

生成関数によるテンソルネットワークの和の 計算方法

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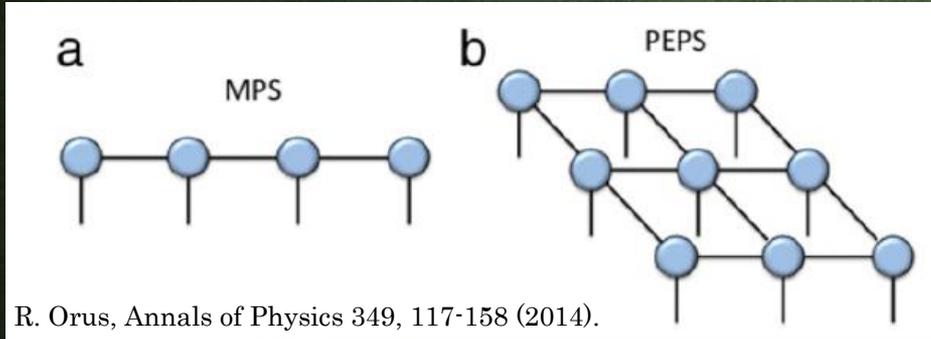
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WLT *et al.*, arXiv:2101.03935 (2021)

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Tensor network:



R. Orus, Annals of Physics 349, 117-158 (2014).

One-particle excitation:

$$|\Phi_k(B)\rangle = \sum_{j=0}^{N-1} e^{-ikj} \hat{T}^j \left(\begin{array}{c} \text{---} \text{[B]} \text{---} \text{[A]} \text{---} \dots \text{---} \text{[A]} \text{---} \\ | \quad | \quad \quad | \quad \quad | \\ s_1 \quad s_2 \quad \dots \quad s_N \end{array} \right)$$

Static structural factor:

$$S^{\alpha,\beta}(k) = \frac{1}{N} \sum_{j,j'=1}^N e^{ik \cdot (r_j - r_{j'})} \langle \hat{O}_j^\alpha \hat{O}_{j'}^\beta \rangle$$

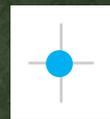
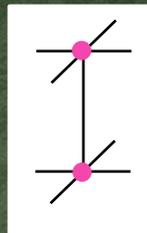
- Tensor network ansatz has become one of the reliable numerical approaches for quantum many-body systems.
- After obtaining the ground state ansatz, we need to calculate some physical observables or solve for low-energy excited states for further analysis.
- Within such evaluation, the summation of a series of tensor graphs is needed, which results in a large computational cost.

Spin static structural factor in 2D:

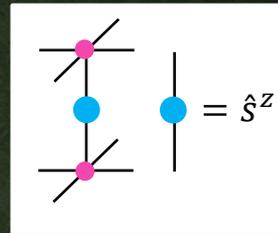
$$S^Z(k) = \frac{1}{N} \sum_{j \neq j'}^N e^{ik \cdot (r_j - r_{j'})} \langle \hat{S}_j^Z \hat{S}_{j'}^Z \rangle$$



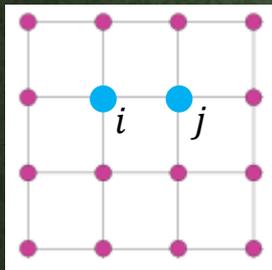
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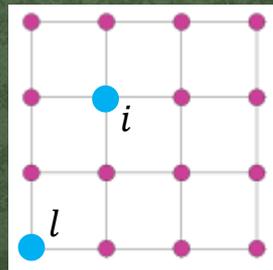
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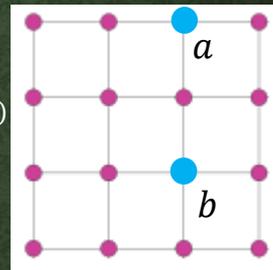
$$= \frac{1}{N} e^{ik \cdot (r_i - r_j)}$$



$$+ e^{ik \cdot (r_i - r_l)}$$



$$+ e^{ik \cdot (r_a - r_b)}$$



+ ...

