



## Research article

# Definitive adsorption states of intermediates on Ru nanocatalysts for progress of ammonia synthesis discovered by modulation excitation spectroscopy under reaction conditions

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## ABSTRACT

Diffuse reflectance infrared Fourier transform spectroscopy combined with modulation excitation spectroscopy realized observation of NH<sub>3</sub> production process on a Ru nanoparticle supported on MgO under operation conditions in the presence of both H<sub>2</sub> and N<sub>2</sub> (−400 °C, 0.1 MPa). Our study clearly indicated that N<sub>2</sub> dissociation is developed via transition from vertically adsorbed states of N<sub>2</sub> (N<sub>2</sub>-Ver) on an on-top Ru site to horizontally adsorbed states via N<sub>2</sub>-Ver on a bridge site, which is a critical step for the progress of AS reactions on Ru catalysts at low temperatures and persists in the presence of both N<sub>2</sub> and H<sub>2</sub>.

## 1. Introduction

Ammonia is a versatile basic raw material for the production of fertilizers and pharmaceuticals,[1] and is also expected to be a next-generation carbon-free hydrogen carrier, which is an essential component for efficient long-distance transport of green and blue hydrogen.[2] Thus, reduction of energy consumption through the application of renewable energy to NH<sub>3</sub> production would further enhance the advantage of NH<sub>3</sub> utilization. In this context, ammonia synthesis (AS) catalysts workable under milder conditions are highly demanded. Aika reported that Ru-based catalysts can efficiently promote ammonia synthesis at lower temperatures and lower pressures (350 °C, 0.1 MPa)[3] compared to those on industrially used Harbor-Bosch catalysts (450 °C, 30 MPa).[4] The high efficiency of Ru catalysts was explained by appropriate bond energy between Ru and N atoms, which is neither too high nor too low for the promotion of NH<sub>3</sub> synthesis, and then can exhibit high reaction rates due to smooth adsorption of N<sub>2</sub> and desorption of NH<sub>3</sub> on the surface.[5] To improve the catalytic activity of Ru-based catalysts, metal oxide,[6] electrides,[7] alkali metals,[8] alkali earth metals,[9] lanthanoid oxide,[10] hydride,[11] oxyhydride[12]

and nanoalloys[13] have been applied, which promotes activity or prevents hydrogen poisoning.[14] For further catalytic enhancement, a more specific understanding of reaction progress on the catalyst surface is probably essential.

The adsorption behavior of reactant molecules on AS catalysts has been studied by mass spectrometry employing isotope labelling for post-reaction gases,[15] infrared spectroscopy (IR) and so on.[16] Despite these intense efforts, identifying active species and active sites are still challenging because of their low concentrations and variable natures of the catalytic surface, where multiple catalytic steps are simultaneously developed under reaction conditions. Here, in this study, we demonstrate the time-resolved in-situ IR measurement under reaction conditions (400 °C, 0.1 MPa) to precisely understand the adsorption states of intermediates on Ru AS catalysts, which can be a basic data for the characterization of various types of AS catalysts.

Previously, we observed adsorbed states of N<sub>2</sub> or H<sub>2</sub> on the Ru AS catalyst using diffuse reflectance infrared Fourier transform spectroscopy (DRIFTS) combined with modulation excitation spectroscopy (MES) under nearly reaction conditions (250 °C, 0.1 MPa).[17] MES realizes highly sensitive observation using phase sensitive detection of

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signals responding only to a periodic excitation generated by repeatedly introducing a reaction gas such as  $N_2$  or  $H_2$  and non-reactive reference gas such as Ar, which cancels off-resonant noise signals.[18] Furthermore, dynamics of the adsorbed species on the surface can be extracted by analyzing delays in phase of signals from adsorbed intermediates compared to that from reference gas.[19] We currently constructed a DRIFTS system that is usable at high temperatures (less than 600 °C) in the presence of reactant gases. We then performed time-resolved DRIFTS by introducing 0.1 MPa reaction gas at 400 °C, namely, a reaction condition. Furthermore, co-adsorption behaviors of  $N_2$  and  $H_2$  on the catalyst surface were examined.

## 2. Results and discussion

### 2.1. Structural and catalytic properties of Ru/MgO

MgO-supported Ru nanoparticle (NP) catalysts (Ru/MgO) were prepared accordingly to a previous report (see SI for details).[13] The morphology of Ru NPs on Ru/MgO was observed by transmission electron microscopy (TEM), which confirmed uniform dispersion of Ru NPs on MgO (Fig. 1a and S1). The average particle size of Ru NPs was approximately 5.5 nm (Figure S2). Powder X-ray diffraction (XRD) patterns of Ru/MgO were well reproduced by overlapping weak diffractions originating from a hexagonal close-packed phase of Ru NPs and strong diffractions from a face-centered cubic phase of MgO, suggesting the formation of Ru NPs in the presence of MgO (Figure S3). Scanning electron microscopy together with energy dispersive X-ray spectroscopy revealed that the loading amount of Ru on the catalyst is 4.3 wt%. Catalytic performance on Ru/MgO was evaluated using a gas-flow reactor at 300–500 °C under the mixed gas flow condition (0.1 and

1.0 MPa,  $N_2:H_2 = 1:3$ ). Ru/MgO exhibited the highest activity at 400 °C and decreased above 400 °C at 0.1 MPa whereas the highest performance at 1 MPa was observed at 450 °C (Fig. 1b). The higher catalytic activity on Ru/MgO at higher total pressures accords to reported results.[13,20] and the reaction order of Ru/MgO for  $N_2$ [21] was found to be analogous to the that obtained in previous studies, whereas the value for  $H_2$  was slightly positive compared to our previous result (Figure S4).[13] The difference is possibly explained by considering difference in the loading amount of Ru; 4.3 wt% in this work but 2 wt% in the previous work. This is because a larger number of active sites formed on 4.3 wt%-Ru/MgO somewhat hinders the poisoning compared to that on 2 wt% Ru/MgO. Activation energies estimated from Arrhenius plots at 0.1 and 1.0 MPa (Figure S5, Table S1) were in the range for those reported on Ru-based catalysts.[6] These results suggested Ru/MgO exhibits normal AS activities.

