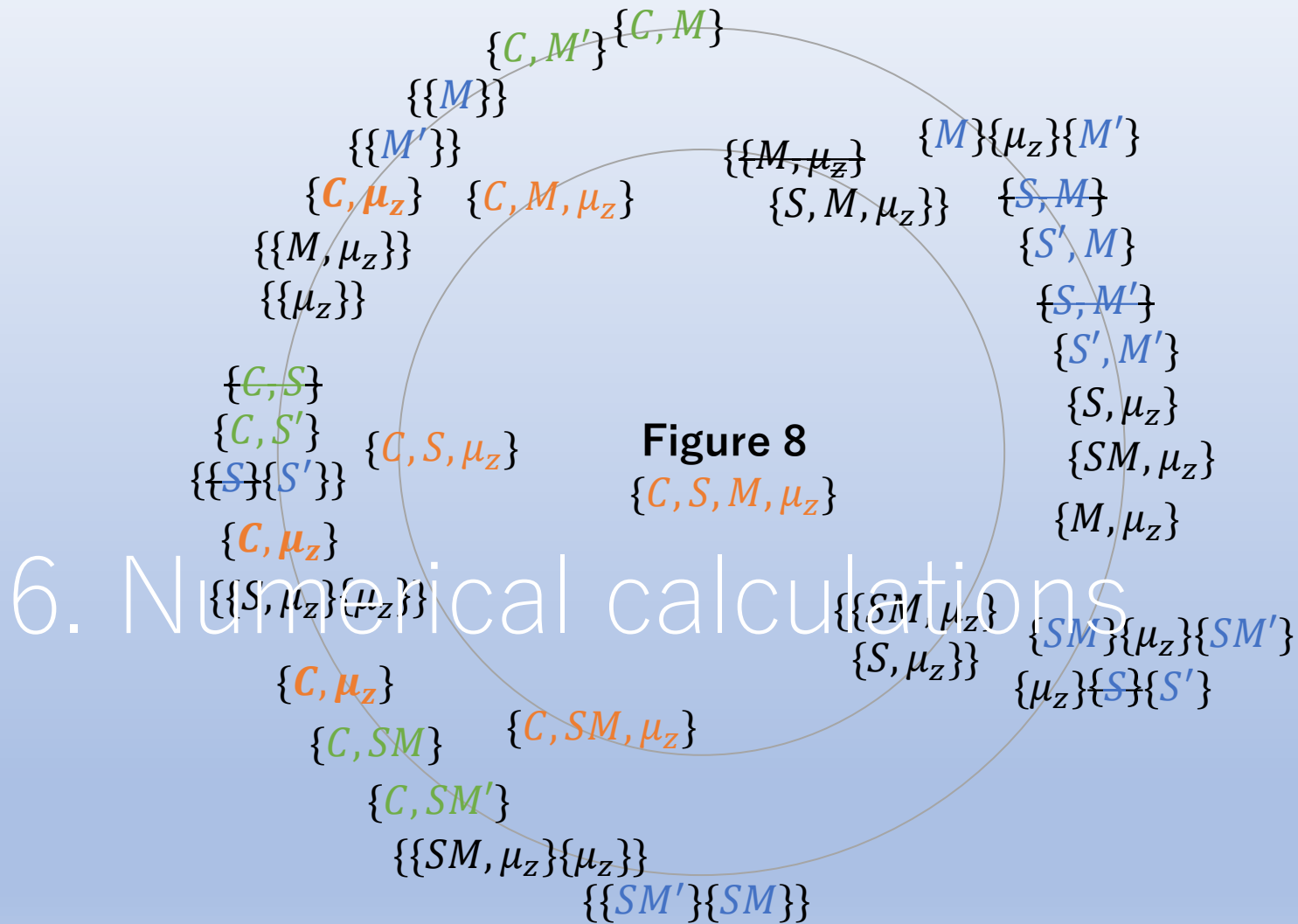
impossible

$\{C, SM'\}$

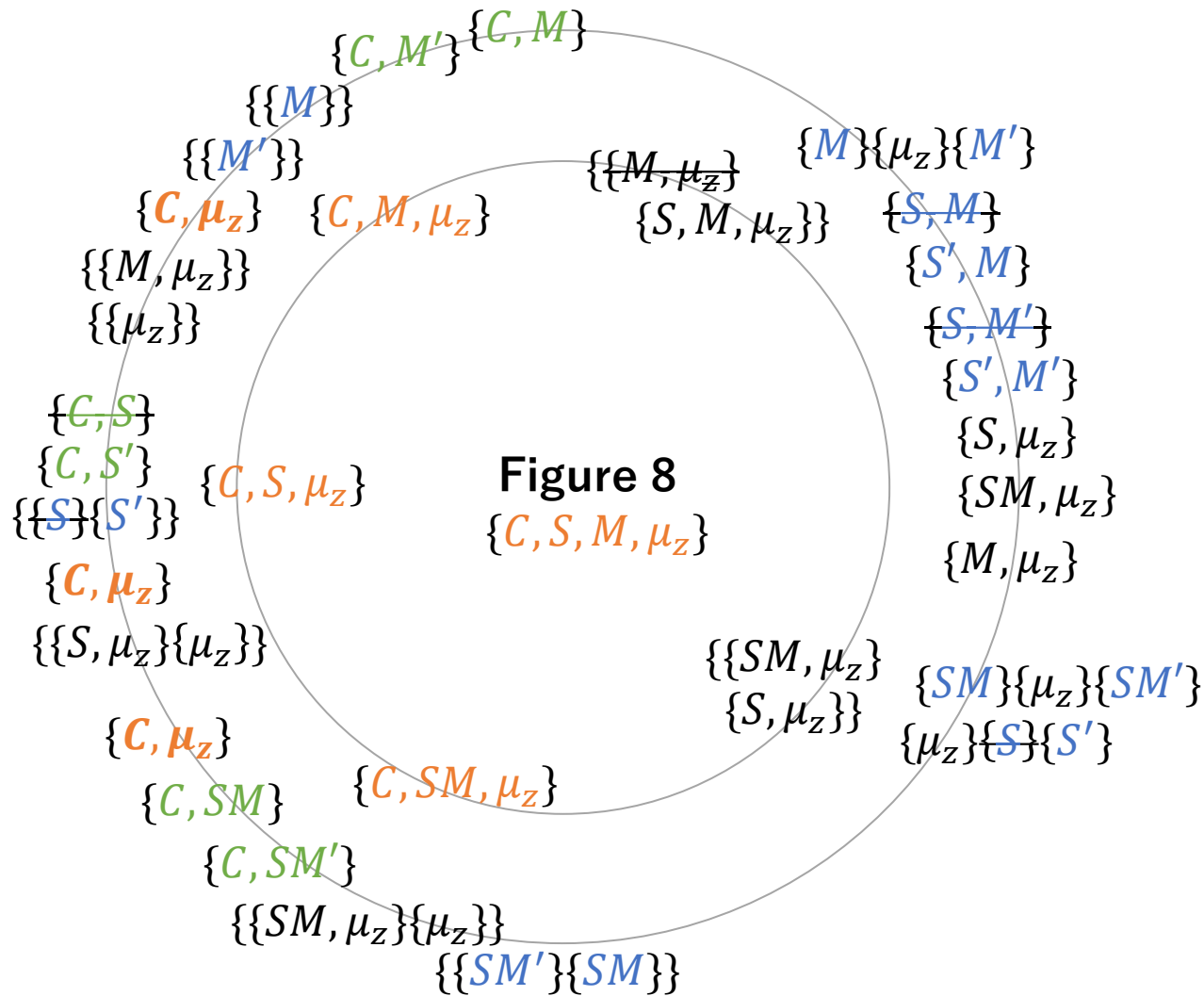
$\{SM, \mu_z\}$ $\{\mu_z\}$

$\{SM'\}$ $\{SM\}$