$N(N - 1)$ copies

- Therefore, in the traditional way one needs to consider many distinct tensor graphs for the final summation.
- However, for a large system or large virtual bond dimension, calculation for single graph can be already cumbersome.

Generating function (GF):

- In field theory, this function is often constructed and by taking the derivative, one can obtain the target values.
- Here, we are going to borrow the same idea and apply it for the tensor network ansatz.

Taking 1D system as an example
(with translational symmetry):

One-particle excitation:

$$|\Phi_k(B)\rangle = \sum_{j=0}^{N-1} e^{-ikj} \hat{T}^j \left[\begin{array}{c} \text{---} \text{---} \text{---} \text{---} \text{---} \text{---} \\ | \quad | \quad | \quad | \quad | \\ \text{---} \text{---} \text{---} \text{---} \text{---} \end{array} \right]$$

$$\frac{\partial |G_\Phi(\lambda)\rangle}{\partial \lambda} \Big|_{\lambda=0}$$



Corresponding GFs:

$$|G_\Phi(\lambda)\rangle = \left[\begin{array}{c} \text{---} \text{---} \text{---} \text{---} \text{---} \text{---} \\ | \quad | \quad | \quad | \quad | \\ \text{---} \text{---} \text{---} \text{---} \end{array} \right]$$

$$MPS_j(\lambda) = A + \lambda e^{-ikr_j} B$$

Static structural factor:

$$S^{\alpha,\beta}(k) = \sum_{j=1}^N e^{ik \cdot (r_1 - r_j)} \langle \hat{O}_1^\alpha \hat{O}_j^\beta \rangle$$

$$\left\langle \frac{\partial \hat{G}_{SF}(\lambda)}{\partial \lambda} \Big|_{\lambda=0} \right\rangle$$

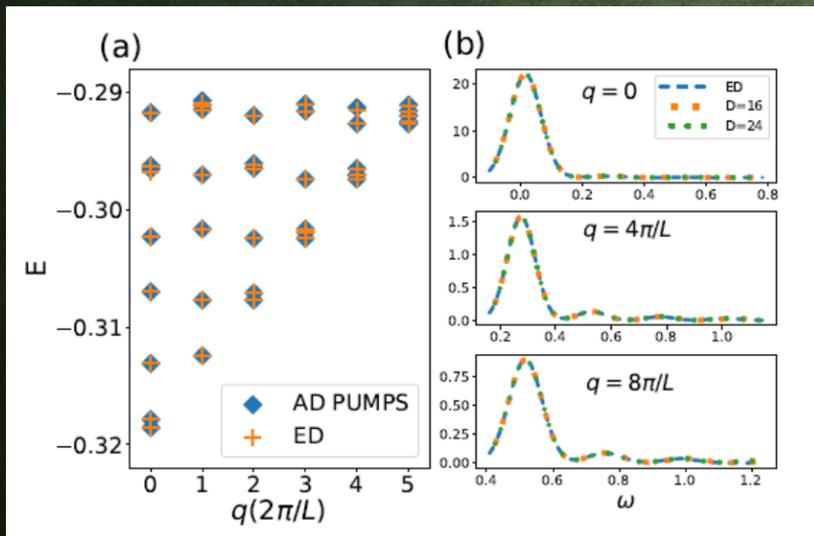


$$\hat{G}_{SF}(\lambda) = \left[\begin{array}{c} \text{---} \text{---} \text{---} \text{---} \text{---} \text{---} \\ | \quad | \quad | \quad | \quad | \\ \text{---} \text{---} \text{---} \text{---} \end{array} \right]$$

$$\hat{O}_j^\beta(\lambda) = I + \lambda e^{-ikr_j} \hat{O}^\beta$$

- With generating functions, now we only need to calculate one or a few tensor graphs.
- Moreover, the derivatives can be evaluated using automatic differentiation(AD), which is often utilized in neural networks.