### 2.2. Adsorption of pure $N_2$ or $H_2$

DRIFTS-MES measurements were performed using a Fourier transform IR instrument equipped with gas switching valves (Fig. 1c). Approximately 25 mg of the sample was placed in a porous sample cup and enclosed in the DRIFTS cell. Prior to the measurement, the sample was pretreated under 10 vol%  $H_2$ -Ar gas flow at 600 °C. We kept flowing of both reaction and reference gases at a flow rate of 80 ccm during the measurements, whereas the mixing ratio of reaction and reference gases was optimized at each temperature. The details of the gas mixing ratios are summarized in Table S2. We measured 30 IR spectra for 78 s (or 135 s) with the flow of reaction gas and analogously with reference gas, where the spectra measured at the end in the reference gas period was used as a background, namely 60 spectra were obtained in one series and

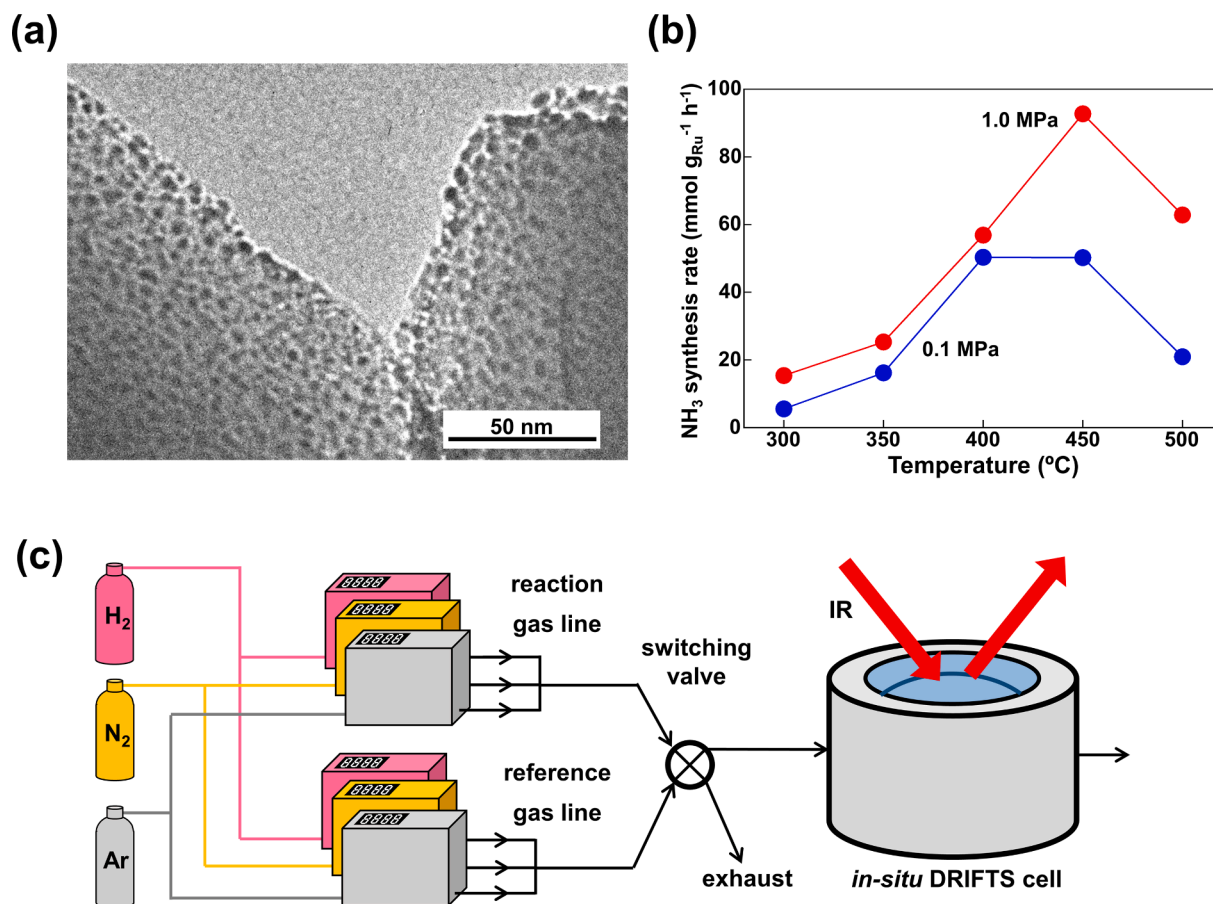


Fig. 1. Sample and measurement setup. (a) TEM image of Ru/MgO. (b) Temperature dependence in catalytic activity for  $NH_3$  synthesis of Ru/MgO. (c) Schematic of the gas-switching type DRIFTS-MES system.

obtained data for 10 series. We produced a spectrum by accumulating signals obtained in the last 8 series.