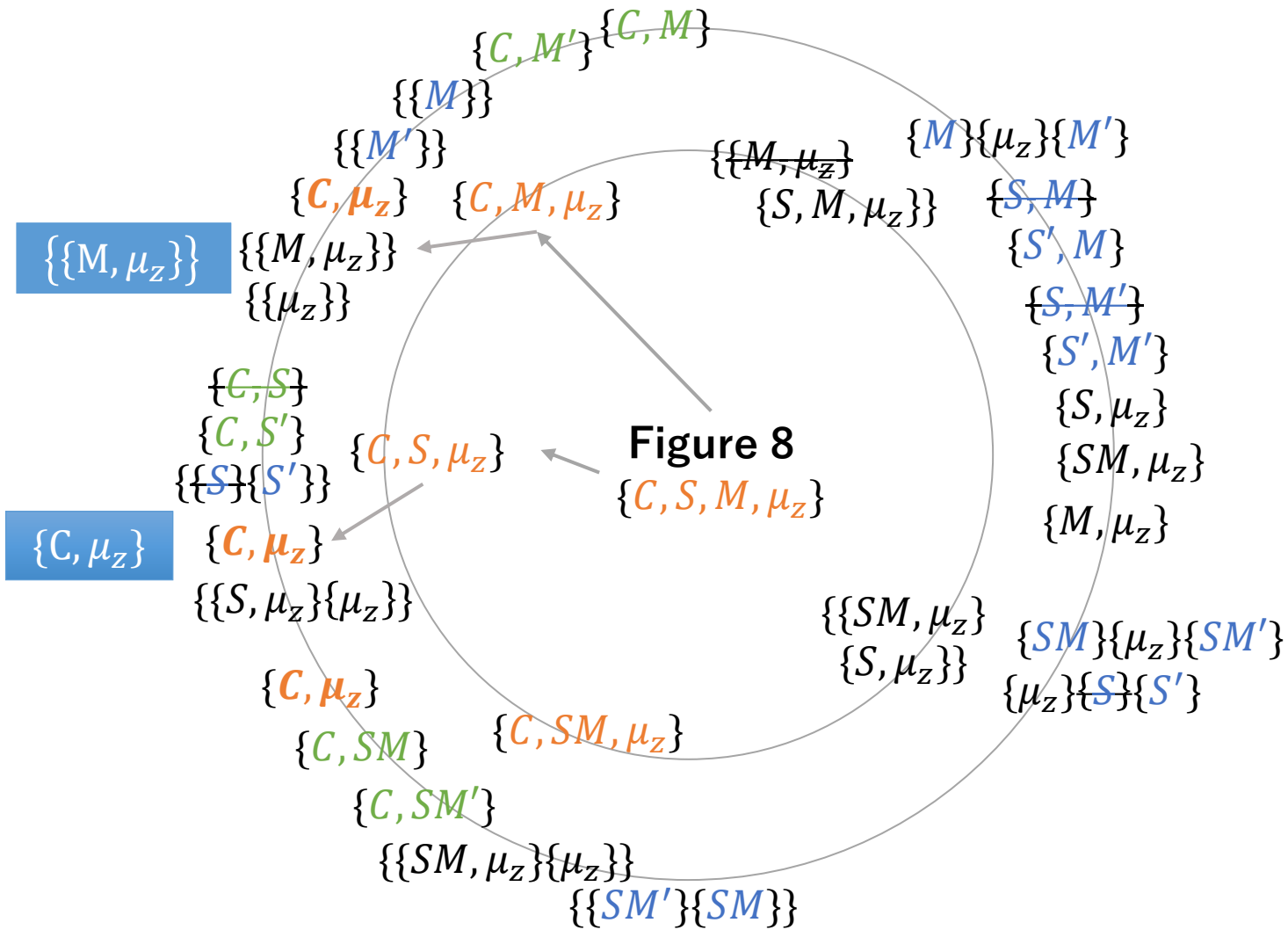




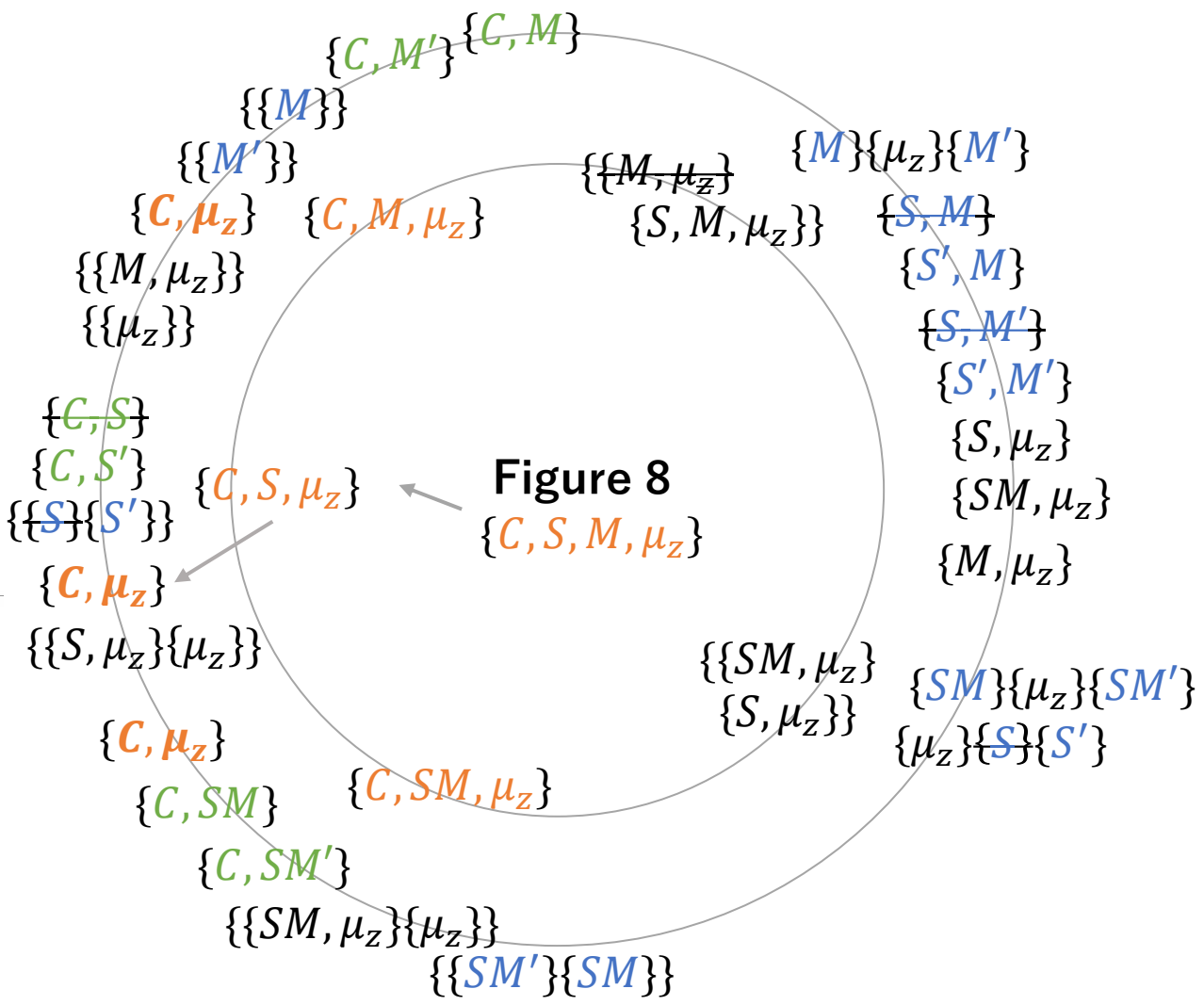
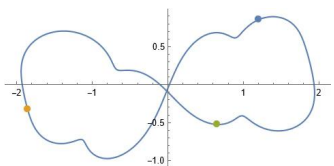
6. Numerical calculations