1D critical Ising chain(L=24):



$$S^\alpha(k, \omega) = \sum_n |M_k^\alpha|^2 \delta(\omega - E_n^k + E_0),$$

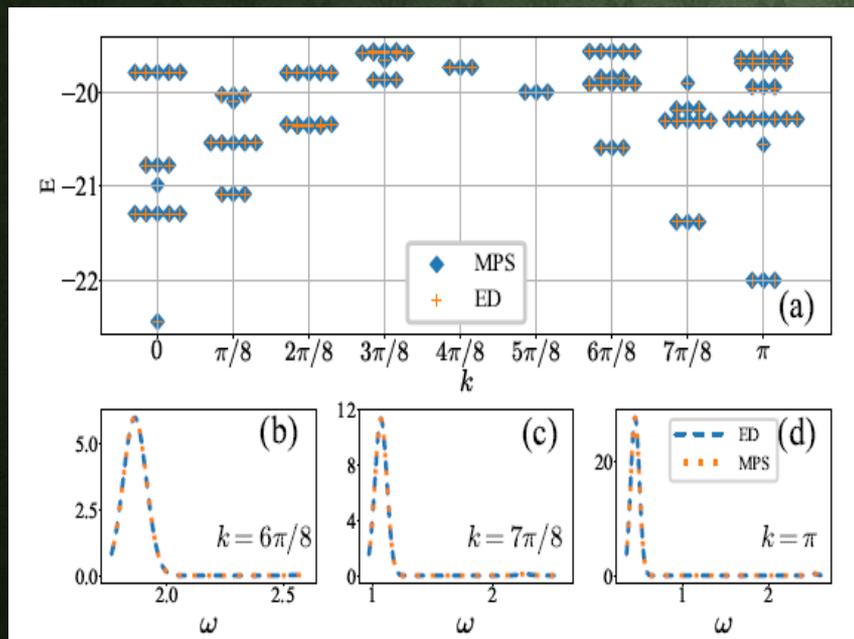
$$M_k^\alpha = \langle \Phi_k(B_n) | S_k^\alpha | \Psi(A) \rangle,$$

$$S_k^\alpha = \frac{1}{\sqrt{N}} \sum_{j=1}^N e^{ikr_j} S_j^\alpha$$

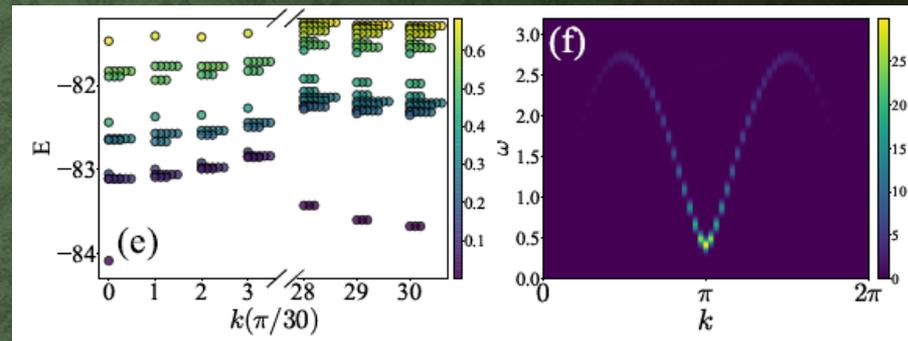
- (a) low-energy spectrum by us and from exact diagonalization(ED).
- (b) spectrum weight(dynamical structural factor) with different bond dimension and ED

1D spin-1 Heisenberg model:

L=16:



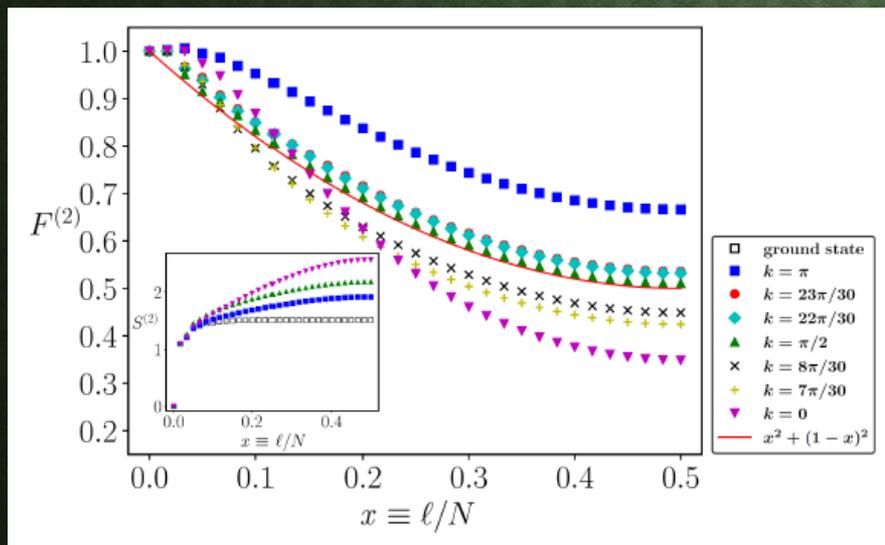
L=60:



- Good benchmark accordance in smaller size with ED.
- Energy spectrum and dynamical structural factor can be obtained for larger system size.
- Haldane gap ≈ 0.4105

1D spin-1 Heisenberg model:

Entanglement entropy:



$$\rho_l^n = \text{Tr}_{\bar{l}} |\psi\rangle\langle\psi| \text{ or } \text{Tr}_{\bar{l}} |\Phi\rangle\langle\Phi|$$

$$S(n) = \frac{1}{1-n} \log \text{Tr}_l \rho_l^n$$

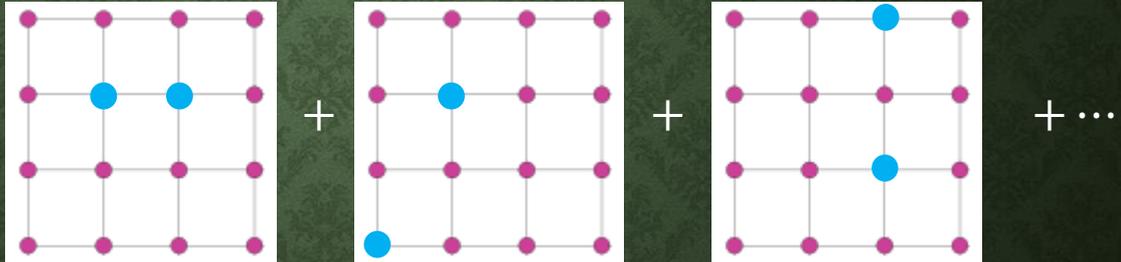
$$F(n) = \frac{\text{Tr} \rho_{\Phi}^n}{\text{Tr} \rho_{\Psi}^n} \quad \begin{array}{l} |\psi\rangle: \text{Ground state} \\ |\Phi\rangle: \text{Excited state} \end{array}$$

- Red line indicates the theoretical prediction that accords with our result for $k = \pi/2$.
- Traditionally this character is hard to calculate with TN due to a large summation of complex TN graphs.

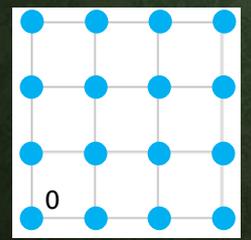
Future goals:

- Although we only demonstrate translationally invariant 1D systems, the idea of generating function can be in principle extended to other scenarios such as in two dimension.
- Moreover, even when we try to calculate the ground state energy for some long-range interactive model, generating function can be also applied.

$$H = \sum_{ij} S_i^z S_j^z$$

$\langle H \rangle =$


 $N(N-1)/2$ terms

$G(H) =$


 Only N terms!!

$\langle H \rangle = \frac{1}{2} \frac{\partial G(H)}{\partial \lambda} \Big|_{\lambda=0}$

+ other terms with different reference point


 S_0^z

 $I + \lambda S_i^z$

- It might be useful for complex models of quantum chemistry

SUMMARY

- We borrow the idea of generating function originating from field theory and apply it to the tensor network ansatz.
- Our results indicated that this is an useful approach and reduces the obstacle while performing TN summation.
- Possible extension includes its application in 2D systems, and we may also utilize the generating function for long-range interactive models.