Fig. 2a shows IR spectra for H<sub>2</sub> adsorption on Ru/MgO after Ar pre-adsorption at 400 °C together with the spectrum previously observed at 250 °C for comparison. Two peaks appeared around 2000 and 1850 cm<sup>-1</sup> at both temperatures (Fig. 2a), which were assigned to adsorptions of H<sub>2</sub> molecules and dissociated H atoms (H-Ru), respectively, based on the result of density functional theory (DFT) calculations.[17] The peak

at 1628 cm<sup>-1</sup> observed at 250 °C was attributed to the bending vibration of OH species on MgO.[17] All signals decreased with time after switching to Ar flow, suggesting that desorption of adsorbed species occurred. It should be emphasized that adsorption and dissociation behaviors of H<sub>2</sub> at 250 and 400 °C looked similar, although no strong signal for the bending vibration of OH species was observed at 400 °C. In the spectra obtained in the presence of N<sub>2</sub>, two obvious peaks appeared around 2071 and 1928 cm<sup>-1</sup> at 250 °C (Fig. 2b), which were already

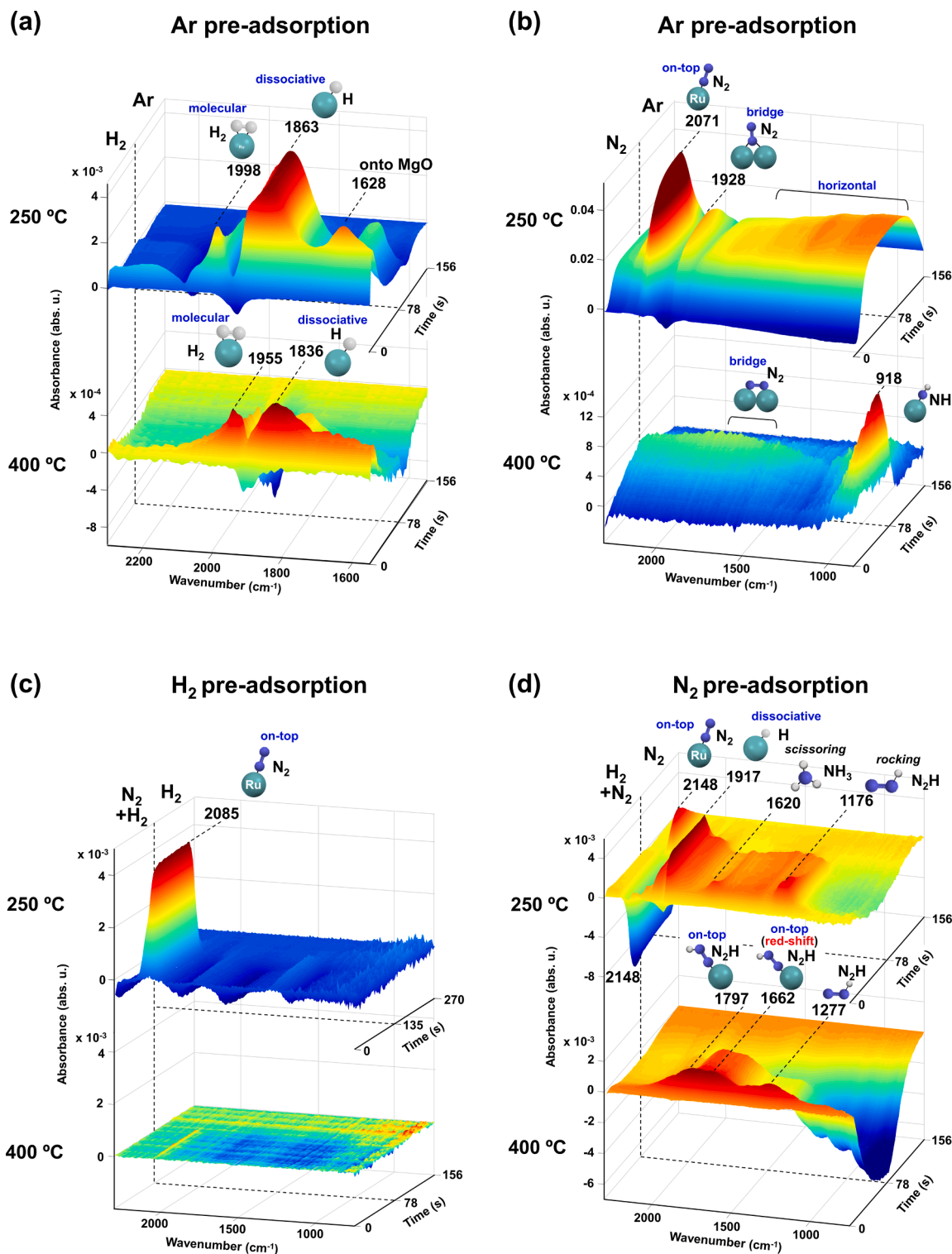


Fig. 2. Time-domain spectra on Ru/MgO surface at 250 and 400 °C under 0.1 MPa for (a) H<sub>2</sub> adsorption after Ar pre-adsorption, (b) N<sub>2</sub> adsorption after Ar pre-adsorption, (c) co-adsorption after H<sub>2</sub> pre-adsorption, and (d) co-adsorption after N<sub>2</sub> pre-adsorption.