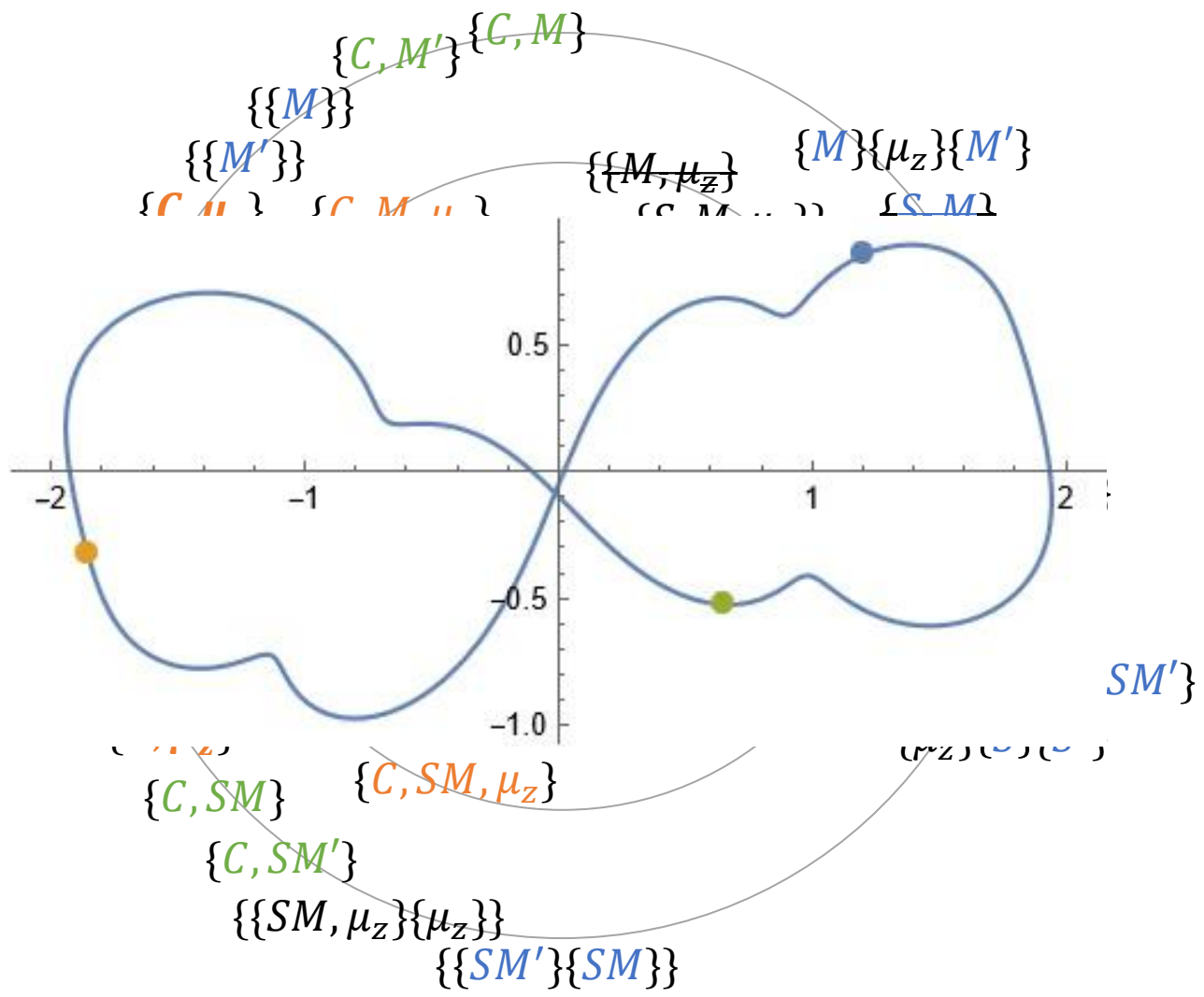
From 2020, we are calculating numerically bifurcations on the second ring.

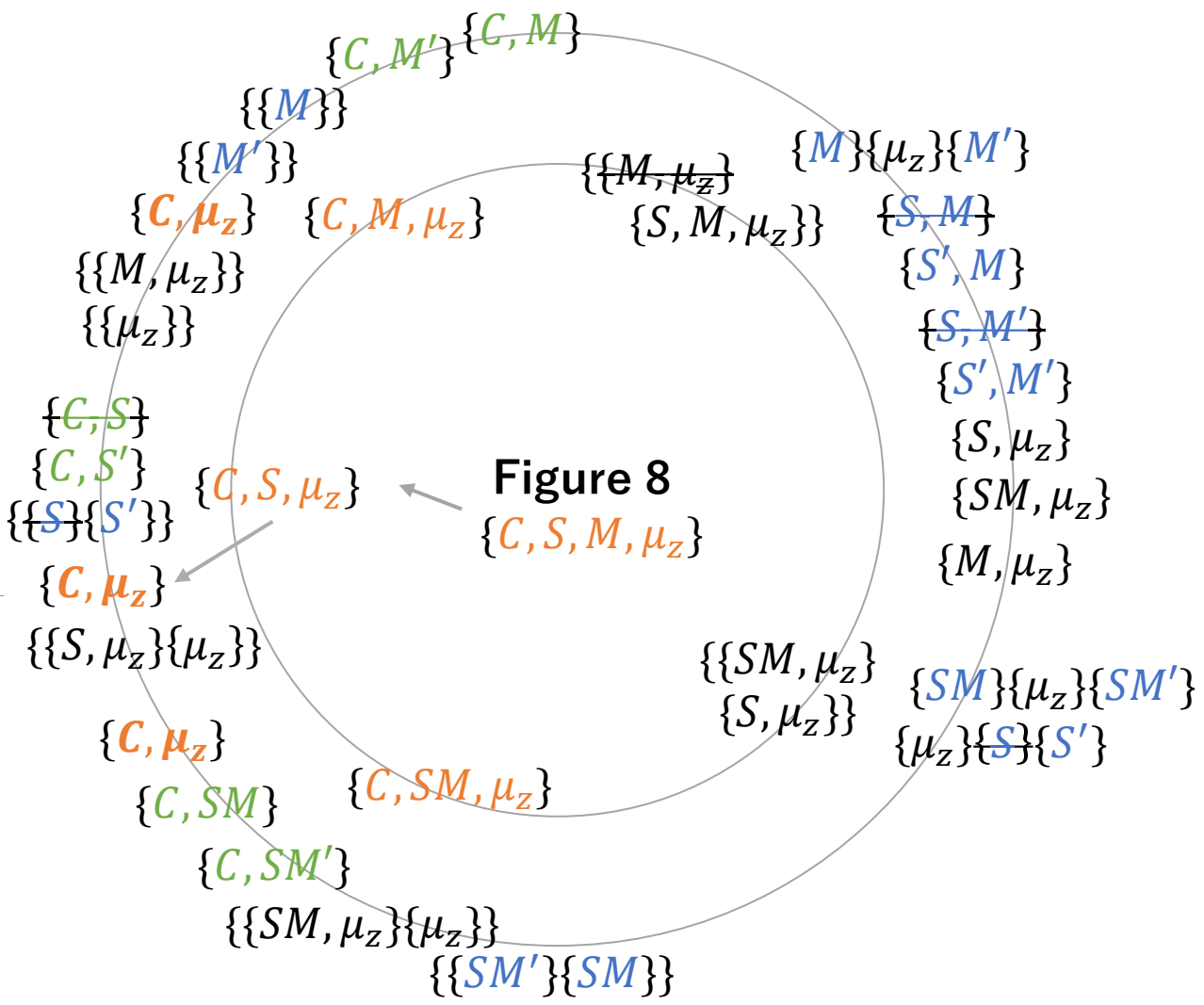
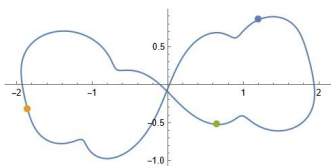


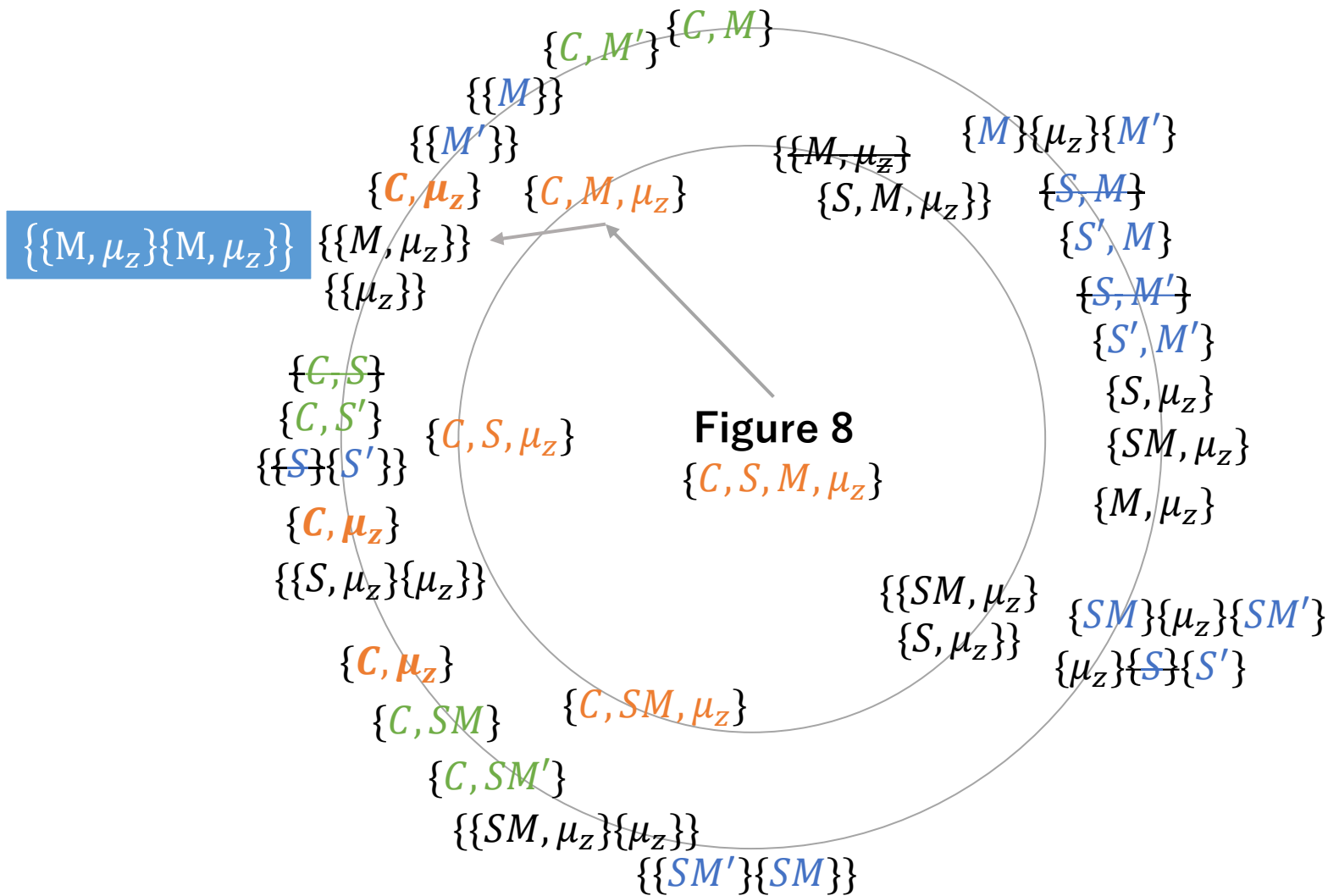
In LJ system, two bifurcations are interesting.