attributed to vibrations of vertical adsorbed  $N_2$  ( $N_2$ -Ver) on an on-top and bridge site of Ru NPs, respectively.[17] In addition, weak absorptions below  $1500\text{ cm}^{-1}$  at  $250\text{ }^\circ\text{C}$  were assigned to vibrational states of horizontally adsorbed  $N_2$  ( $N_2$ -Hor).[17] At  $400\text{ }^\circ\text{C}$ , a distinct sharp peak was observed at  $918\text{ cm}^{-1}$ , which did not exist at  $250\text{ }^\circ\text{C}$ . Thus, the adsorption states in the flow of  $N_2$  at  $400\text{ }^\circ\text{C}$  are completely different from those at  $250\text{ }^\circ\text{C}$ . In IR spectrum for  $N_2$  adsorption observed at  $400\text{ }^\circ\text{C}$  in the wider wavenumber range (Figure S6), we can see obvious peaks at 3409, 1088, and  $918\text{ cm}^{-1}$  and a broad peak with a peak around  $1844\text{ cm}^{-1}$ . DFT calculations confirmed that the peak around  $3409\text{ cm}^{-1}$  well matches  $3404\text{--}3471\text{ cm}^{-1}$  for N–H bonds in NH species adsorbed on various sites of Ru(0001). Similarly,  $3467\text{ cm}^{-1}$  was calculated for the antisymmetric vibration of N–H bonds in the  $NH_2$  adsorbed on Ru (0001), and even the vibration of N–H bonds in  $NH_3$  ( $3487\text{ cm}^{-1}$ ) is comparable to the observed value of  $3409\text{ cm}^{-1}$ . Calculated results for the vibrations of  $NH_x$  species on Ru(0001) are summarized in Figure S7. Here, we need to consider H sources for the NH formation. In this experiment, we treated catalysts with  $H_2$  before all measurements and then confirmed the presence of OH groups formed on MgO during the pretreatment, which is the only possible H source in this experiment. Namely, at relatively lower temperatures such as  $250\text{ }^\circ\text{C}$ , OH groups are strongly immobilized on the Ru surface, whereas H atoms of OH groups are possibly dissociated and used for  $NH_3$  formation at  $400\text{ }^\circ\text{C}$ . The peak observed at  $1844\text{ cm}^{-1}$  corresponds to vibrations for H–Ru on the on-top site ( $1852\text{ cm}^{-1}$ ) and for  $N_2$ -Hor on the bridge site ( $1851\text{ cm}^{-1}$ ) of  $Ru_{153}$ , as reported previously.[17] To further assign the undefined 1088 and  $918\text{ cm}^{-1}$  peaks, we focused on the delay of the adsorption intensity in phase, which was obtained by phase sensitive detection (PSD). PSD can notify the delay of response of the signal compared to the time for the introduction of a reaction gas ( $N_2$  and  $H_2$ ), as delay of the signal in phase ( $\Psi_{\text{PSD}}(\theta)$ , see SI) where  $\Psi_{\text{PSD}}(\theta) = 0$  implies that signal intensity completely follows concentration of the reference gas. The  $\Psi_{\text{PSD}}(\theta)$  value corresponds to the time taken to form the bond associated with its wavenumber. Regarding  $N_2$  adsorption (Figure S8), we found IR peaks observed at 3409, 1844, 1088 and  $918\text{ cm}^{-1}$  were characterized with  $\Psi_{\text{PSD}}(\theta) = 53, 85, 99,$  and  $53^\circ$ , respectively. If some peaks are identified with the same  $\Psi_{\text{PSD}}(\theta)$ , it strongly suggests that they are originated from identical adsorbed species. Thus,  $918\text{ cm}^{-1}$  signal with  $\Psi_{\text{PSD}}(\theta) = 53^\circ$  can be correlated to  $3409\text{ cm}^{-1}$  signal with  $\Psi_{\text{PSD}}(\theta) = 53^\circ$  and is accordingly assigned to the N–H stretching vibration. DFT calculation suggested that Ru–N vibration of adsorbed NH species located on the top site of Ru(0001) has a vibrational frequency of  $822\text{ cm}^{-1}$  (Figure S7), which agrees with the experimental value of  $918\text{ cm}^{-1}$ . The DFT calculations suggest that the experimental value of  $918\text{ cm}^{-1}$  corresponds to the Ru–N vibration of NH adsorbed on the top site of Ru(0001), with a calculated vibrational frequency of  $822\text{ cm}^{-1}$  (Figure S7). Although a slightly larger discrepancy exists between the calculated and observed wavenumbers for this specific vibrational mode compared to others, our comprehensive analyses of all possible adsorption sites in our model yield a reasonable conclusion. It is important, however, to acknowledge the challenge of fully matching the real experimental conditions and the idealized calculation model.

We hence distinguished the  $1088\text{ cm}^{-1}$  peak from N–H stretching vibration, whereas DFT calculation suggested that the peak is related to the wagging vibration of  $NH_3$  on Ru(0001), which is characterized with  $1126\text{ cm}^{-1}$  or Ru–N vibration involving an N atom binding to the top site on Ru(0001) with  $1036\text{ cm}^{-1}$  (Figure S7). Thus, the formation of  $NH_x$  species and Ru–N bonds were confirmed at  $400\text{ }^\circ\text{C}$  although all observed peaks were not exactly assigned. Accordingly, we excluded  $N_2$ -Ver on the on-top and bridge sites, implying that  $N_2$  molecules are dissociated on the Ru NP surface at  $400\text{ }^\circ\text{C}$  and AS can proceed via the reaction between dissociated N and H, namely *dissociative* mechanism.[4].