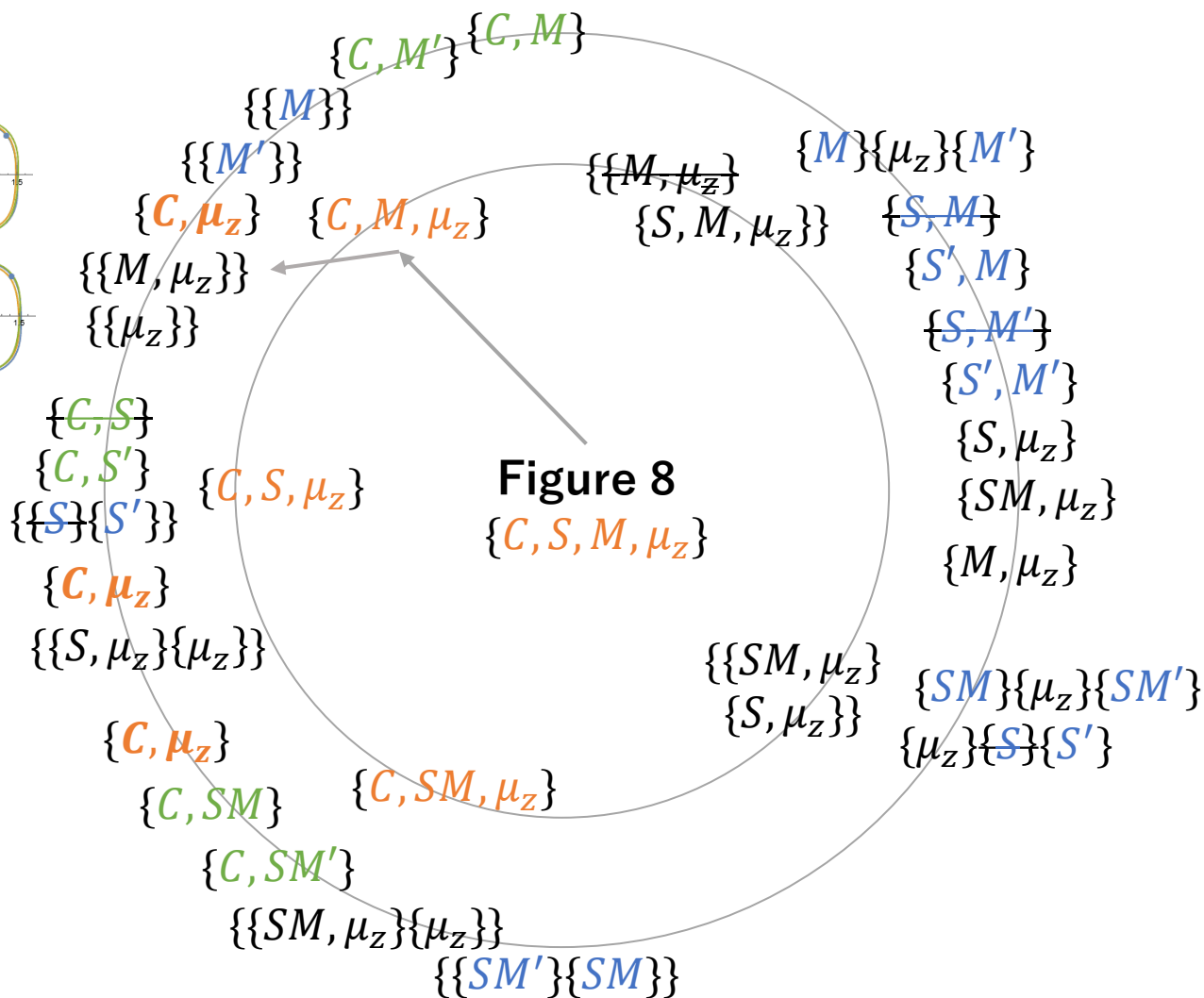
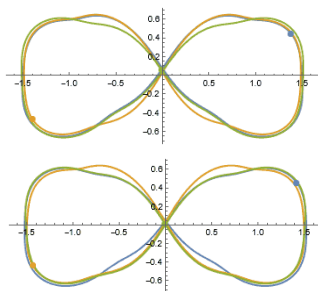
This is choreographic C without any spatial symmetry in the plane.







This, abbreviation of repetition, has two eigenfunctions with the same symmetry $\{M, \mu_z\}$ in two-dimensional representation, and



yields two solutions with the same symmetry.