### 2.3. Co-adsorption of $N_2$ or $H_2$

Next, we observed the co-adsorption of  $N_2$  and  $H_2$  by introducing a

mixed gas of  $N_2$  and  $H_2$ . When the mixed gas was introduced after the  $H_2$  pre-adsorption, a strong, sharp peak at  $2085\text{ cm}^{-1}$  was solely observed at  $250\text{ }^\circ\text{C}$  (Fig. 2c top), whereas somewhat negative intensities were observed below  $2000\text{ cm}^{-1}$ , suggesting that some adsorption sites for H species are used for  $N_2$  adsorption. The  $2085\text{ cm}^{-1}$  peak is assignable to  $N_2$ -Ver on the on-top Ru site,[17] similarly to  $2071\text{ cm}^{-1}$  peak observed after introducing Ar (Fig. 2b top). We then exemplified the effect of the dissociated H atoms on the adsorption energy and N–N vibrational frequencies of  $N_2$ -Ver as shown in Figure S9, revealing that the N–N vibrational frequency increases when H atoms approach a vertically adsorbed  $N_2$ . [22] Thus, we can say that the blue shifted  $2085\text{ cm}^{-1}$  peak for  $N_2$ -Ver on the on-top site under the mixed gas condition suggests certain interaction of the  $N_2$  with an H atom that withdraws electrons from a Ru NP, which hinders the dissociation of the N–N bond. This result was already predicted by the previous DFT calculations for the  $Ru_{153}$  system.[22] Furthermore, the signal intensity for  $N_2$ -Ver at  $2085\text{ cm}^{-1}$  decreased more quickly than that of the  $2071\text{ cm}^{-1}$  peak observed after the Ar pre-adsorption (Fig. 2b top), which also suggests a weak interaction between  $N_2$  and the  $H_2$ -treated surface. Co-adsorption after the  $H_2$  pre-adsorption showed only very weak signals around  $2000\text{ cm}^{-1}$  at  $400\text{ }^\circ\text{C}$  (Fig. 2c), which suggests that  $NH_3$  formation hardly occurs even at relatively high temperatures on the surface that is once exposed to pure  $H_2$ . Thus, the Ru surface fully covered with H species cannot activate  $N_2$  molecules for the progress of AS reaction even at  $400\text{ }^\circ\text{C}$ .

In contrast, various absorption peaks were observed when a mixed gas was introduced after the pre-adsorption with  $N_2$  at both temperatures. DFT calculations suggested that a strong peak observed around  $1917\text{ cm}^{-1}$  at  $250\text{ }^\circ\text{C}$  (Fig. 2d top) is assignable to H–Ru on the top of a Ru vertex atom ( $1921\text{ cm}^{-1}$ ), which is located in close proximity to another H atom. The H–H distance corresponding to this frequency was calculated to be  $2.114\text{ \AA}$ , as shown in Figure S10. In addition, we found weak peaks around  $1620$  and  $1176\text{ cm}^{-1}$  at  $250\text{ }^\circ\text{C}$  (Fig. 2d top), which accord with wavenumbers predicted in the previous DFT studies;  $1642\text{--}1659\text{ cm}^{-1}$  for some configurations of  $N_2$ -Hor on bridge sites of  $Ru_{153}$  and  $1162\text{--}1183\text{ cm}^{-1}$  for a few  $N_2$  molecules interacting on the 4F sites of  $Ru_{153}$ . [17] Similarly,  $1164\text{--}1187\text{ cm}^{-1}$  for the H–Ru formed on the 3F sites of  $Ru_{153}$  is also another possible mode.[22] Meanwhile, DFT calculations conducted in this study for  $NH_3$  adsorbed on the top site of Ru(0001) showed that scissoring ( $1581\text{ cm}^{-1}$ ) vibrational modes of  $NH_3$  are in good agreement with the observed peaks at  $1620\text{ cm}^{-1}$  (Fig. 2d top). Similarly, we found that the  $1176\text{ cm}^{-1}$  peak is related to rocking vibrations of  $N_2H$  horizontally adsorbed on the  $Ru_{153}$  surface ( $N_2H$ -Hor), of which vibrational wavenumbers were calculated to be in the range from  $1112$  to  $1195\text{ cm}^{-1}$ . The detailed DFT results are summarized in Fig. S11a. Thus, AS reactions are developed via the formation of horizontally adsorbed  $N_2H$  at  $250\text{ }^\circ\text{C}$  as shown in Fig. 3a, which is denoted as *horizontally associative* mechanism.

It should be noted that a sharp negative peak and broad negative intensities were observed at  $2148\text{ cm}^{-1}$  and below  $1000\text{ cm}^{-1}$  at  $250\text{ }^\circ\text{C}$  in the flow of the mixed gas after the  $N_2$  pre-adsorption (Fig. 2d top). Based on a previous report, we have correlated the  $2148\text{ cm}^{-1}$  peak to  $N_2$ -Ver on a facet of Ru NPs, namely, wavenumbers were calculated to be in the range from  $2163$  to  $2192\text{ cm}^{-1}$ , [17] which is larger than those for  $N_2$ -Ver on a vertex. It should be noted that we did not observe  $N_2$ -Ver on a facet after  $H_2$  pre-adsorption (Fig. 2c top). This is because all absorption sites on facets are already occupied by H species during  $H_2$  pre-adsorption. The negative regions below  $1000\text{ cm}^{-1}$  correspond to vibrations for  $N_2$ -Hor at 3F or 4F sites,[17] where  $N_2$  molecules are highly activated. Thus, it is acceptable that vertically adsorbed  $N_2$  molecules on a facet are transferred to horizontally adsorbed ones and then converted to NH species. Interestingly, the negative absorbance at  $2148\text{ cm}^{-1}$  became positive just after switching to pure  $N_2$  (Fig. 2d top), then decreased to zero at 156 s whereas the constantly positive signals were observed below  $1100\text{ cm}^{-1}$ . Thus, we found the Ru sites on the facet becomes vacant during  $NH_3$  formation in the presence of both  $N_2$  and  $H_2$  and can act as recoverable adsorption sites for  $N_2$ , which drives