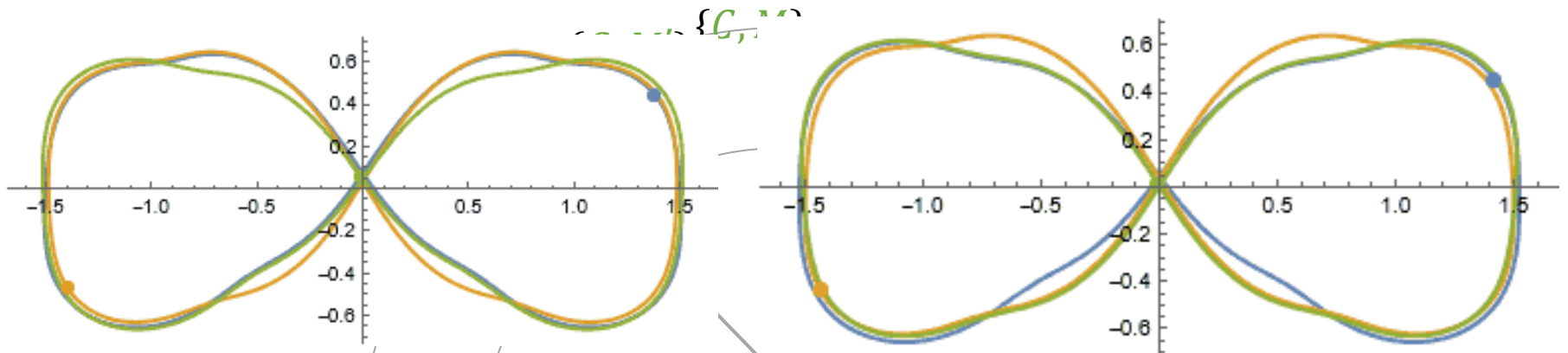
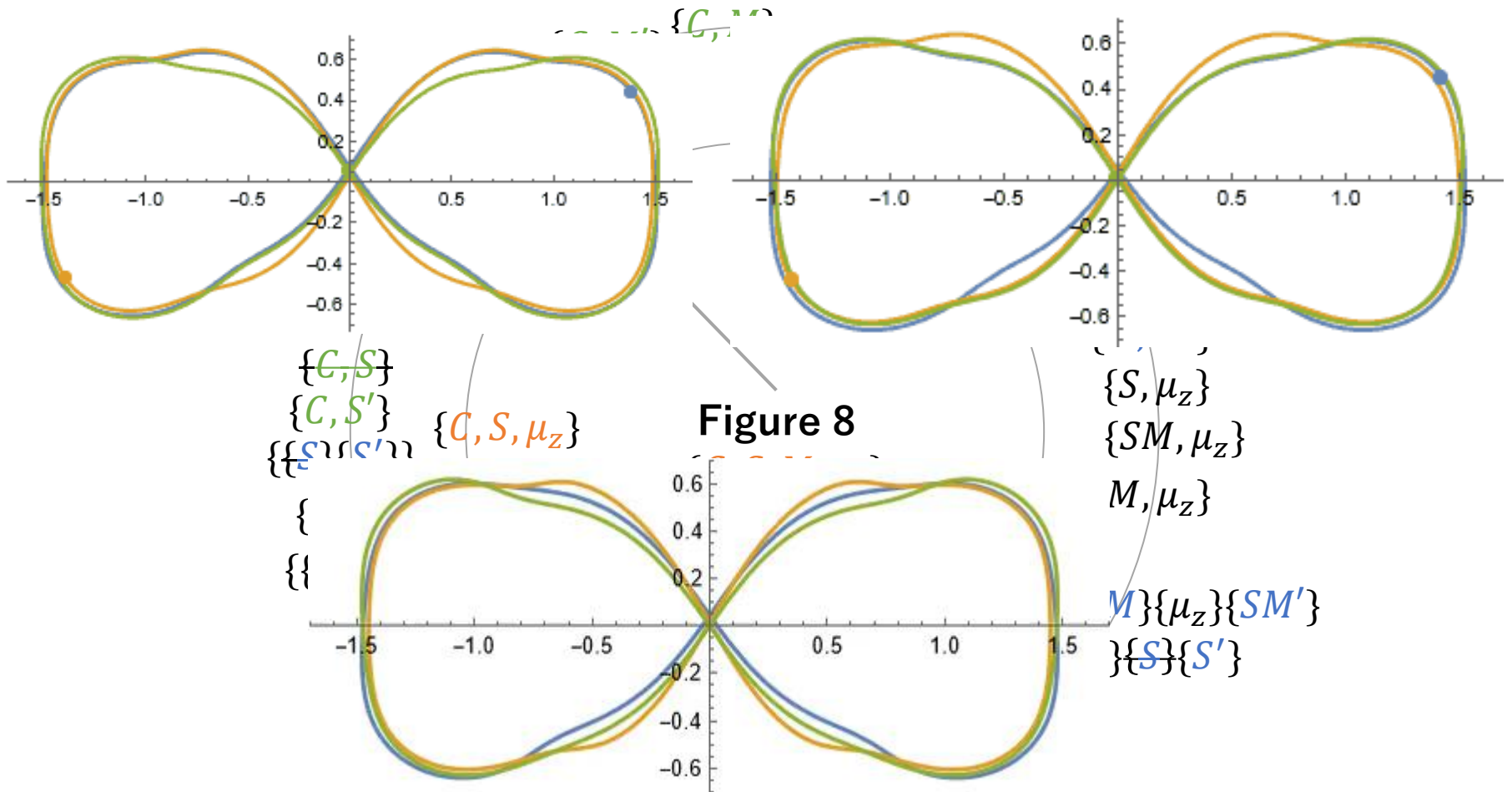


Figure 8

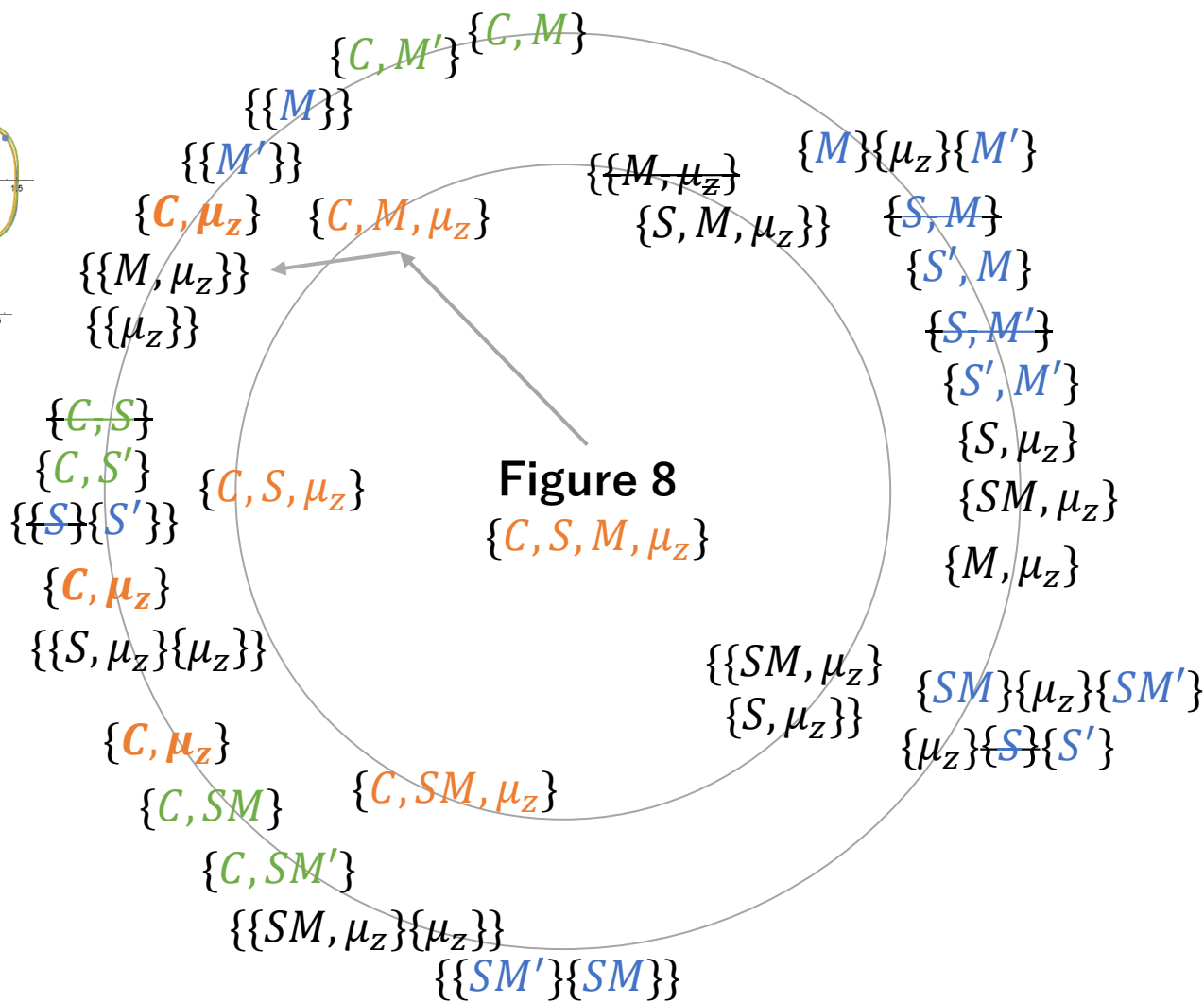
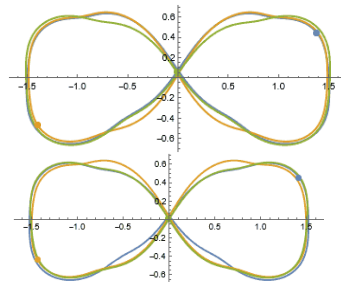
- $\{C, S\}$
- $\{C, S'\}$
- $\{S, S'\}$
- $\{C, \mu_z\}$
- $\{S, \mu_z\}$
- $\{\mu_z\}$
- $\{C, \mu_z\}$
- $\{C, SM\}$
- $\{C, SM'\}$
- $\{SM, \mu_z\}$
- $\{\mu_z\}$
- $\{SM'\}$
- $\{SM\}$
- $\{C, S, \mu_z\}$
- $\{C, S, M, \mu_z\}$
- $\{C, SM, \mu_z\}$
- $\{S, \mu_z\}$
- $\{SM, \mu_z\}$
- $\{SM, \mu_z\}$
- $\{SM\}$
- $\{\mu_z\}$
- $\{SM'\}$
- $\{S\}$
- $\{S'\}$