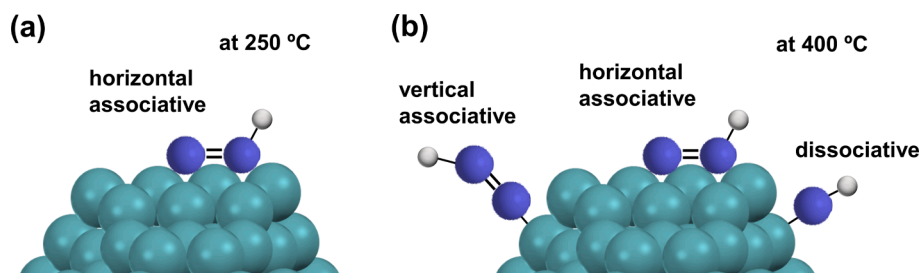


Fig. 3. Possible intermediate species in AS reactions at (a) 250 °C and (b) 400 °C.

the continuous AS reaction.

At 400 °C, we could see a peak around 1277 cm<sup>-1</sup>, which can be assigned to the N–N bond vibration (1284 cm<sup>-1</sup>) of N<sub>2</sub>H-Hor at a bridge site between a vertex and edge atoms of Ru<sub>153</sub> (Figure S11a). However, a broad peak having a maximum of around 1797 and 1662 cm<sup>-1</sup> did not match the results from the DFT calculations for H, H<sub>2</sub>, NH<sub>3</sub>, NH<sub>2</sub>, and NH species on Ru. Instead, DFT calculations could assign both two peaks to the N–N bond vibration of the vertically adsorbed N<sub>2</sub>H (N<sub>2</sub>H-Ver) on a vertex atom of Ru<sub>153</sub> (Figure S11b). Depending on the position of remaining dissociated H atoms, the N–N vibrational frequency can become more red-shifted to 1688 cm<sup>-1</sup>, which is in good agreement with the experimental peak at 1662 cm<sup>-1</sup>. These results suggest that AS reactions proceed not only via *horizontally associative* mechanism but also via the formation of vertically adsorbed N<sub>2</sub>H species, namely, *vertically associative* mechanism, at 400 °C (Fig. 3b).

As we already mentioned, we could not observe vibration for N<sub>2</sub>-Ver at 400 °C after the Ar and N<sub>2</sub> pre-adsorption (Fig. 2b and 2d bottom), suggesting the lifetime of N<sub>2</sub> species related to N<sub>2</sub>-Ver is considerably short.

#### 2.4. AS reaction process developed on Ru/MgO

Fig. 4a shows time-dependent change in absorbance of the 1176 cm<sup>-1</sup> peak observed under the mixed gas conditions at 250 °C, which was assigned to N<sub>2</sub>H-Hor on Ru (Fig. 2d, top). Remarkably, the absorbance reached the maximum at 13 s and became constant, implying that the system reached a steady state under the current reaction conditions. Then, we can draw a picture for the progress of AS reaction on Ru/MgO. In the presence of N<sub>2</sub> and H<sub>2</sub>, continuous NH<sub>3</sub> formation is realized because 3 or 4 vacant sites are produced because three H and one N species are removed from the adsorption sites such as three Ru on-top sites for H-Ru, one for N<sub>2</sub>-Ver and two for N<sub>2</sub>-Hor during NH<sub>3</sub>

formation (Fig. 4b). Thus, we assume that the progress of AS reactions is relevant to the transformation of vertically adsorbed N<sub>2</sub> from an on-top to a bridge site, which is composed of two Ru atoms. Fig. 5a describes the time-dependent change in the absorbance of peaks which were assigned to N<sub>2</sub>-Ver on the on-top of a facet or edge, N<sub>2</sub>-Ver on the bridge, H-Ru, and NH<sub>3</sub>. Interestingly, the intensity for N<sub>2</sub>-Ver on the bridge exhibited similar trend to those for H-Ru and NH<sub>3</sub>, implying that N<sub>2</sub>-Ver on the bridge is relevant to NH<sub>3</sub> formation, whereas the other two signals for N<sub>2</sub>-Ver on the on-top sites showed different trends. To quantitatively evaluate the amount of N<sub>2</sub> intermediate species, Kubelka-Munk transform for the signal intensity was performed (inset of Fig. 5a), and the line-fitting in the range from 115 and 143 s was conducted. The results showed that the inclinations of the signal intensity for N<sub>2</sub>-Ver on the on-top of Ru atom located on a facet ( $\Delta_{\text{top}}$ ) and the bridge ones ( $\Delta_{\text{bridge}}$ ) are  $-4.79 \times 10^{-4}$  and  $2.08 \times 10^{-4}$  K-M unit/s, respectively, with a ratio of  $\Delta_{\text{top}}/\Delta_{\text{bridge}} = -2.30$ . This value matches the model where the elimination of two adsorption species on the on-top sites is needed for emergence of N species on one bridge site (Fig. 5b). In addition, our previous PSD analysis suggested that increase of signal intensity for N<sub>2</sub>-Ver on the on-top site is considerably faster and slow transition of vertically adsorbed N<sub>2</sub> from an on-top to bridge site occurs.[17] Thus, we think that transformation of N<sub>2</sub> from the on-top to bridge site is very important for the formation of horizontally adsorbed N<sub>2</sub>, which is critical species for the progress of AS on Ru/MgO at 250 °C.