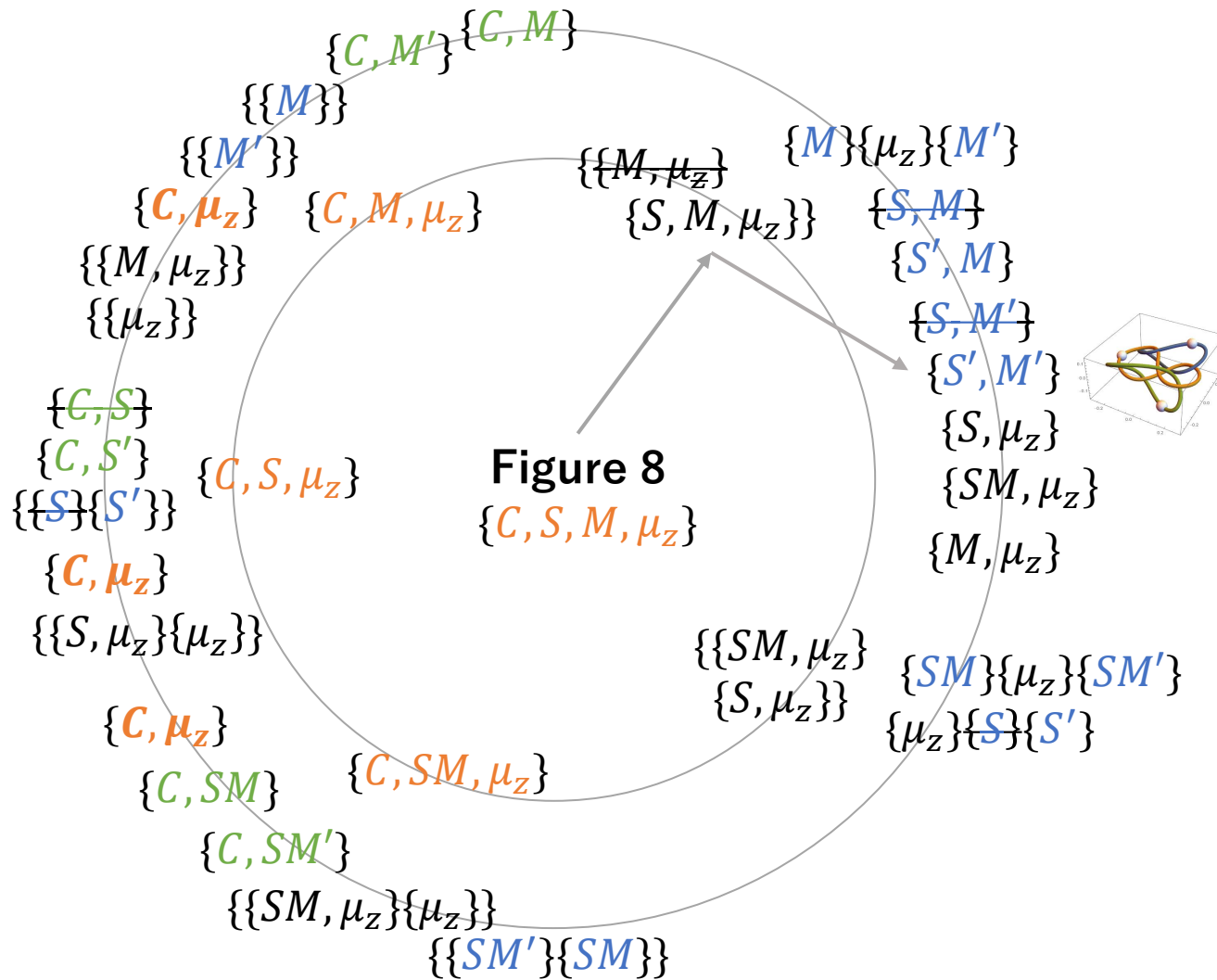
But until now, we do not know what chooses these solution



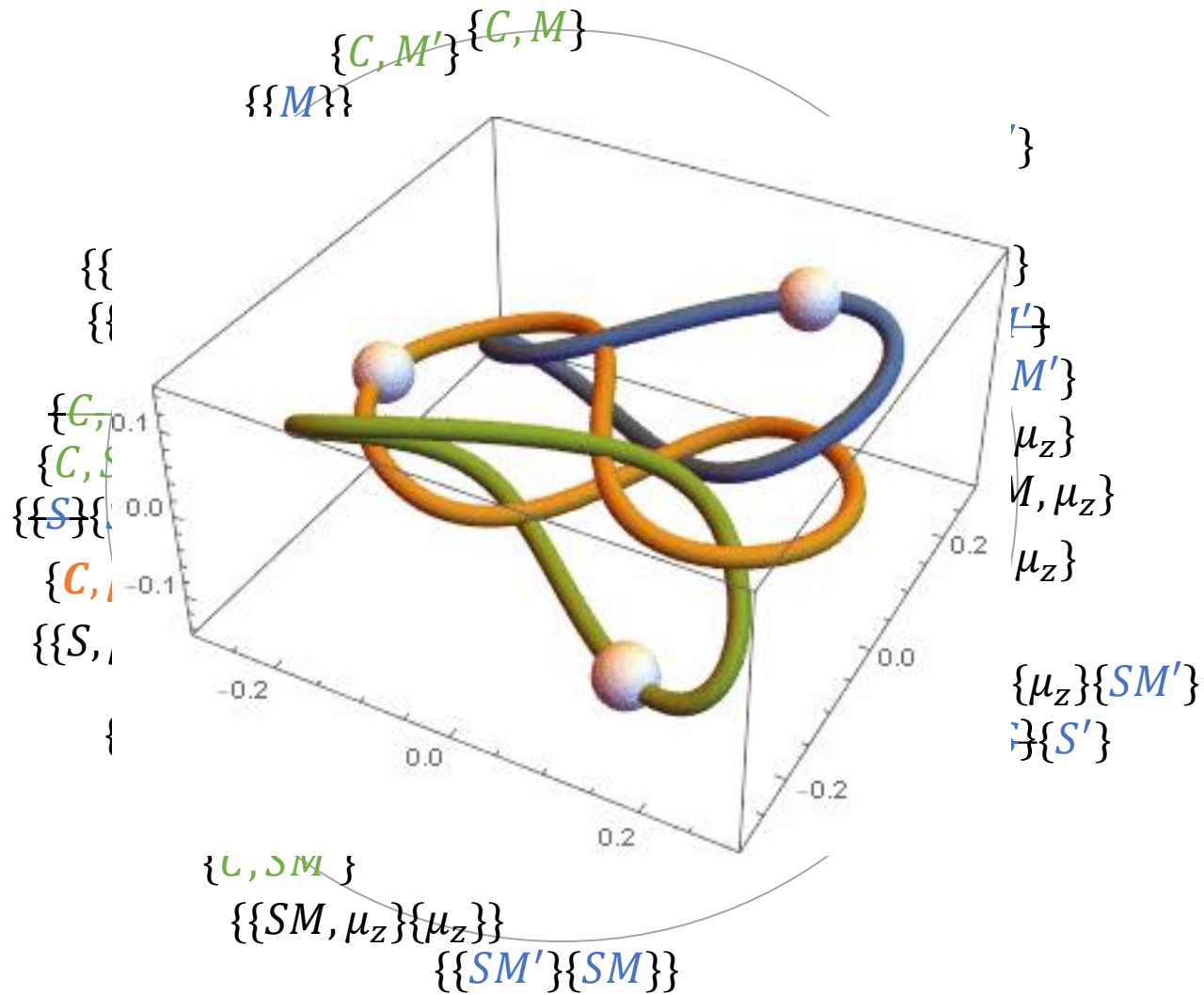
change of $q + \cos\theta\phi_1 + \sin\theta\phi_2$ in θ
 $\{\{SM\}\{SM'\}\}$

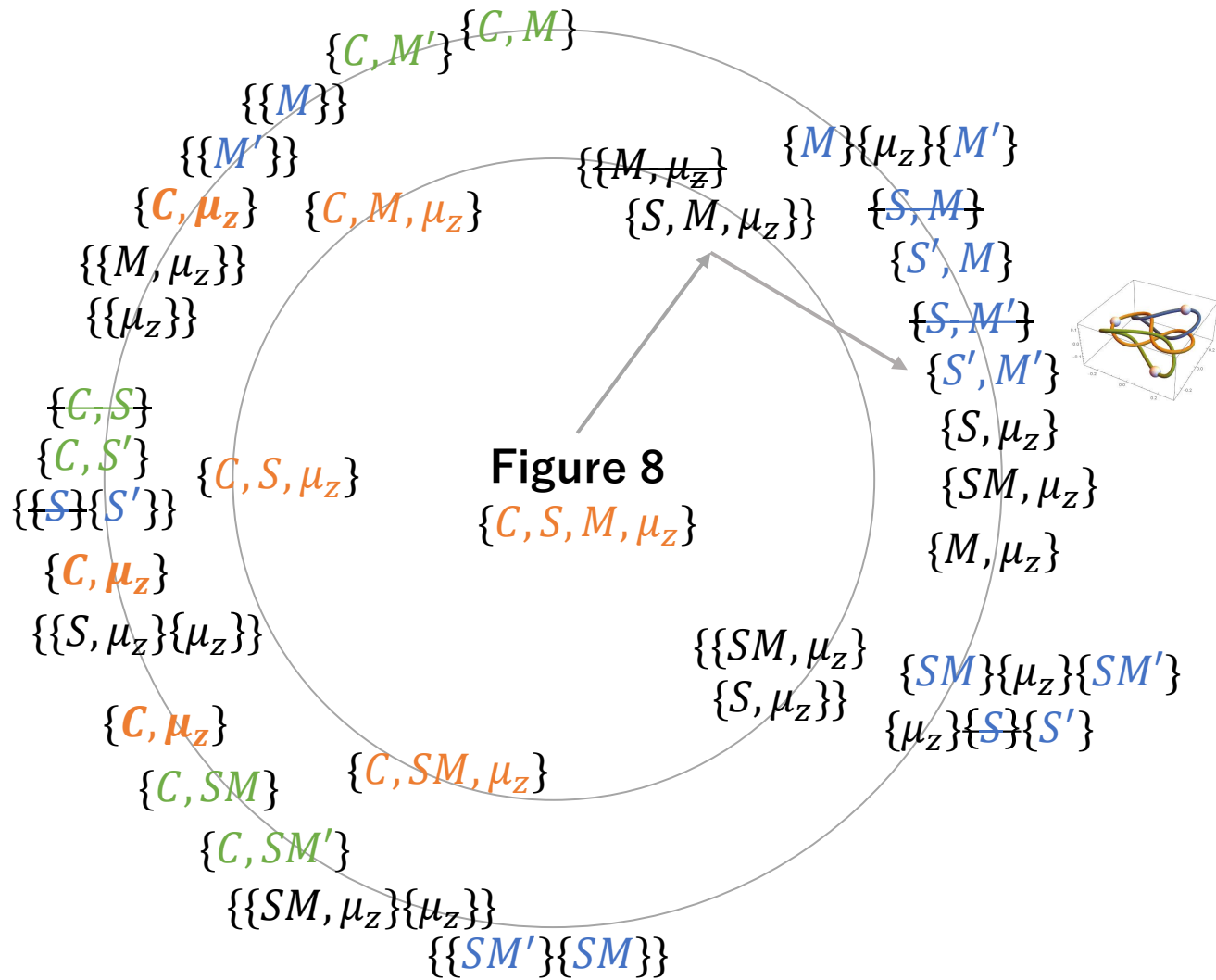
from the linear combinations of eigenfunctions. The movie is showing change of representation in θ .

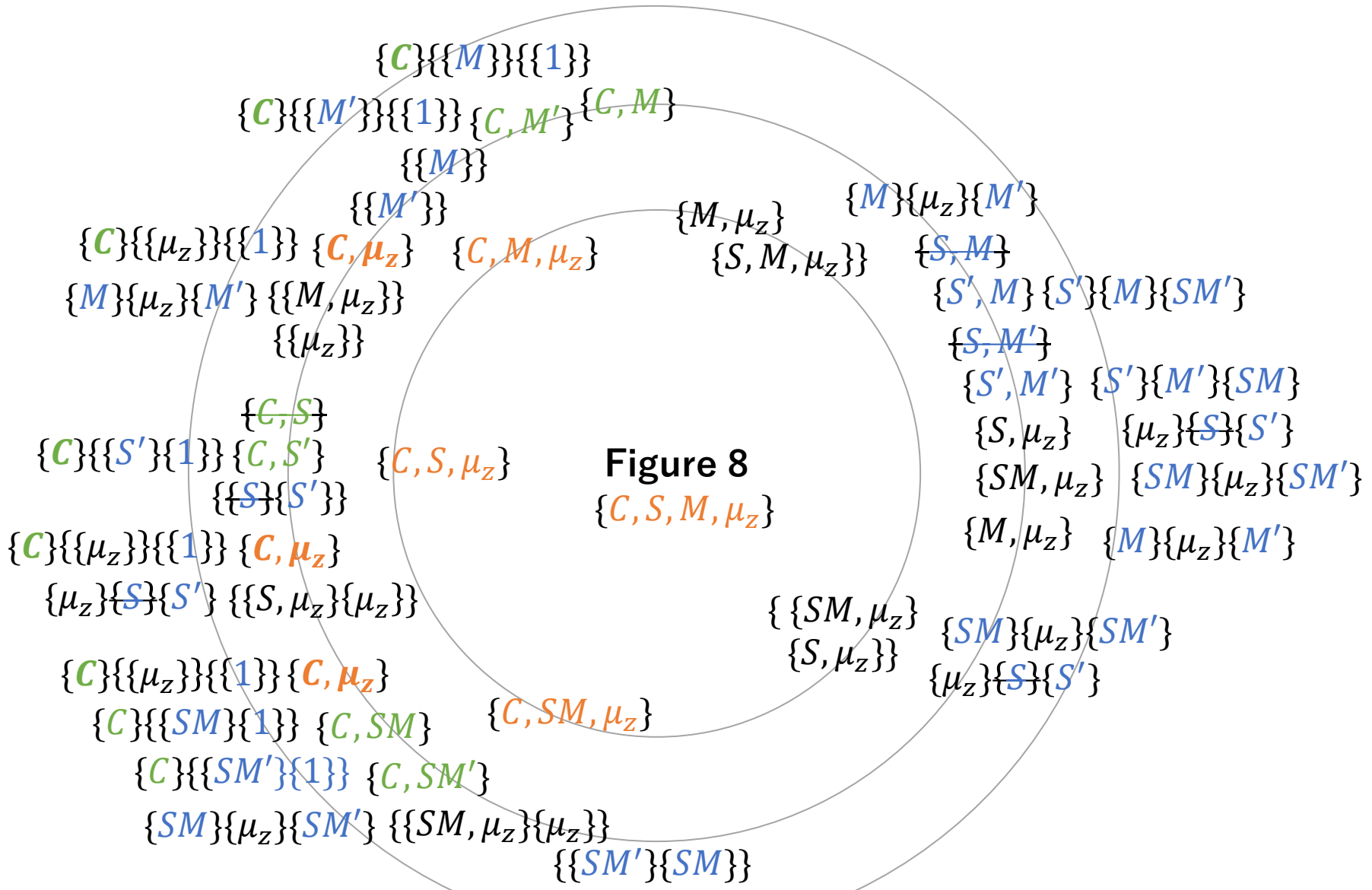




In homogeneous system, we also found interesting bifurcation solutions, a non-planar solution.