### 3. Conclusion

We have successfully observed adsorbed species changing with time on Ru/MgO under reaction conditions by in-situ DRIFTS-MES and DFT calculations, which elucidates not only the types of intermediates but also their adsorption states and sites on Ru/MgO. In particular, different thermal conditions were found to involve different reaction pathways of

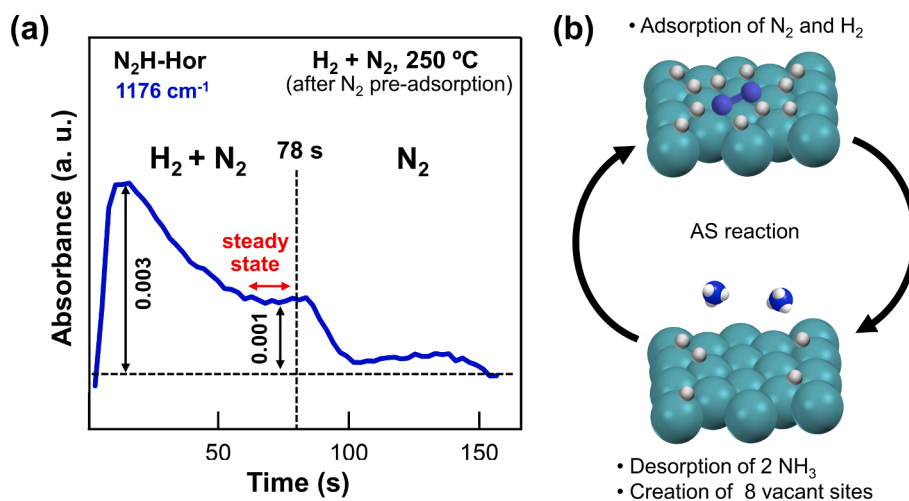


Fig. 4. (a) Time dependence of the absorbance for wagging NH<sub>3</sub> peak (1176 cm<sup>-1</sup>) under the mixed gas condition at 250 °C after N<sub>2</sub> pre-adsorption. (b) Schematic image of the progress of AS reaction.

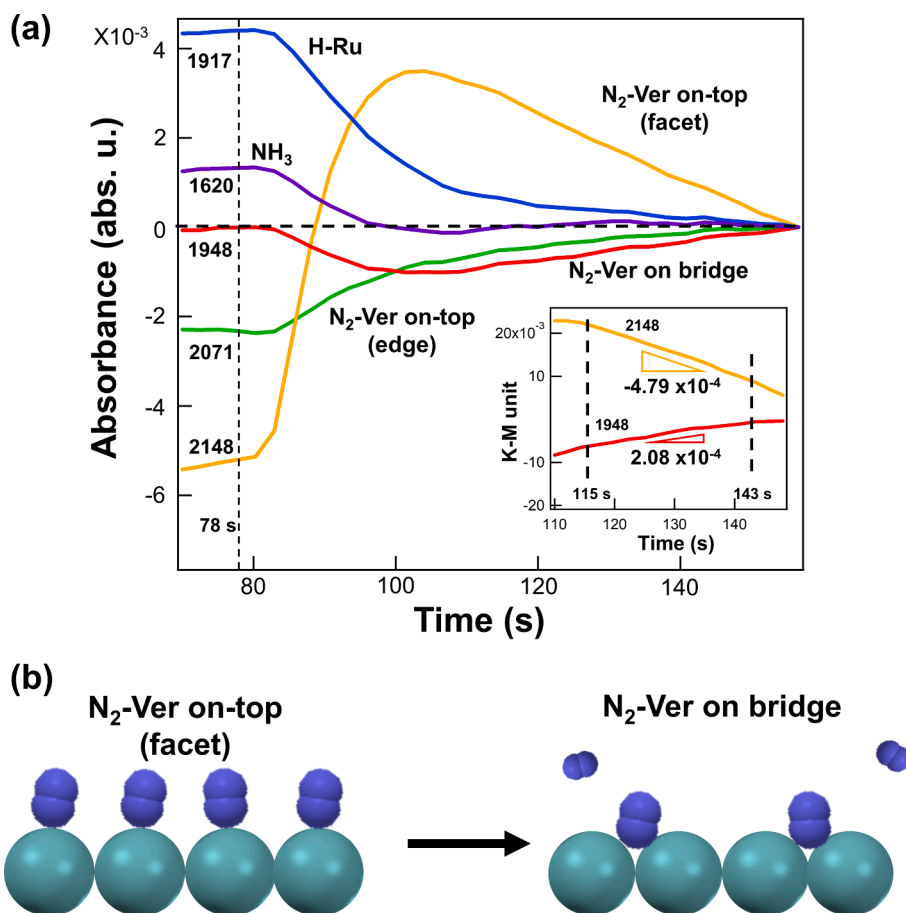


Fig. 5. (a) Time dependence of the absorbance for each adsorption state after the introduction of mixed gas at 250 °C (Fig. 2d top). Intensity of 2148 and 1948  $\text{cm}^{-1}$  peaks in Kubelka-Munk unit, and their inclinations from 115 to 143 s are shown in inset. (b) Schematic image for transfer of vertically adsorbed  $\text{N}_2$  from the on-top (facet) to bridge sites.

AS reaction on Ru NPs; only the **horizontally associative** mechanism runs at 250 °C, whereas three considerable mechanisms such as **dissociative**, **vertically** and **horizontally associative** mechanisms are developed at 400 °C. Therefore, the catalyst design for accelerating **horizontally associated** mechanism is a key to promote AS reaction at lower temperatures. Furthermore, our PSD analysis revealed that the formation of vertically adsorbed  $\text{N}_2$  on the bridge site is a hidden step to accelerate AS reaction. These exact descriptions about intermediate species dynamically changing on the catalyst surface would strongly support for rational design to create highly efficient catalysts.