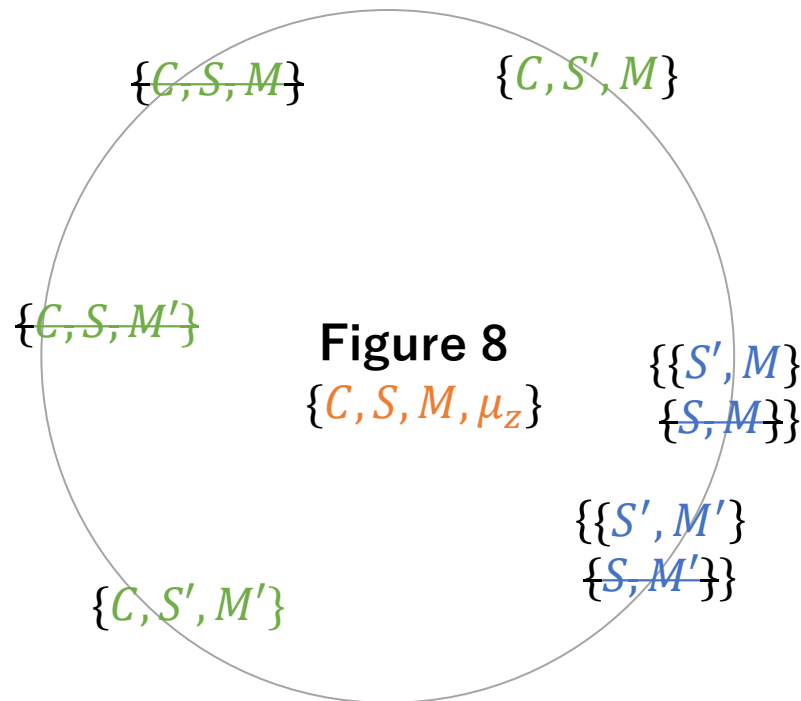




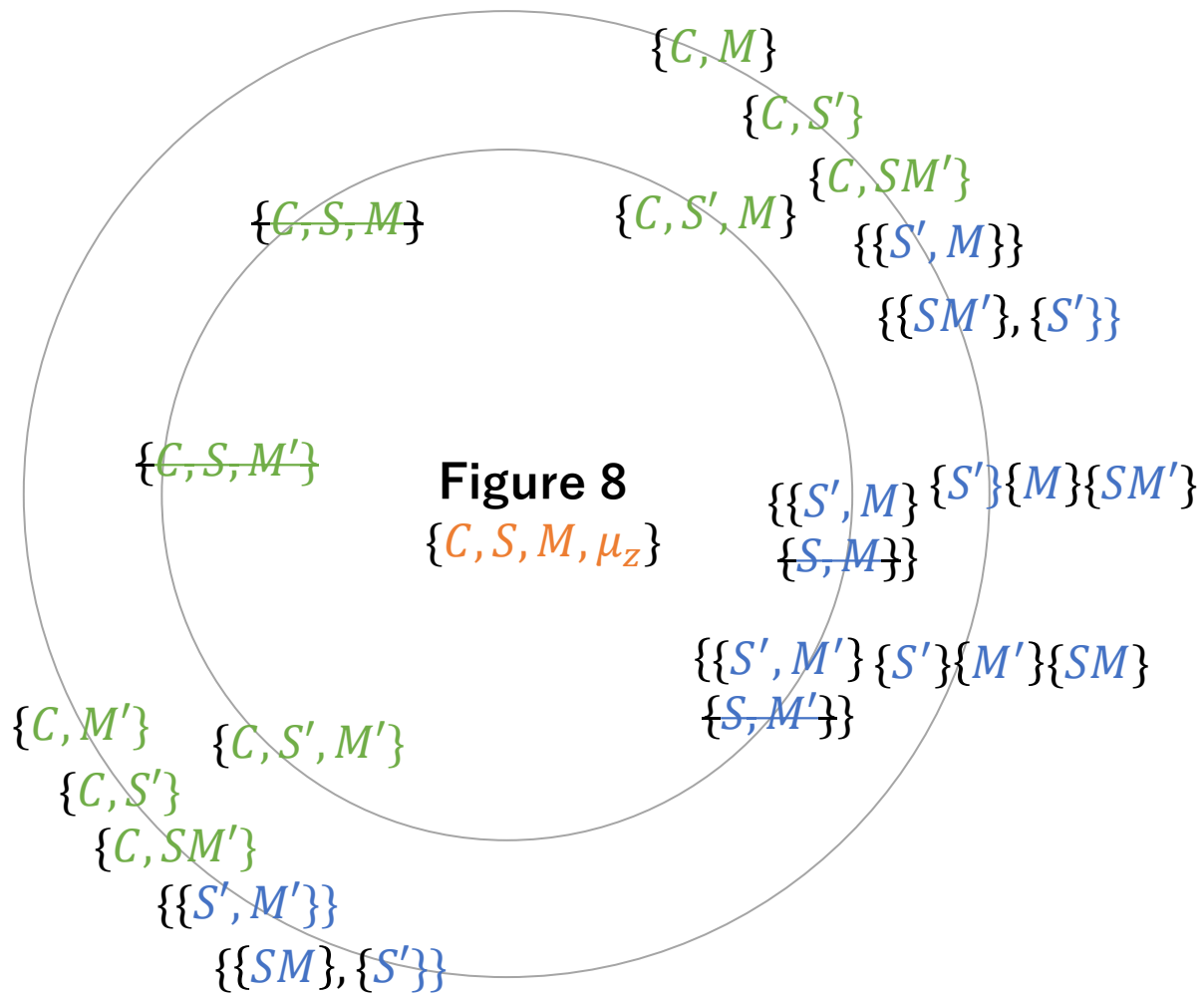


Finally, this is a full diagram for planar bifurcations from figure 8 choreography. On the third ring, the third bifurcations are added.

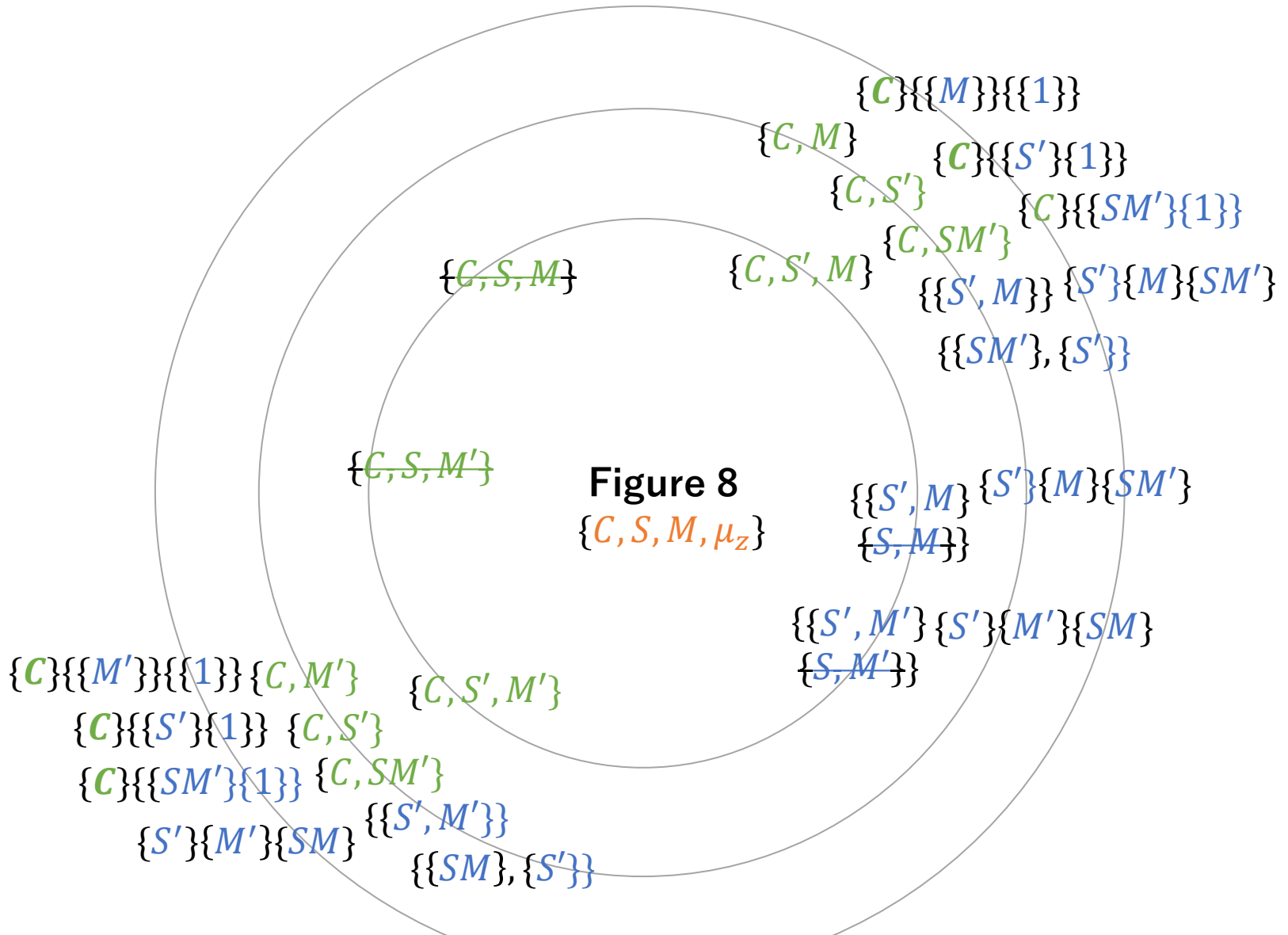
7. Bifurcation of figure eight to the non-planar motion



This is a diagram for direct bifurcations from figure eight choreography to non-planar solutions. Until now we found no this type of solution numerically.



However, they can bifurcate,



and bifurcate.

8. Future work

We want to

- complete numerical search for bifurcation both for homogeneous and LJ system and
- understand bifurcation for two-dimensional representation, especially for degenerate symmetry case.

Thank you.