#### 4. Computational details

All calculations were based on the plane-wave DFT method implemented in the Vienna *ab initio* simulation package (VASP 5.6.2). [23] Perdew – Burke – Ernzerhof parametrization under the generalized gradient approximation was employed as the exchange–correlation functional together with the projector-augmented wave method. [24] Spin-polarized calculations were performed throughout the study with a plane-wave cutoff energy of 600 eV. The convergence criterion for all calculations was set as the point at which the difference in the total energy between the two ionic steps was less than  $10^{-5}$  eV/atom ( $10^{-6}$  eV/atom for self-consistent-field iterations). The optimization of the hexagonal close-packed (hcp) bulk Ru was performed with  $16 \times 16 \times 10$  Monkhorst – Pack k-point mesh for the Brillouin zone integration, where all the atoms and the crystal volume were allowed to relax. After the optimization, the calculated lattice parameters,  $a = 2.715$  and  $c = 4.279$  Å, were in good agreement with the experimental values  $a =$

2.706 and  $c = 4.282$  Å. [25] Subsequently, the bulk was cleaved, and the Ru(0001) slab was constructed with seven atomic layers, from which the coordinates of the atoms in the bottom three layers of the slabs were kept frozen while the coordinates of the atoms in the four top layers were fully relaxed. The vacuum between the slabs was set to span a range of 17 Å to ensure no significant interaction between the slabs. The choice of Ru(0001) as the surface orientation is based on its exceptional stability among Ru surfaces. Previous studies have successfully employed a nanoparticle model consisting of two facets, {0001} and  $\{10\bar{1}1\}$  to identify the adsorption states of  $\text{H}_2$  and  $\text{N}_2$ . [17] Additionally, this model has helped in understanding the impact of H poisoning on Ru NPs. [22]. The optimization of the Ru(0001) was performed with  $4 \times 4 \times 1$  Monkhorst – Pack k-point mesh for the Brillouin zone integration. The interaction of the  $\text{NH}_x$  species was approximated by the adsorption energy,  $E_{\text{ads}}$ , calculated as:

$$E_{\text{ads}} = E_{\text{Ru+adsorbate}} - E_{\text{Ru}} - E_{\text{adsorbate}} \quad (1)$$

where  $E_{\text{Ru+adsorbate}}$ ,  $E_{\text{Ru}}$ , and  $E_{\text{adsorbate}}$  denote the total energies of  $\text{NH}_3$ ,  $\text{NH}_2$ ,  $\text{NH}$ , or  $\text{N}$  atoms interacting with the Ru(0001), the total energy of the pristine Ru(0001), and the total energies of isolated  $\text{NH}_3$ ,  $\text{NH}_2$ ,  $\text{NH}$ , and a  $\text{N}$  atom, respectively. According to this definition, negative adsorption energies denote more stable adsorbate–Ru interactions.

The optimization of the  $\text{N}_2\text{H}$  interacting with the hcp structure of the 1.58 nm diameter Ru nanoparticle,  $\text{Ru}_{153}$ , was performed at the  $\Gamma$  point in reciprocal space owing to the significant spatial extent of the system, where all the atoms were allowed to relax. Similarly for the interaction of  $\text{N}_2$  on  $\text{Ru}_{153}$  pre-covered with two H atoms. The distance between neighboring image Ru nanoparticles was set as 15 Å to avoid interaction

between periodic images. The  $E_{\text{ads}}$  of  $\text{N}_2\text{H}$  on  $\text{Ru}_{153}$  was calculated as in equation (1), where  $E_{\text{Ru+adsorbate}}$ ,  $E_{\text{Ru}}$ , and  $E_{\text{adsorbate}}$  denote the total energies of  $\text{N}_2\text{H}$  interacting with the  $\text{Ru}_{153}$ , the total energy of the  $\text{Ru}_{153}$ , and the total energies of isolated  $\text{N}_2\text{H}$ , respectively. The structures of the  $\text{N}_2$  together with the two H atoms binding to  $\text{Ru}_{153}$  are shown in Figure S11a.

The  $E_{\text{ads}}$  of  $\text{N}_2$  on the H precovered  $\text{Ru}_{153}$  was calculated as:

$$E_{\text{ads}} = E_{\text{Ru+2H+N}_2} - E_{\text{Ru+2H}} - E_{\text{N}_2}, \quad (2)$$

where  $E_{\text{Ru+2H+N}_2}$ ,  $E_{\text{Ru+2H}}$ , and  $E_{\text{N}_2}$  denote the total energies of  $\text{N}_2$  interacting with the  $\text{Ru}_{153}$  pre-covered with two H atoms, the total energy of the  $\text{Ru}_{153}$  with two H atoms, and the total energy of isolated  $\text{N}_2$ , respectively. The structures of the  $\text{N}_2$  together with the two H atoms binding to  $\text{Ru}_{153}$  are shown in Figure S9.

The vibrational calculations were performed using central differences with atomic displacements of  $\pm 0.015 \text{ \AA}$  in all three Cartesian directions. The dynamical matrix was constructed by displacing all the adsorbates' atoms that are interacting with the Ru nanoparticle and Ru (0001). For the  $\text{NH}_x$  species adsorbed on Ru(0001), the N-Ru atoms were also displaced. The self-consistent-field iterations convergence criterion was set to  $10^{-8} \text{ eV/atom}$  for all the vibrational calculations.

### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

### Data availability

No data was used for the research described in the article.

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### Appendix A. Supplementary material

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.jcat.2023.07.020>.

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